

# THz near-field microscopy technique and applications with sub-wavelength aperture probes

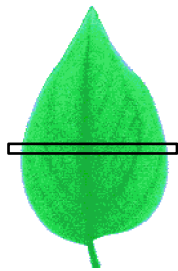
Oleg Mitrofanov

*University College London*

**[o.mitrofanov@ucl.ac.uk](mailto:o.mitrofanov@ucl.ac.uk)**

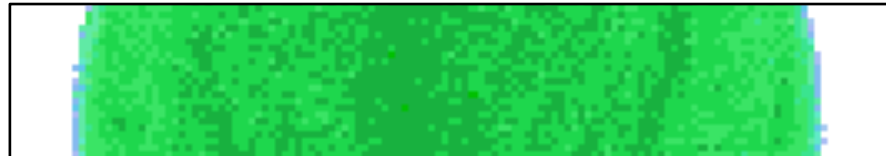
1. *Introduction: some near-field optics concepts*
2. *Historical perspective: THz imaging technology*
3. *THz imaging applications:*  
*E-field mapping in space and time*
4. *THz surface waves and metallic resonators*
5. *Aperture-type THz near-field probes*





Hu and Nuss  
(1995)

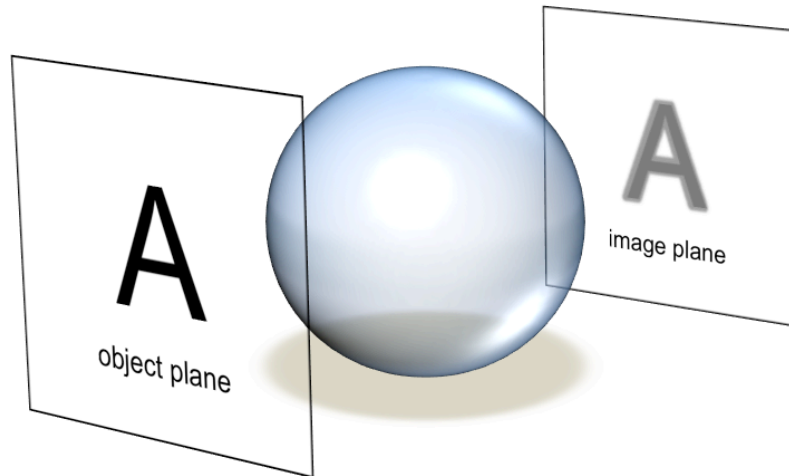
$$\lambda = 300 \mu\text{m}$$



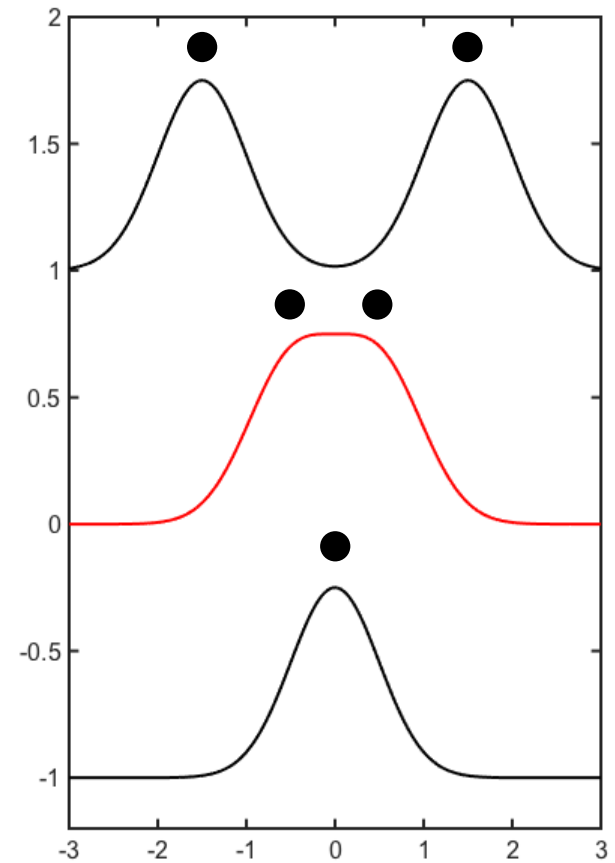
$$d = \frac{\lambda}{2n \sin \alpha}$$

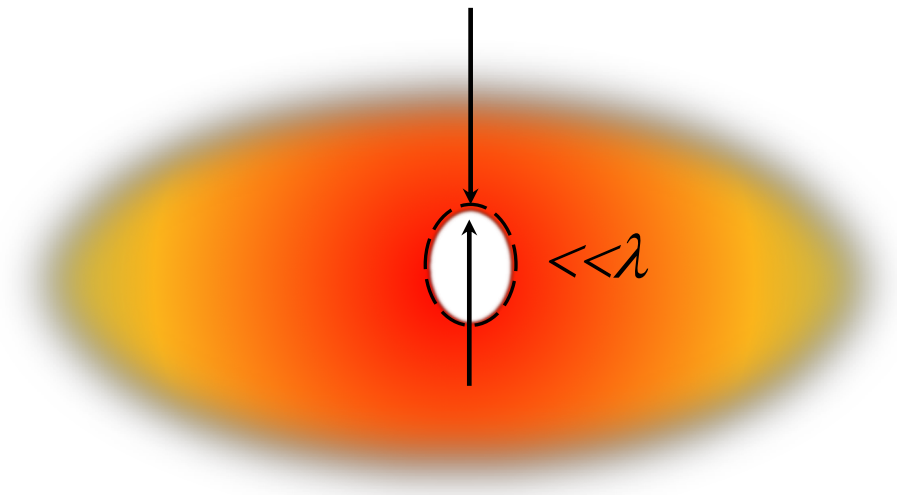
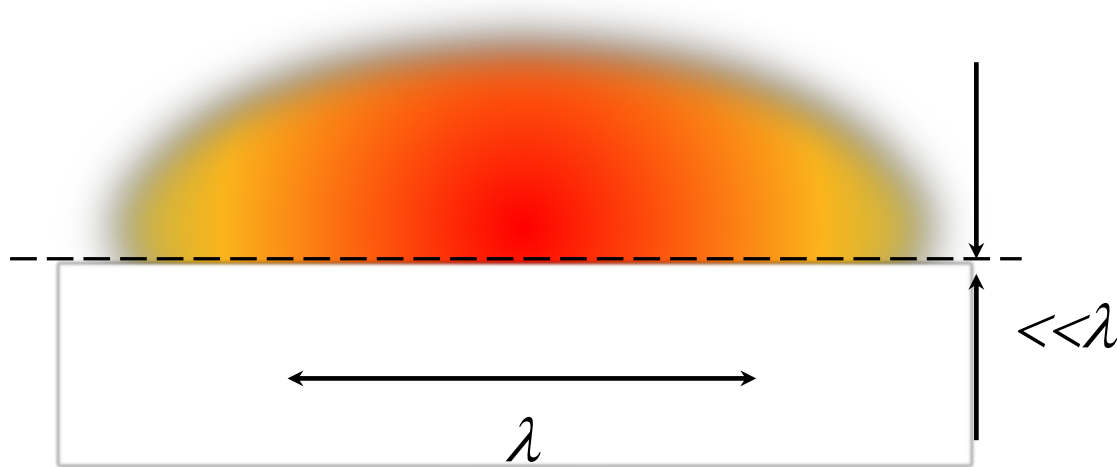
Ernst Abbe (c. 1873)

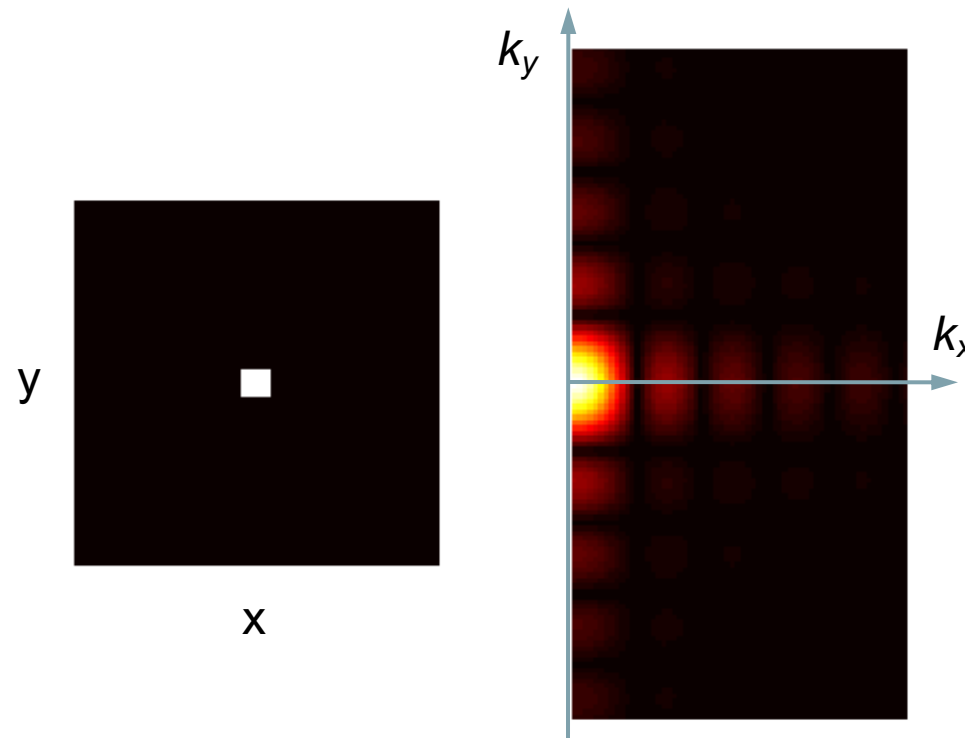
## Point-spread function of imaging system



Resolving two points







Spatial distribution

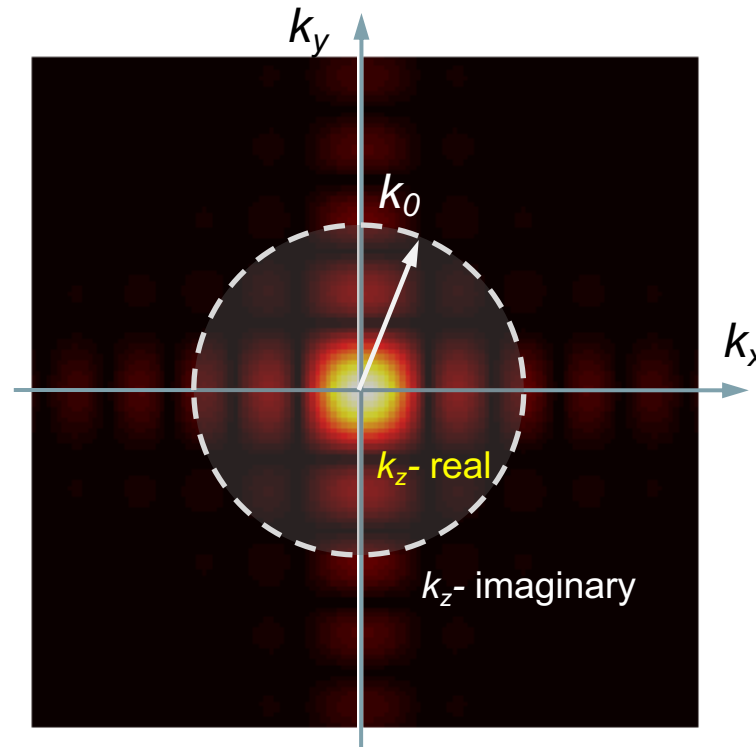


Spatial frequency distribution

*Fourier Transform*

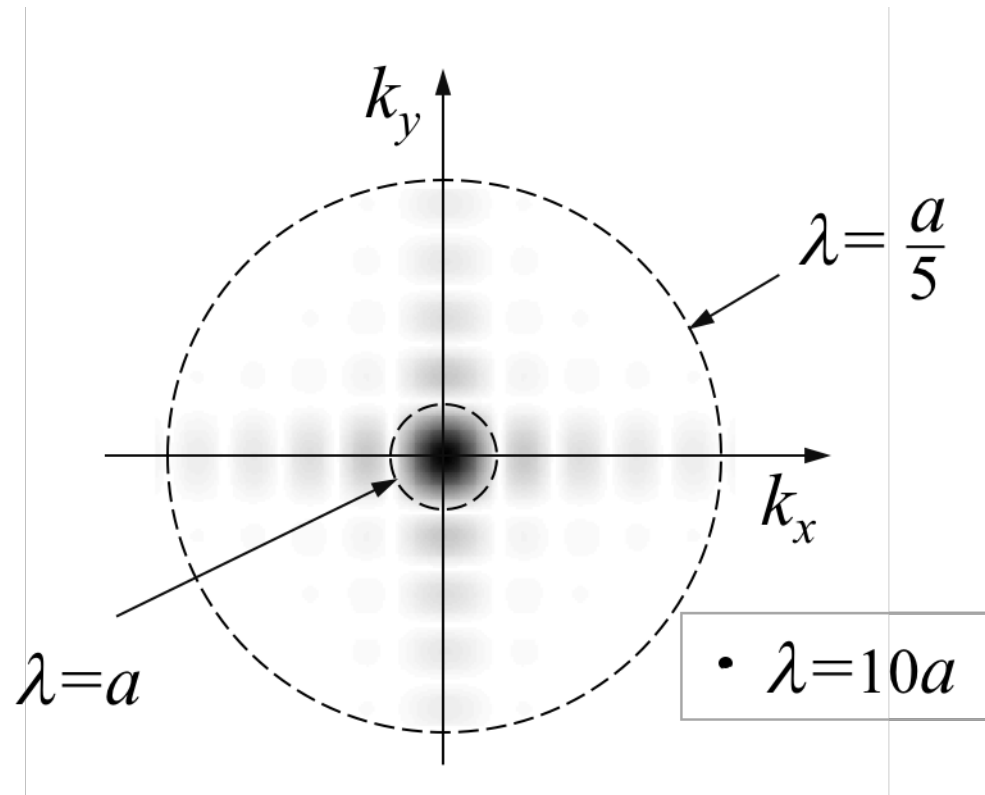
$$E = E_0 \exp(ik_0 r - i\omega t)$$

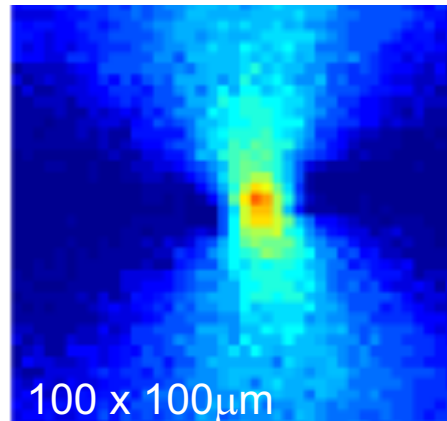
$$k_0^2 = k_x^2 + k_y^2 + k_z^2$$



Imaginary  $k_z$  components – evanescent waves

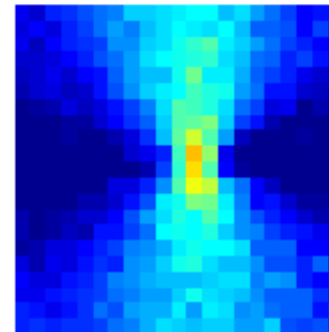
- localized near surface



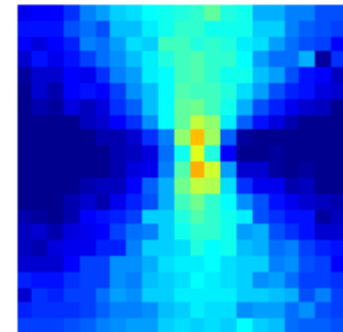


$\lambda \sim 200 \mu\text{m}$

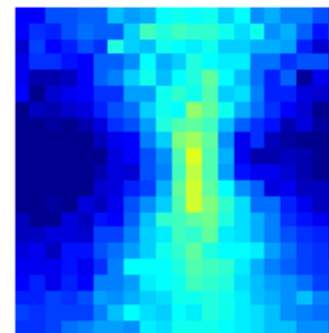
$\Delta z = 0 \mu\text{m}$



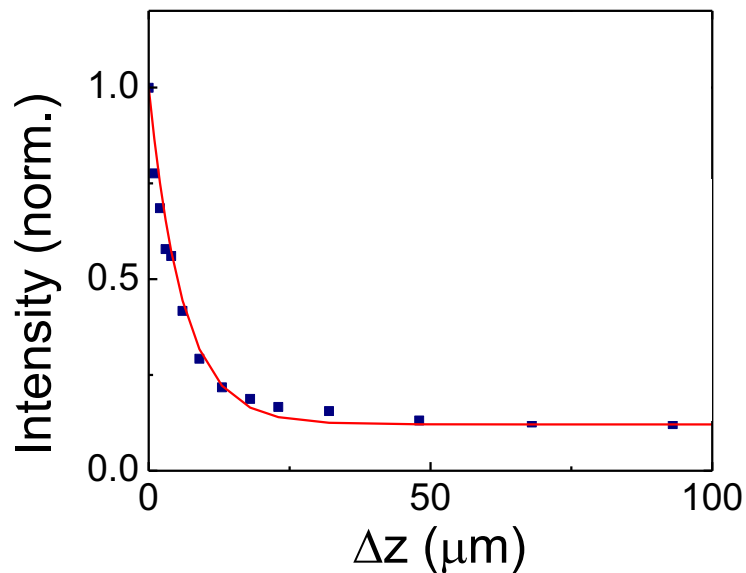
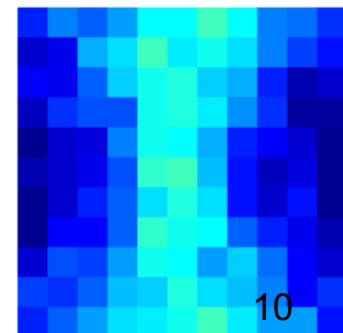
$\Delta z = 1.5 \mu\text{m}$



$\Delta z = 4 \mu\text{m}$



$\Delta z = 8 \mu\text{m}$

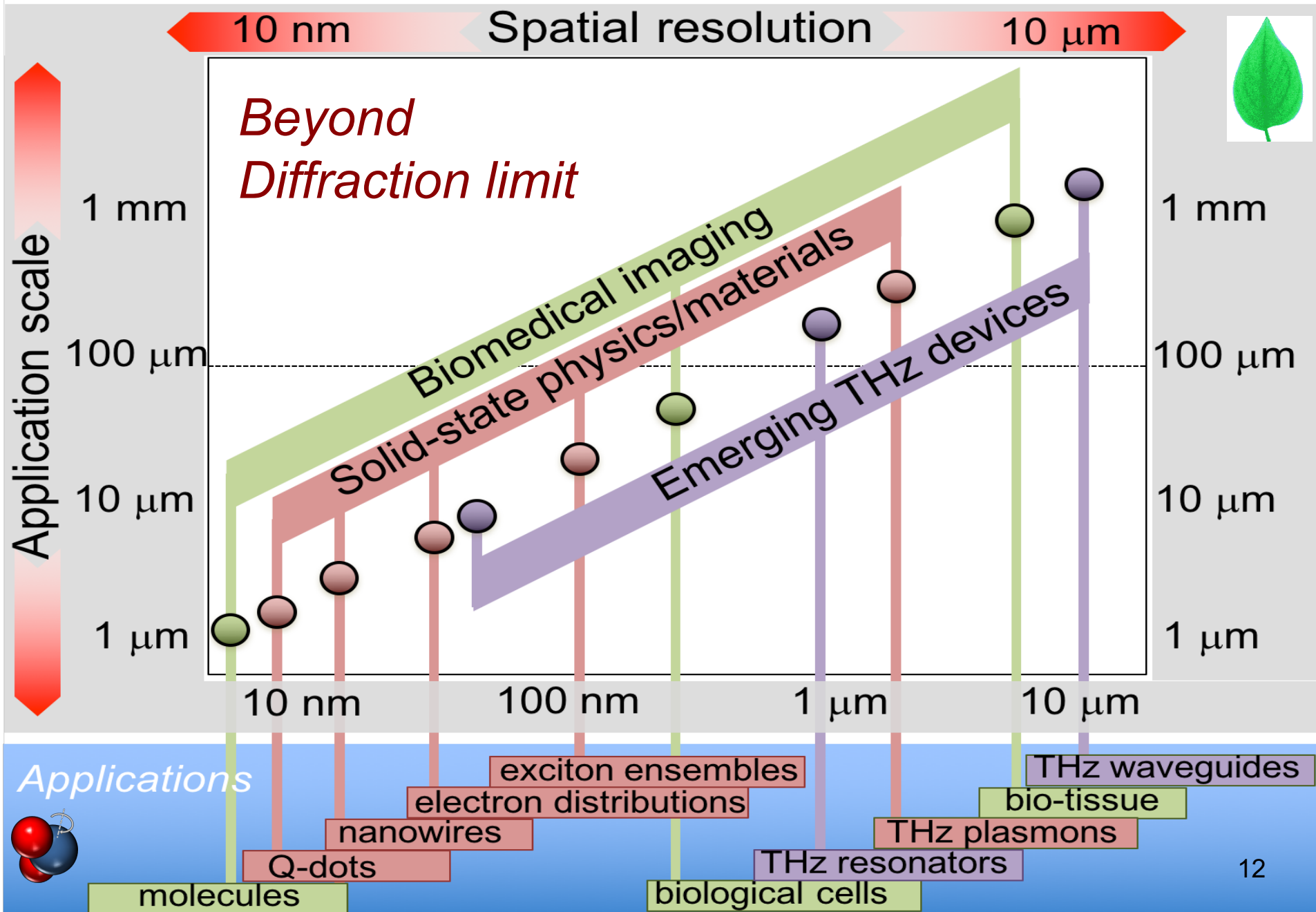


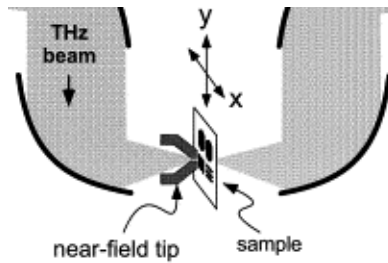


Breaking the diffraction limit in the THz range

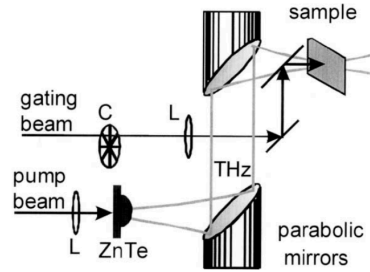
$$\frac{\lambda}{r} = 1 - 100,000$$

# THz Near-field Imaging Technology and Applications

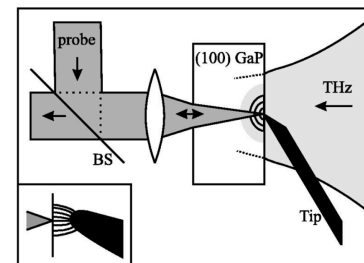




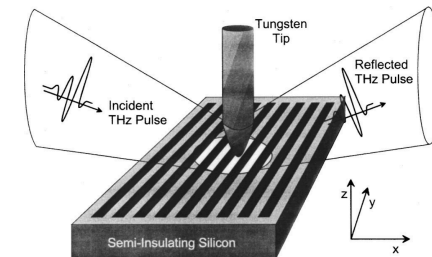
Hunsche et al., 1998



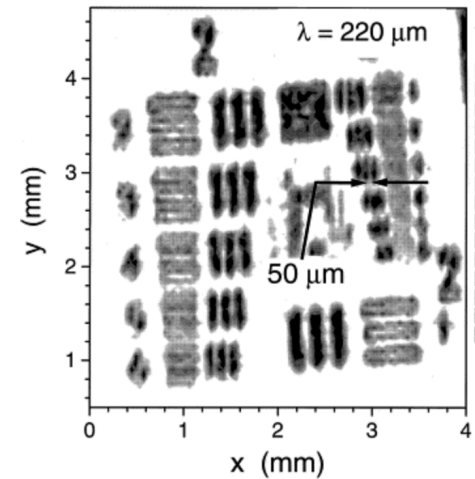
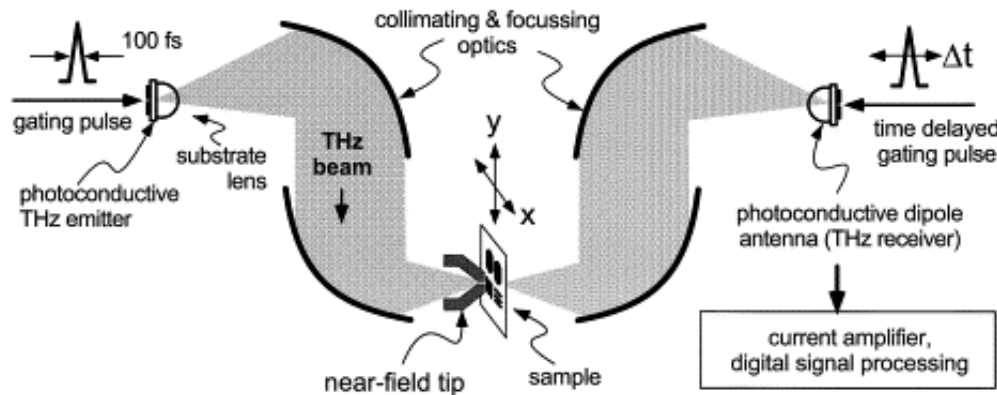
Chen et al., 2000



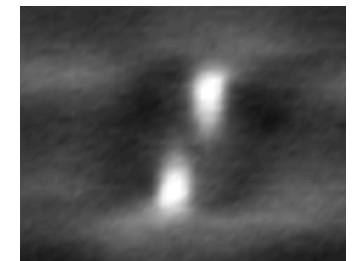
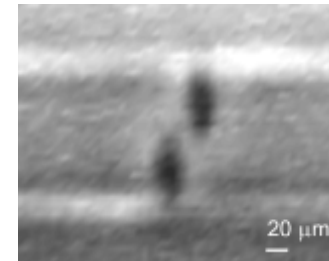
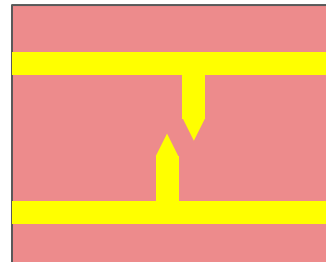
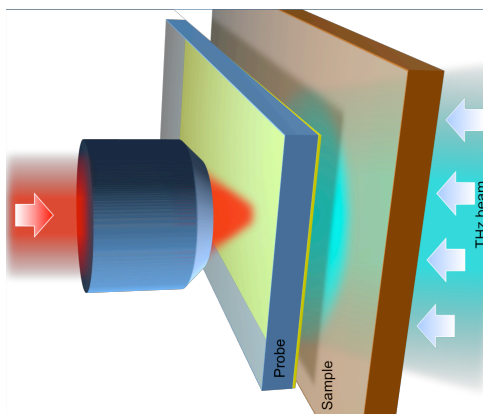
Van der Valk et al., 2002



Chen et al., 2003

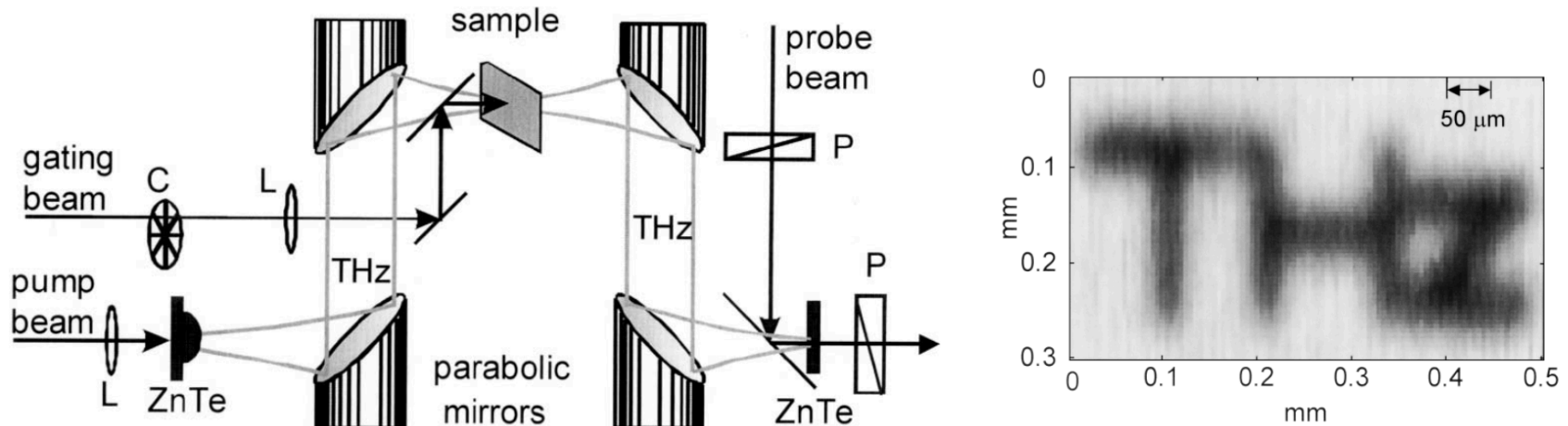


Hunsche et al. *OPTICS COMM* (1998)



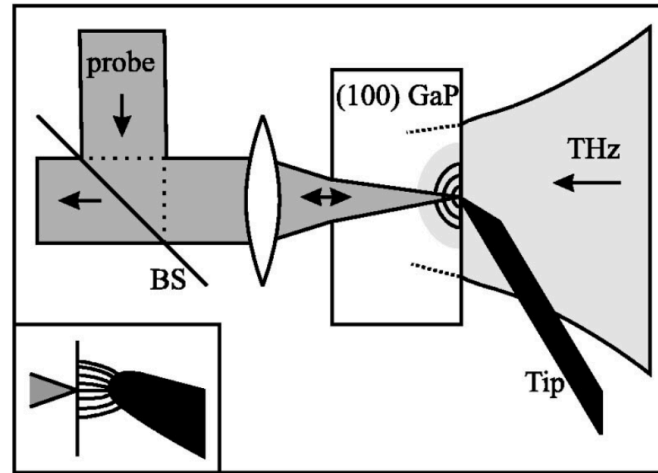
Mitrofanov et al. *J. STQE* **103**, 600 (2001)

## Near-field THz source modulation - 'Dynamic Aperture' (THz-TDS)



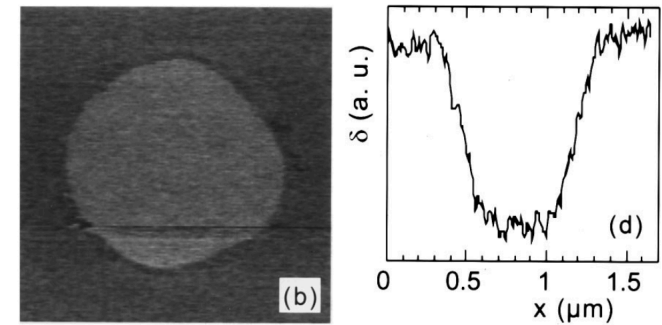
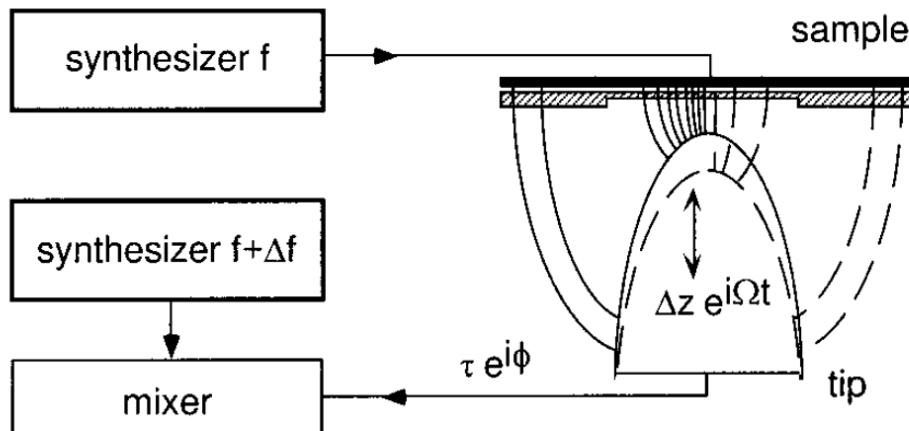
Chen et al. *OPTICS LETTERS* 25 (2000)

## Electro-optic near-field probes (THz-TDS)



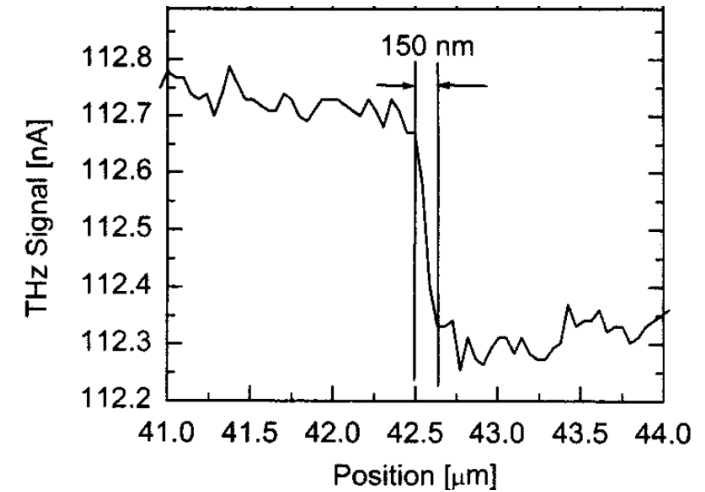
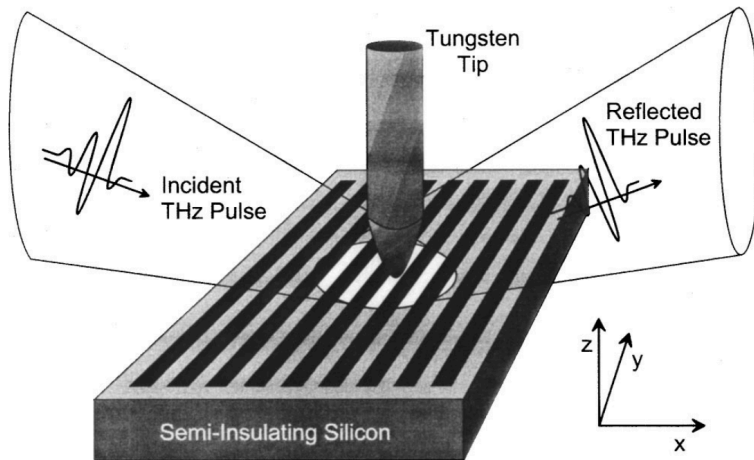
Van der Valk et al., APPL. PHYS. LETT. (2002)

## STM Tip - Microwave transmission (1 GHz)



B. Knoll, F. Keilmann et al., APPL. PHYS. LETT. 70, 2667 (1997)

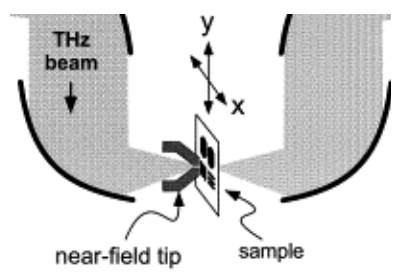
## Scattering Tip near-field microscopy (THz-TDS)



Chen et al., APPL. PHYS. LETT. 83, 3009 (2003)



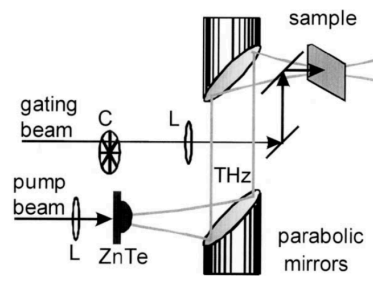
# Progress in development of THz microscopy



Hunsche et al., 1998

Near-field probes with integrated THz detectors

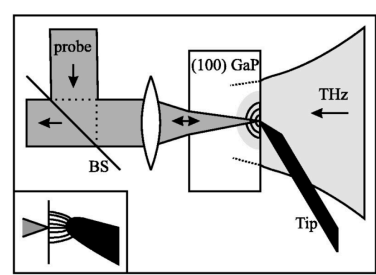
Sub-wavelength THz generation



Chen et al., 2000

Use of patterns instead of apertures

Signal processing: adaptive imaging and compressive sensing

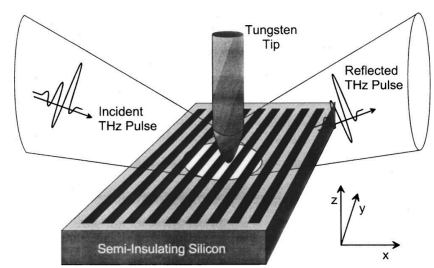


Van der Valk et al., 2002

EO materials/ultrathin crystals

High-E THz sources

Spectral filtering

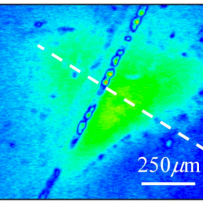
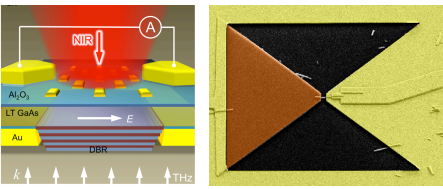


Chen et al., 2003

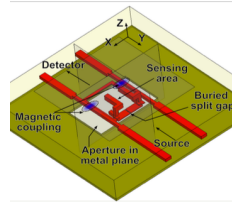
Higher order modulation techniques

Surface plasmons

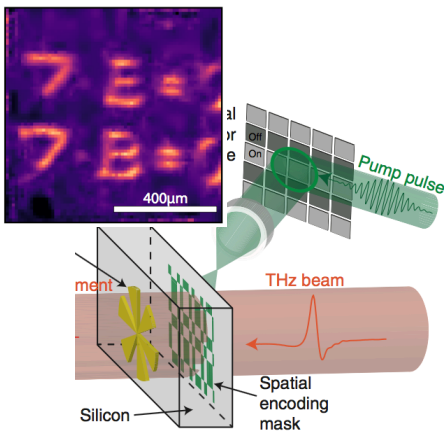
Detection of a THz driven tunneling current



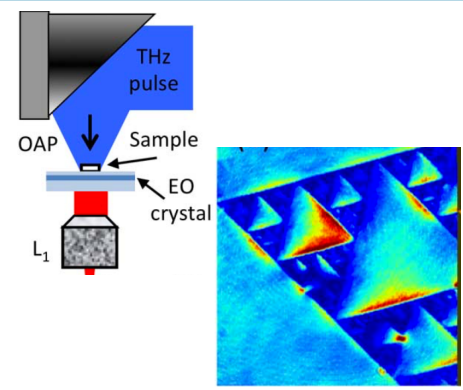
Serita (2012)



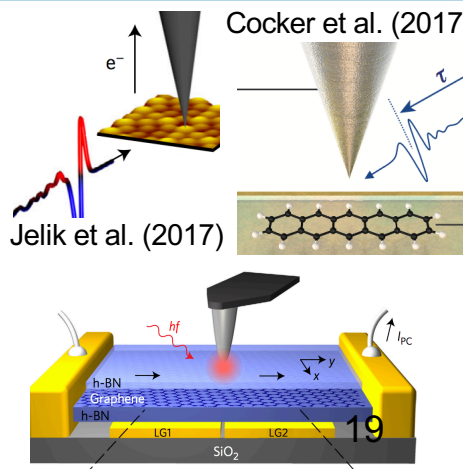
Grzyb (2016)



Rayko et al., 2016



Blanchard & Tanaka, 2016



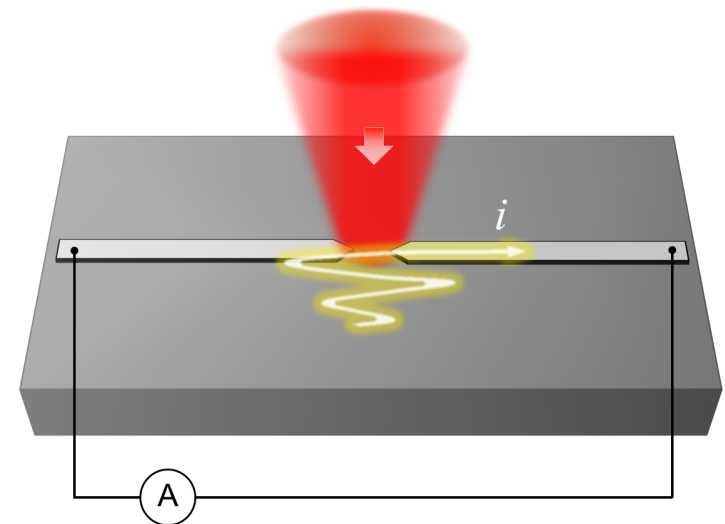
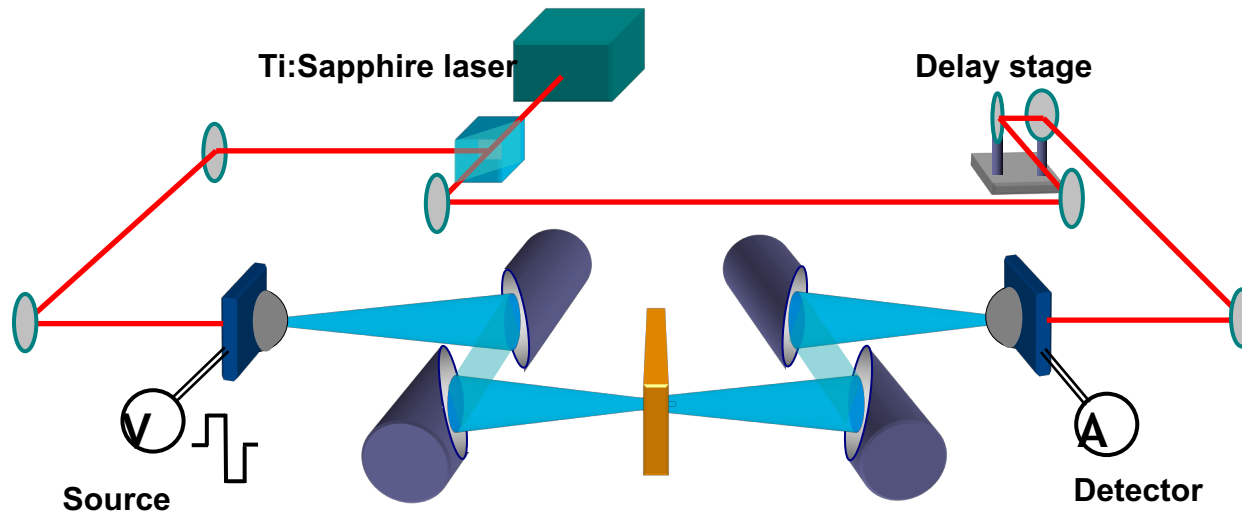
Alonso-Gonzalez et al. (2017)

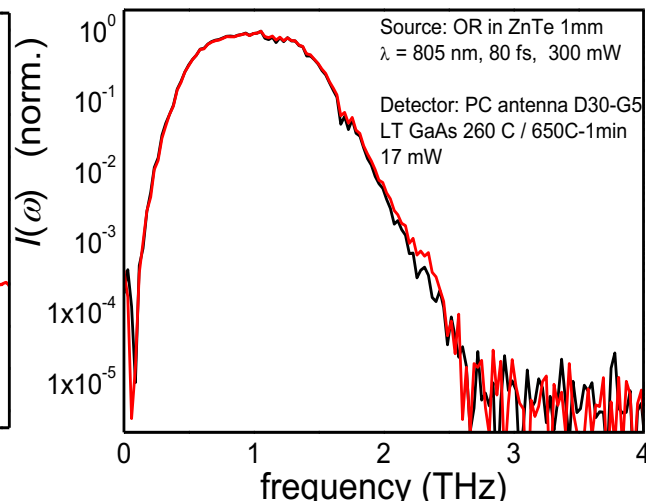
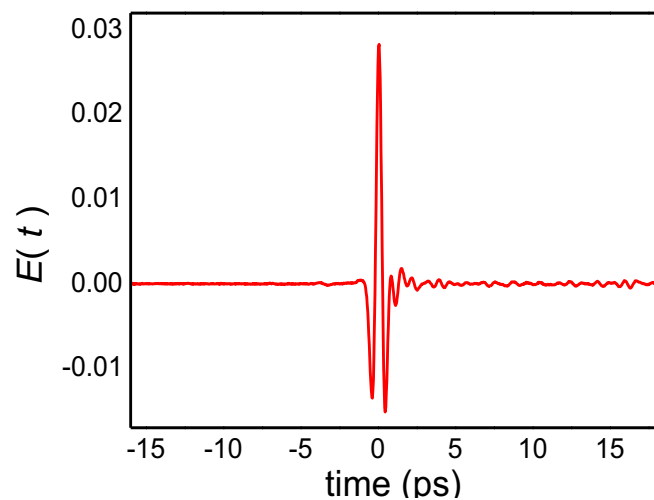
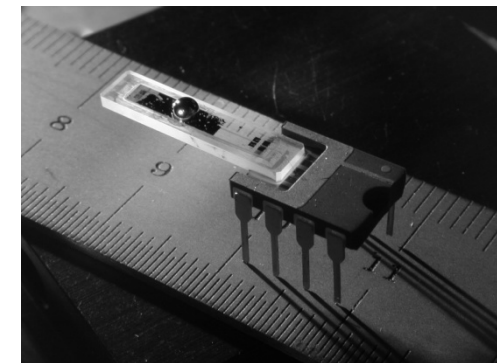
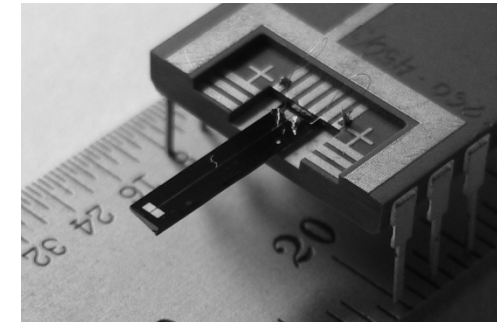
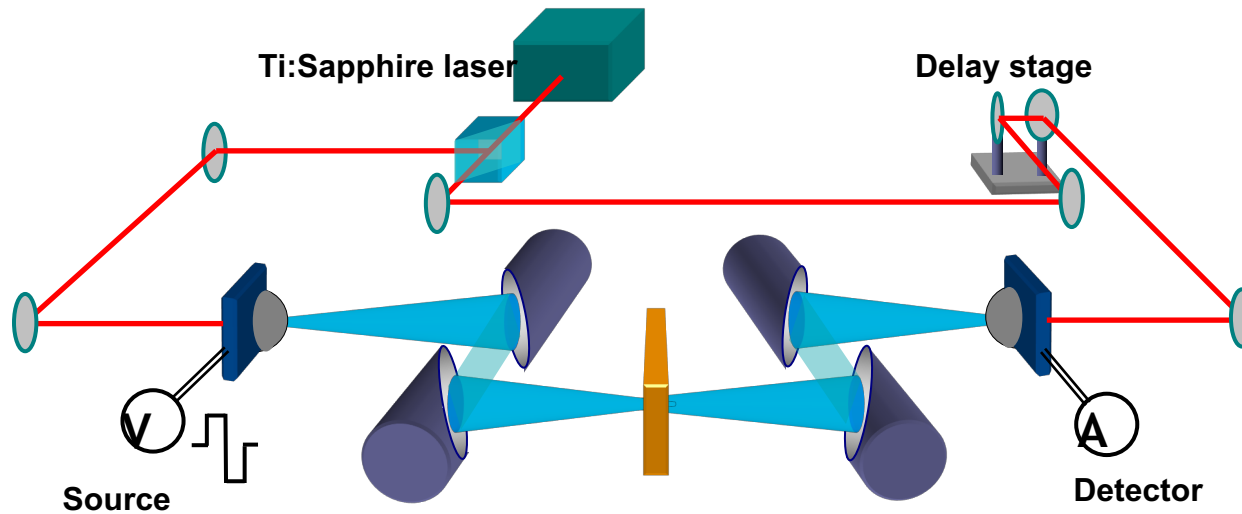
Cocker et al. (2017)

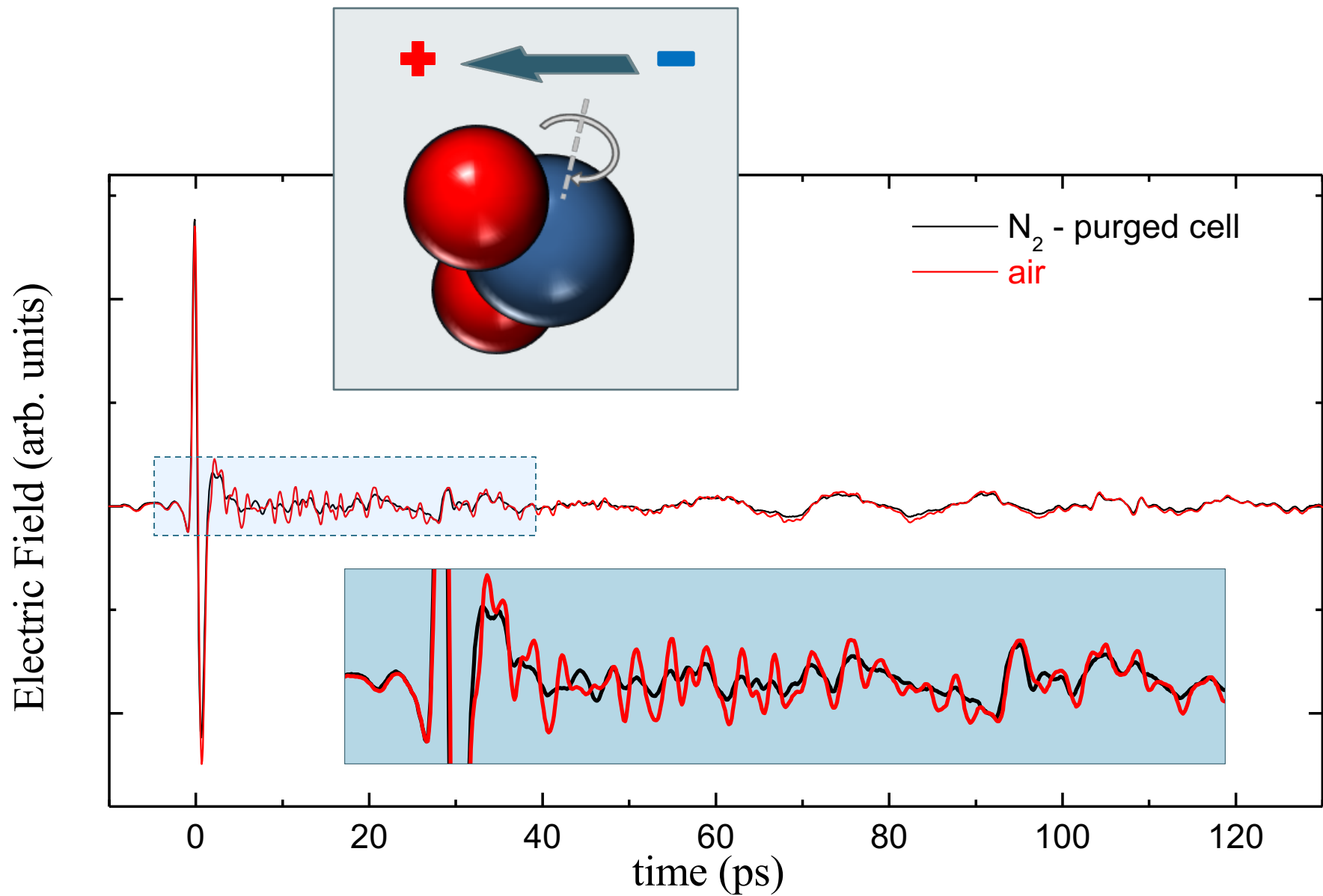
Jelik et al. (2017)

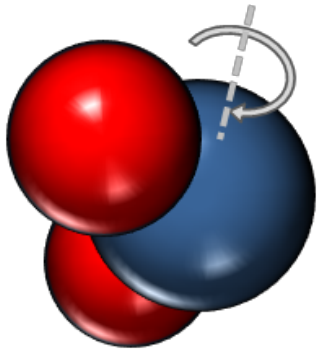
*THz imaging applications:*

*E-field mapping in space and time*



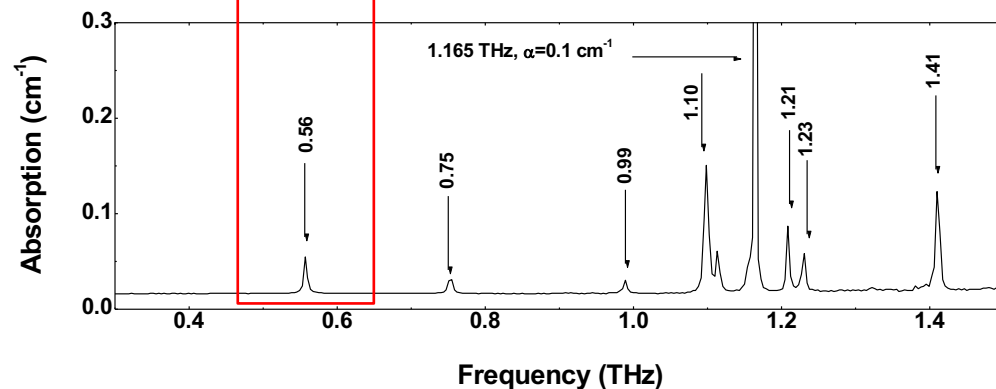
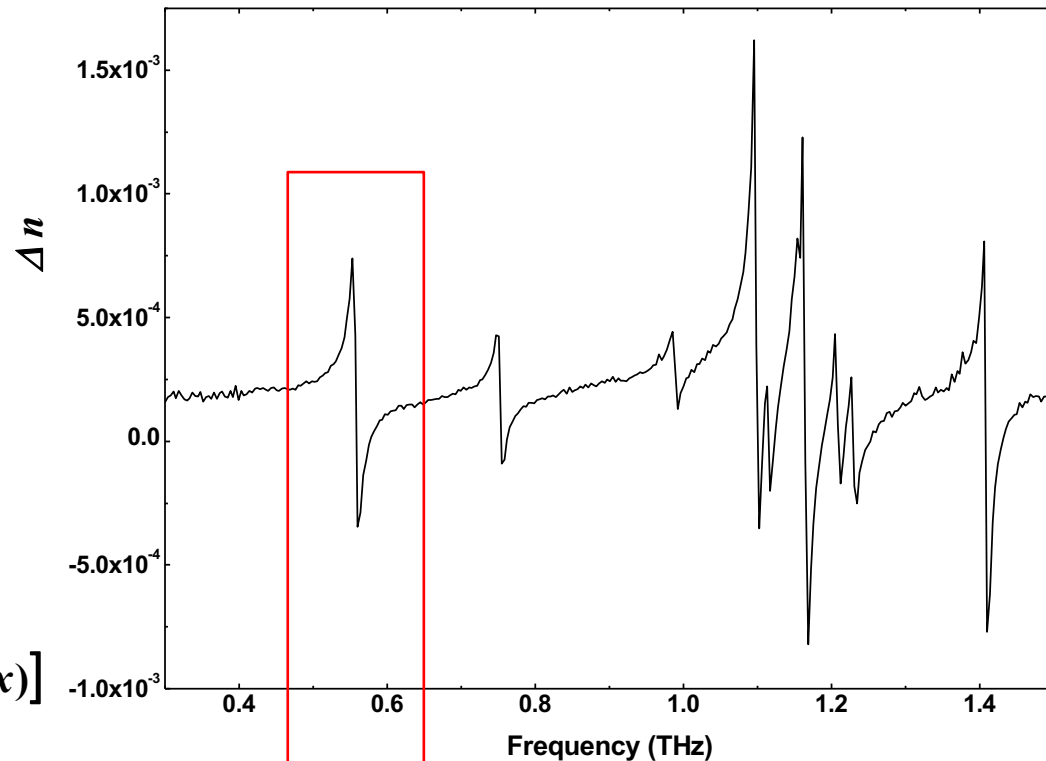


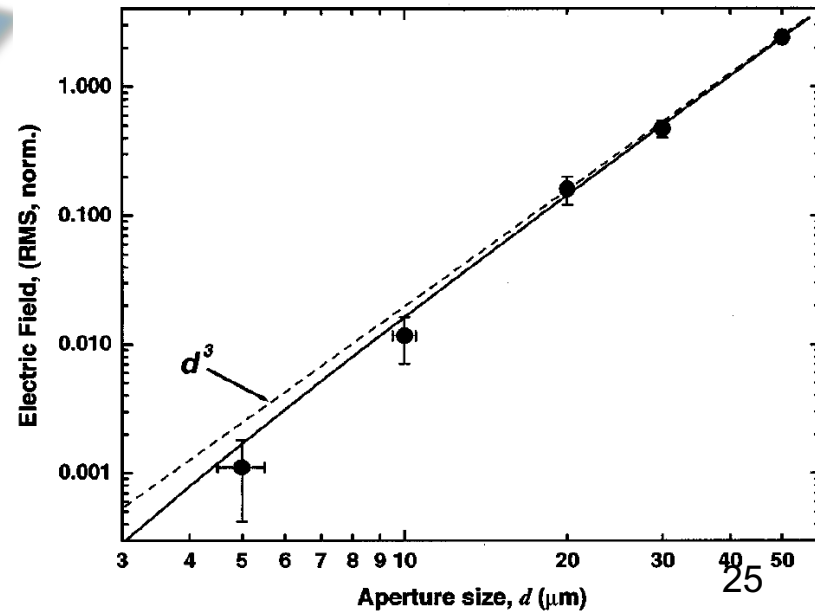
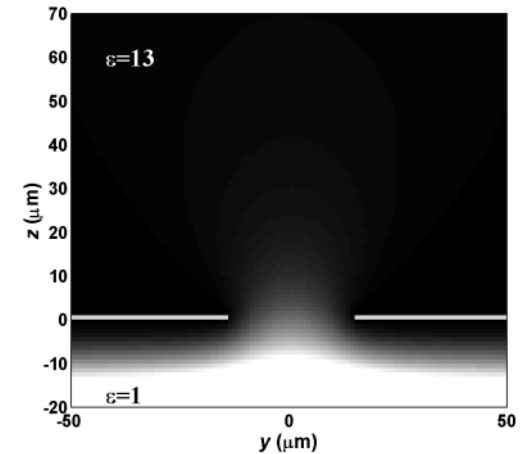
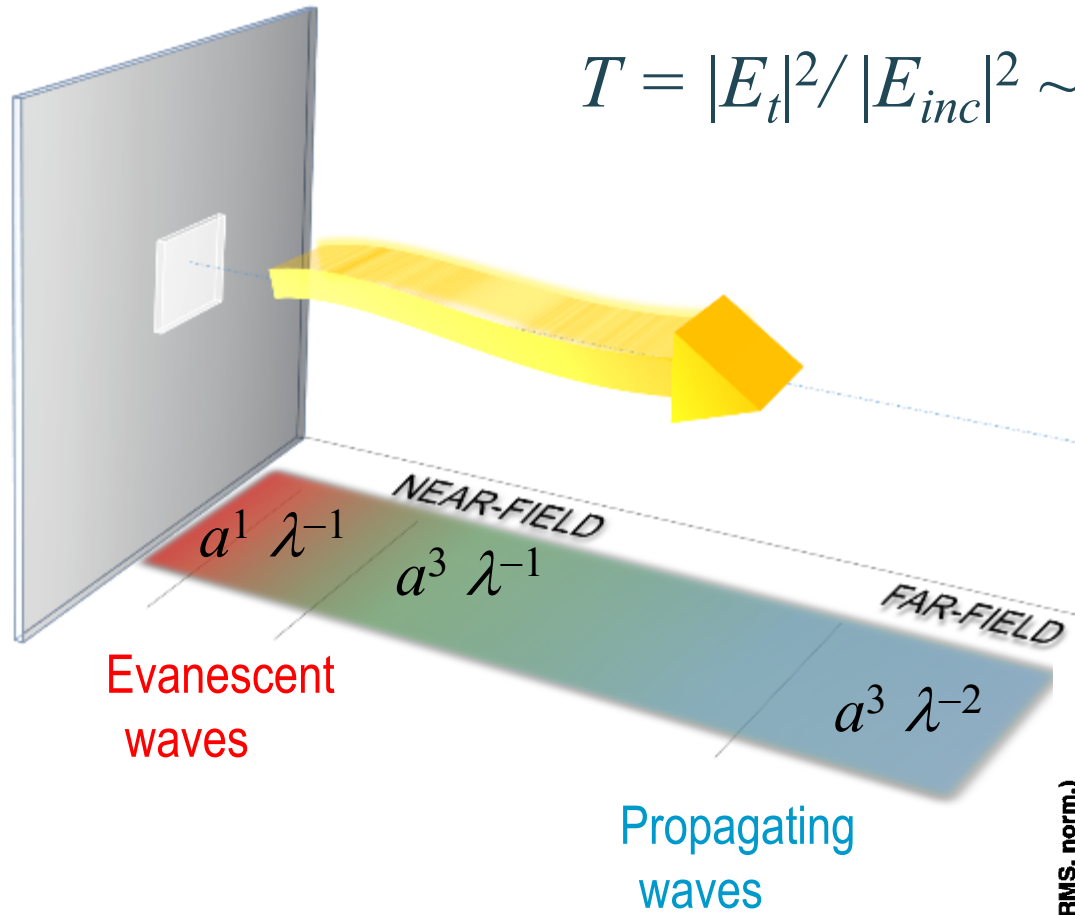




$$\text{Re}[\hat{n}(\omega)] = \frac{c}{\omega l} [\varphi(\omega, x+l) - \varphi(\omega, x)]$$

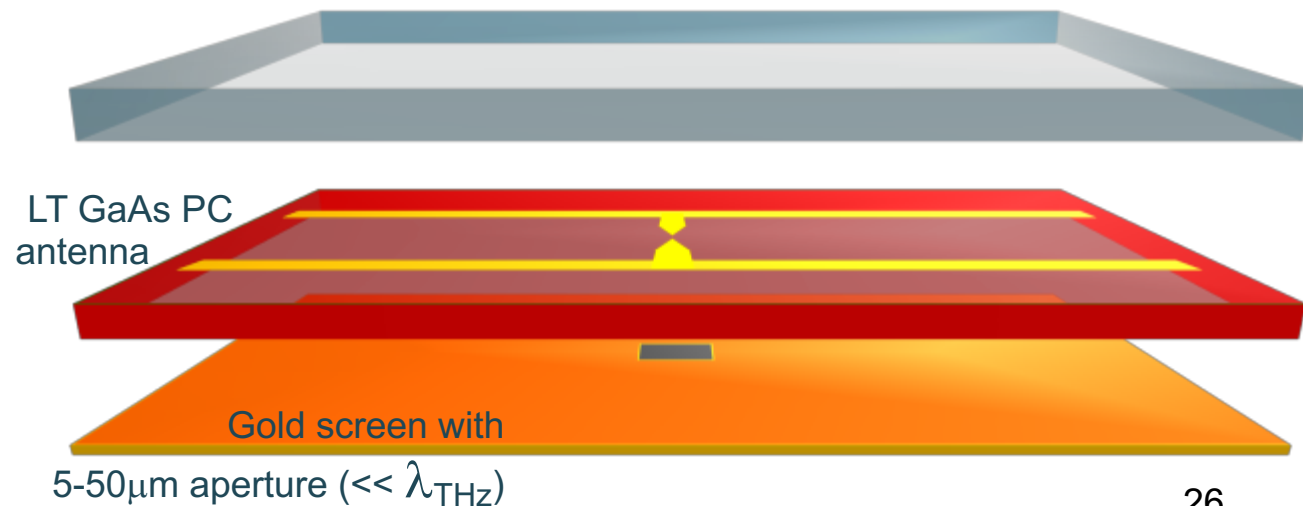
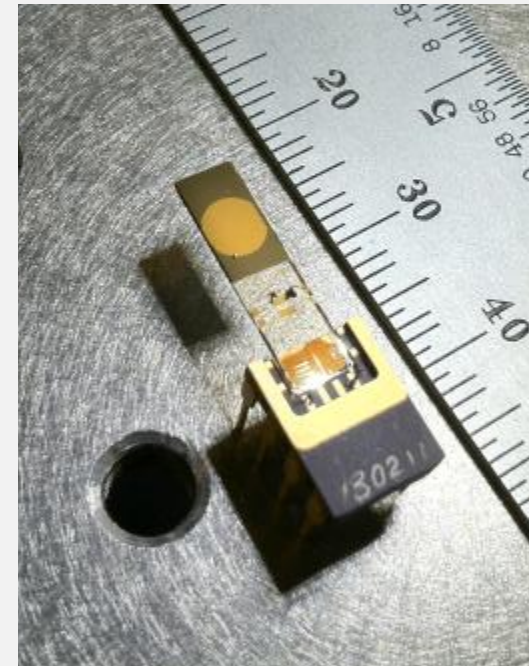
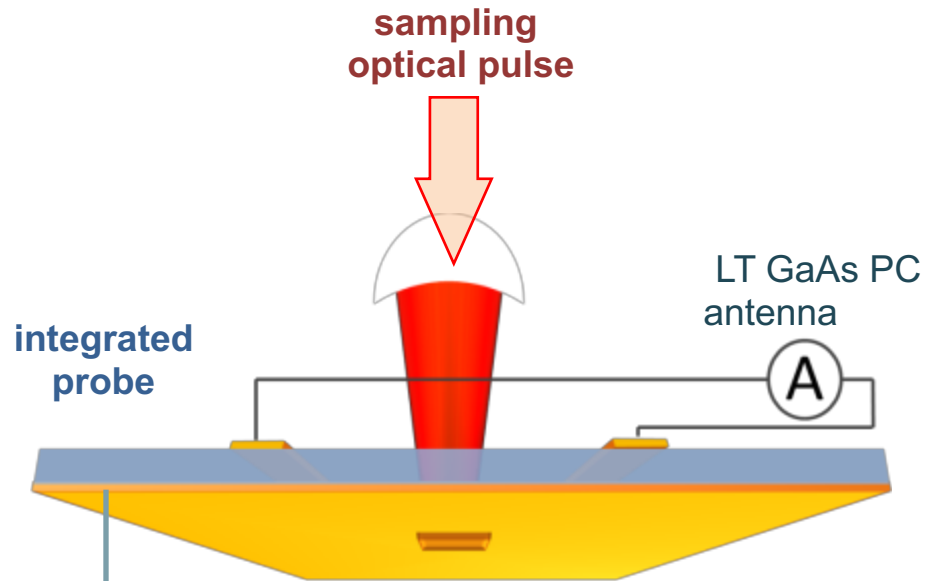
$$\text{Im}[\hat{n}(\omega)] = \frac{c}{\omega l} \left[ \ln \left( \frac{E(\omega, x+l)}{E(\omega, x)} \right) \right]$$



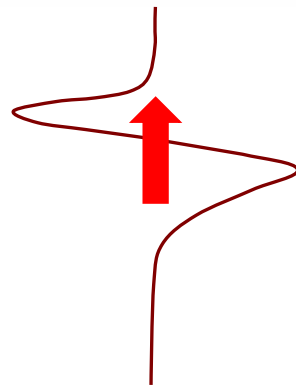
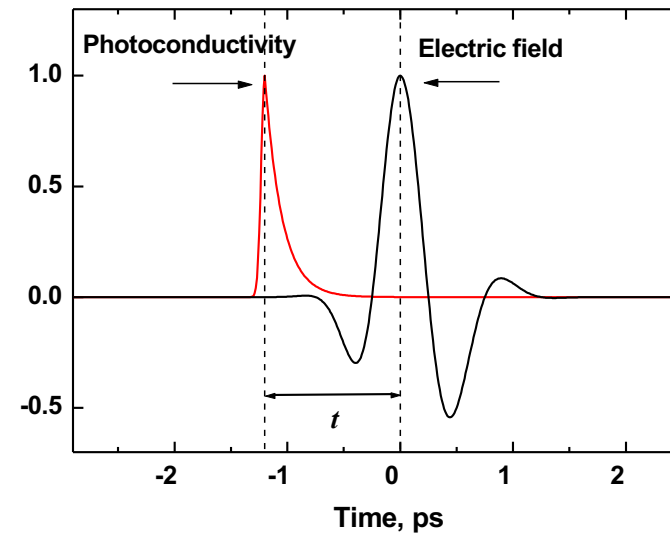
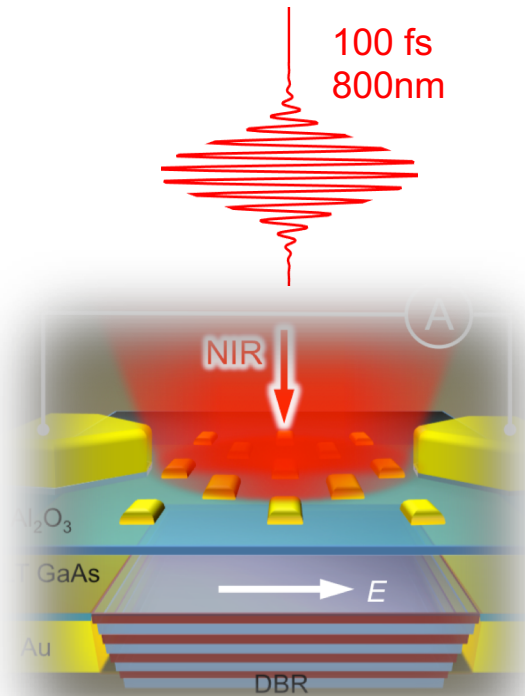


Mitrofanov et al., APL **77**, 3496 (2000)  
APL **79**, 907 (2001)

Adam, J IRMTW **32**, 976 (2011)



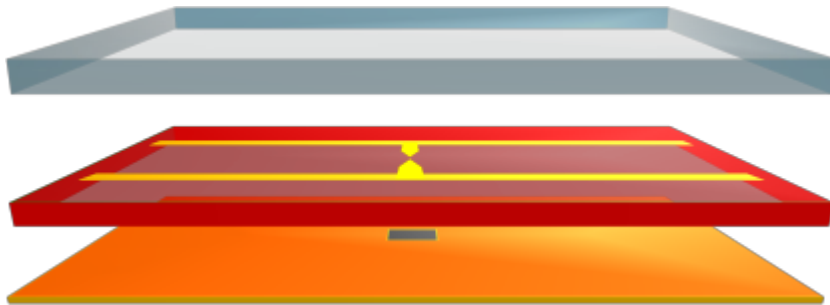




$$I_A(t) = \int \sigma(t' - t) E_{THz}(t') dt'$$

$$E(t), E(\omega), \phi(\omega)$$

material:	LT GaAs	
gating pulse:	$\lambda_c = 800 \text{ nm},$	$\tau = 100 \text{ fs}$
Bandwidth:	3 THz	27

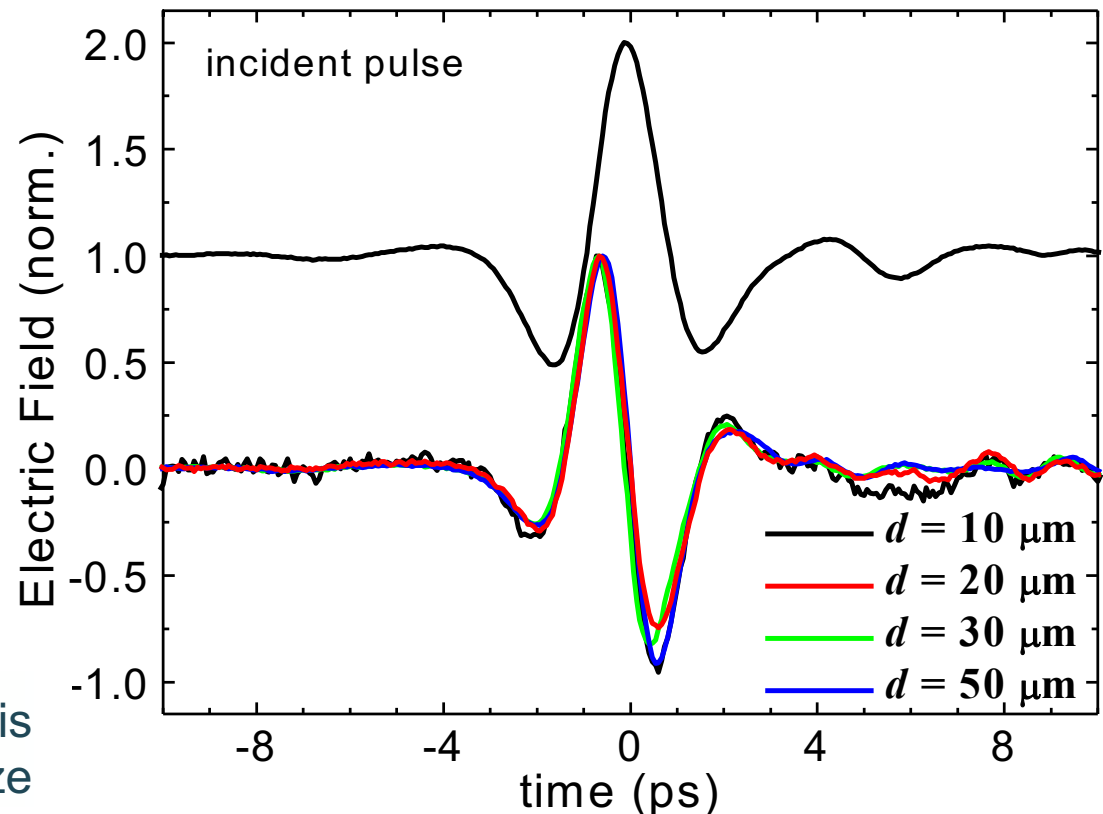


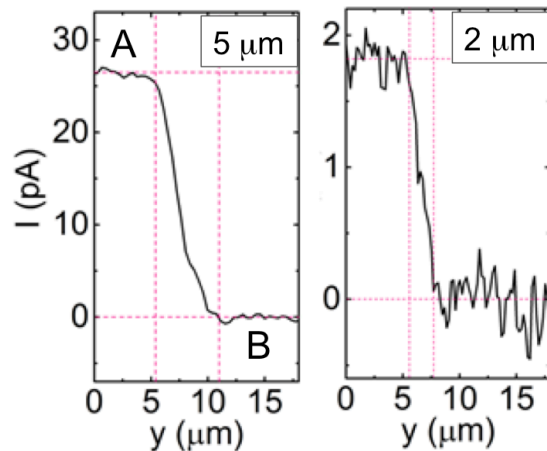
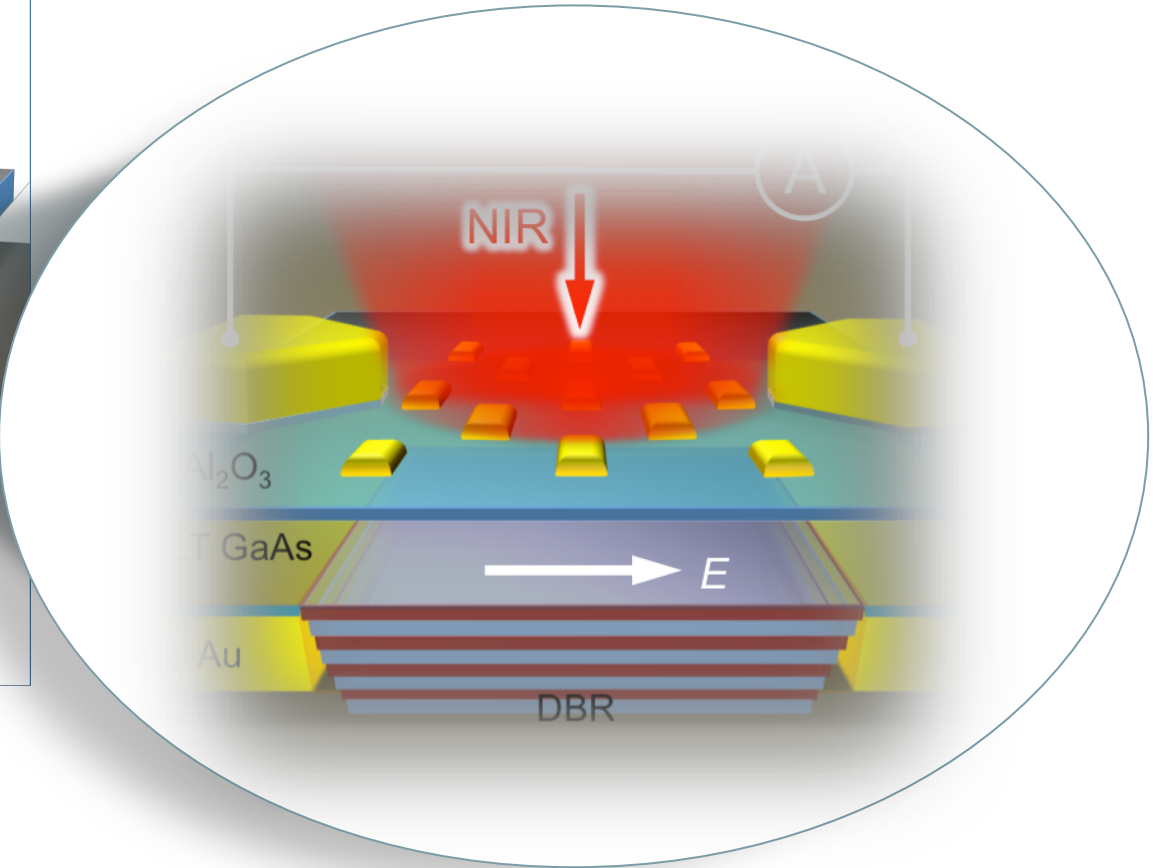
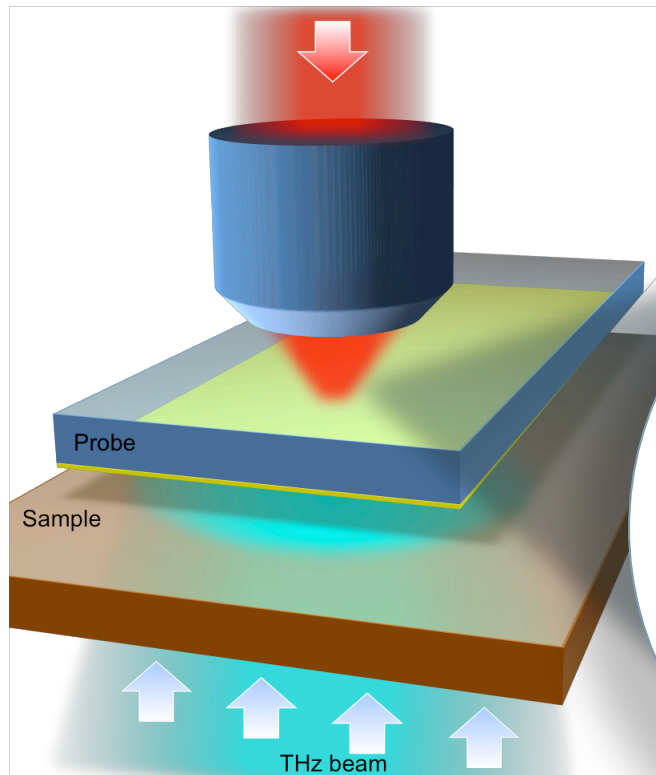
Aperture much smaller  
than the wavelength

$$E(\omega) \sim \omega$$

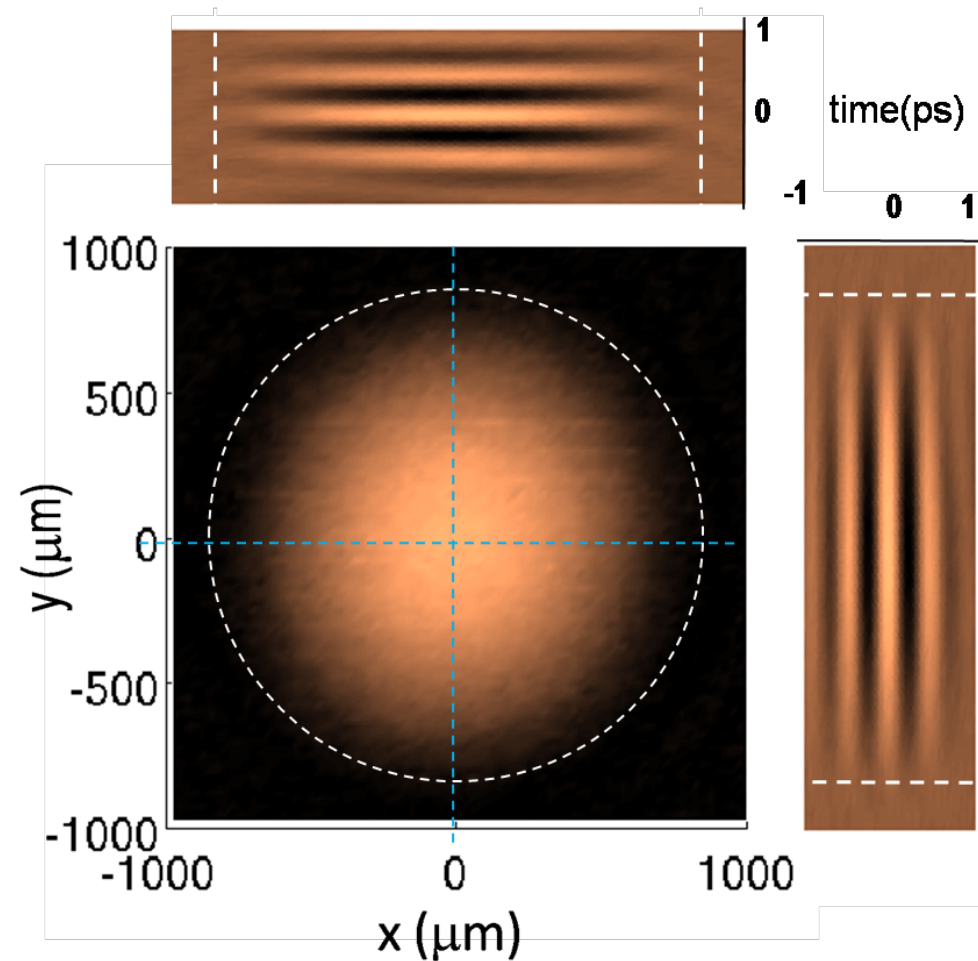
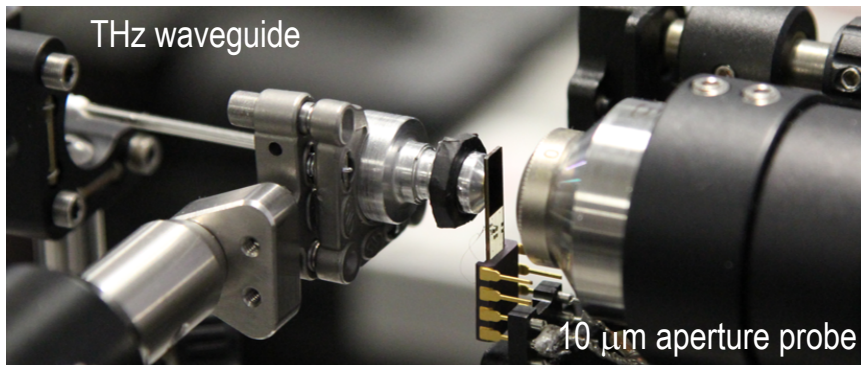
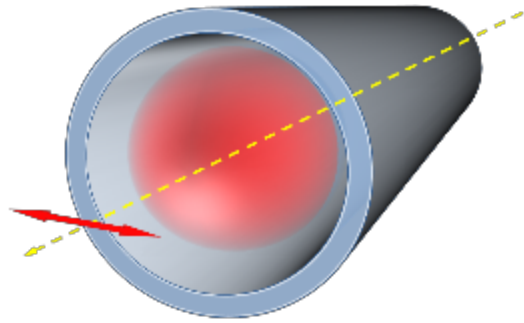
$$\phi(\omega) = \pi/2$$

Waveform deformation is  
independent of the aperture size

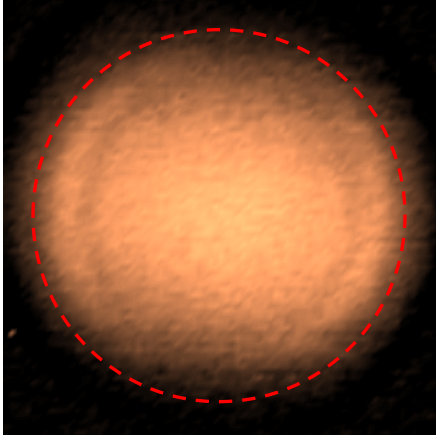




*ACS Photonics* (2015)

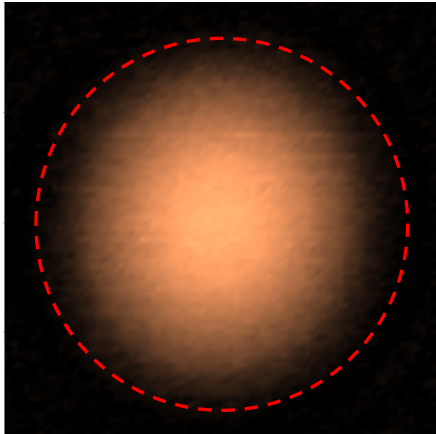


## Dominant mode profiles



**TE<sub>11</sub>**

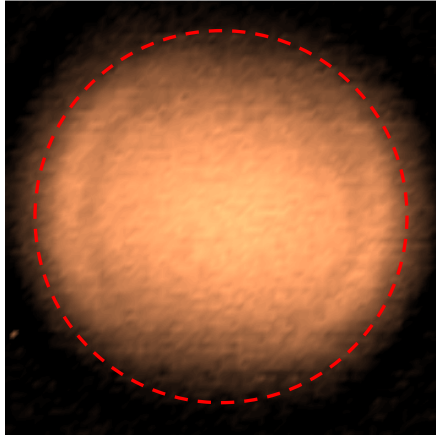
hollow metallic waveguide



**HE<sub>11</sub>**

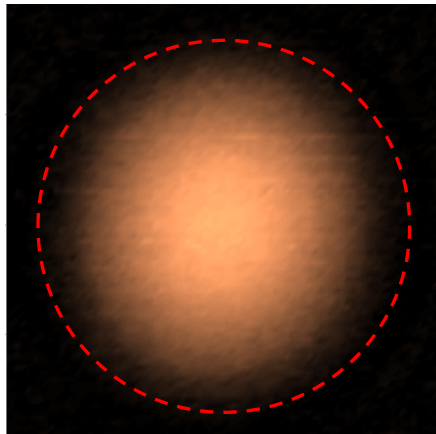
similar waveguide with a PS inner coating

## Dominant mode profiles



$TE_{11}$

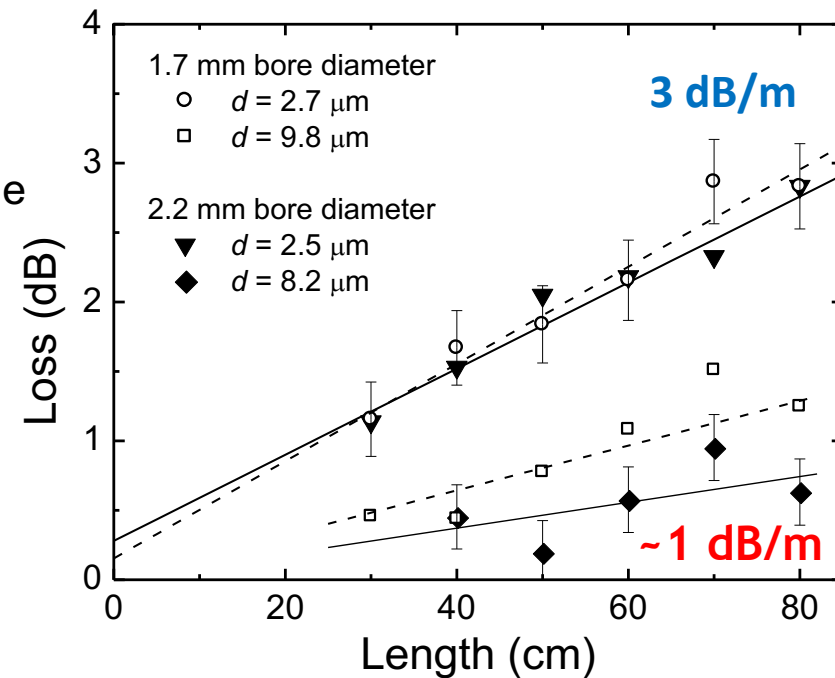
hollow metallic waveguide

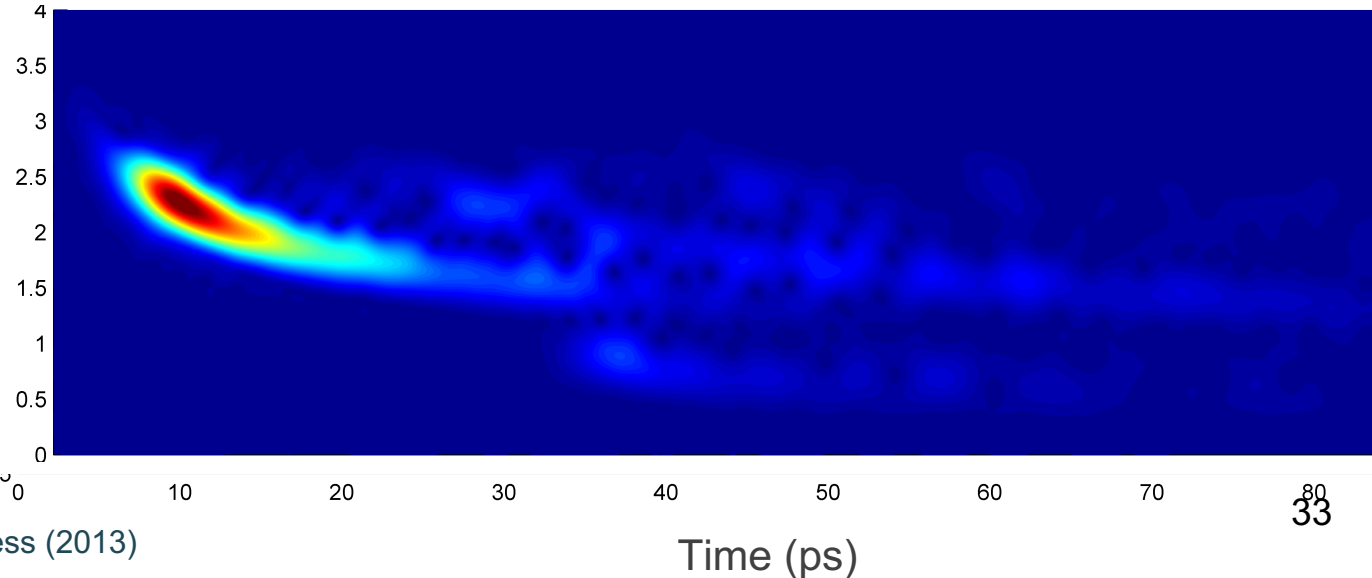
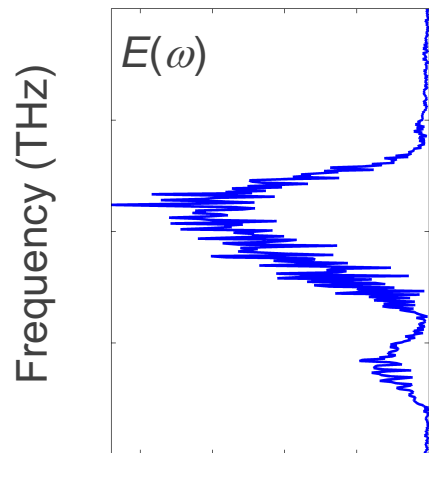
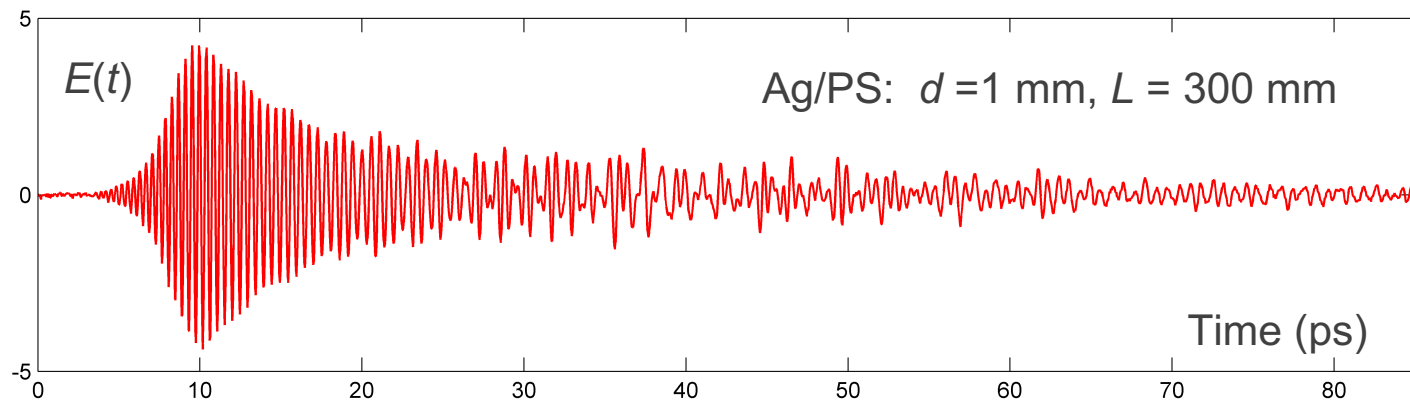
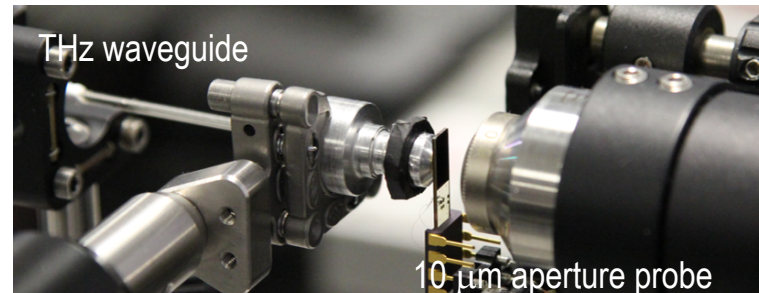
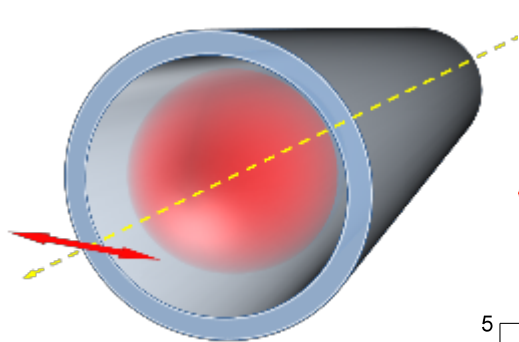


$HE_{11}$

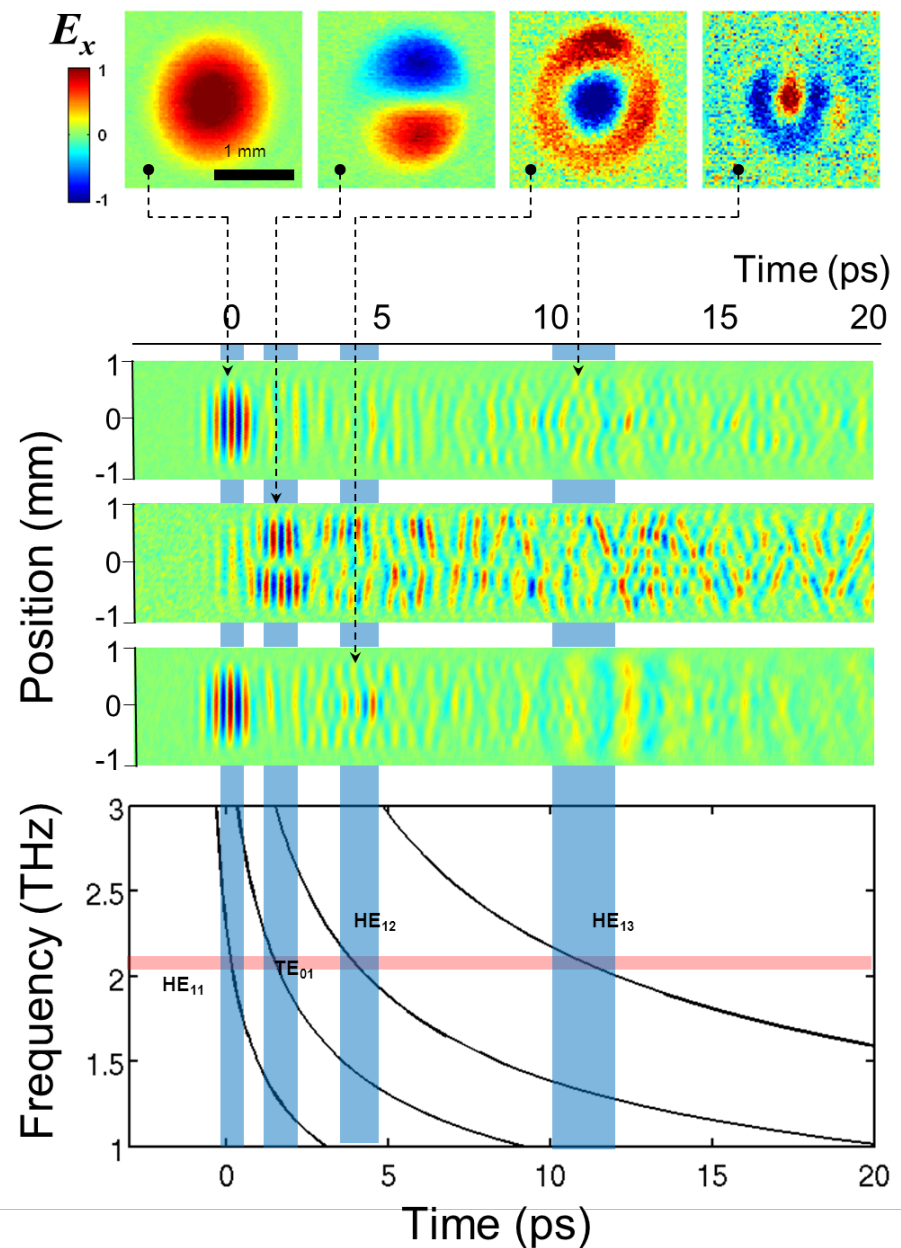
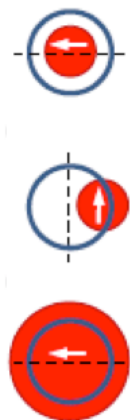
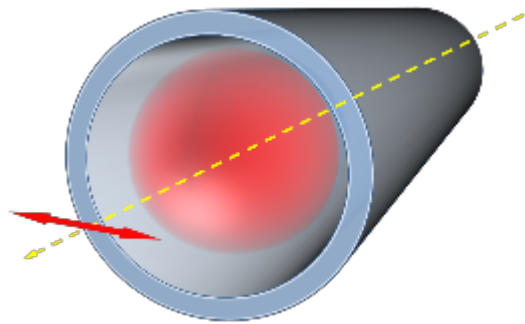
similar waveguide with a PS inner coating

## Transmission Loss at 2.5THz





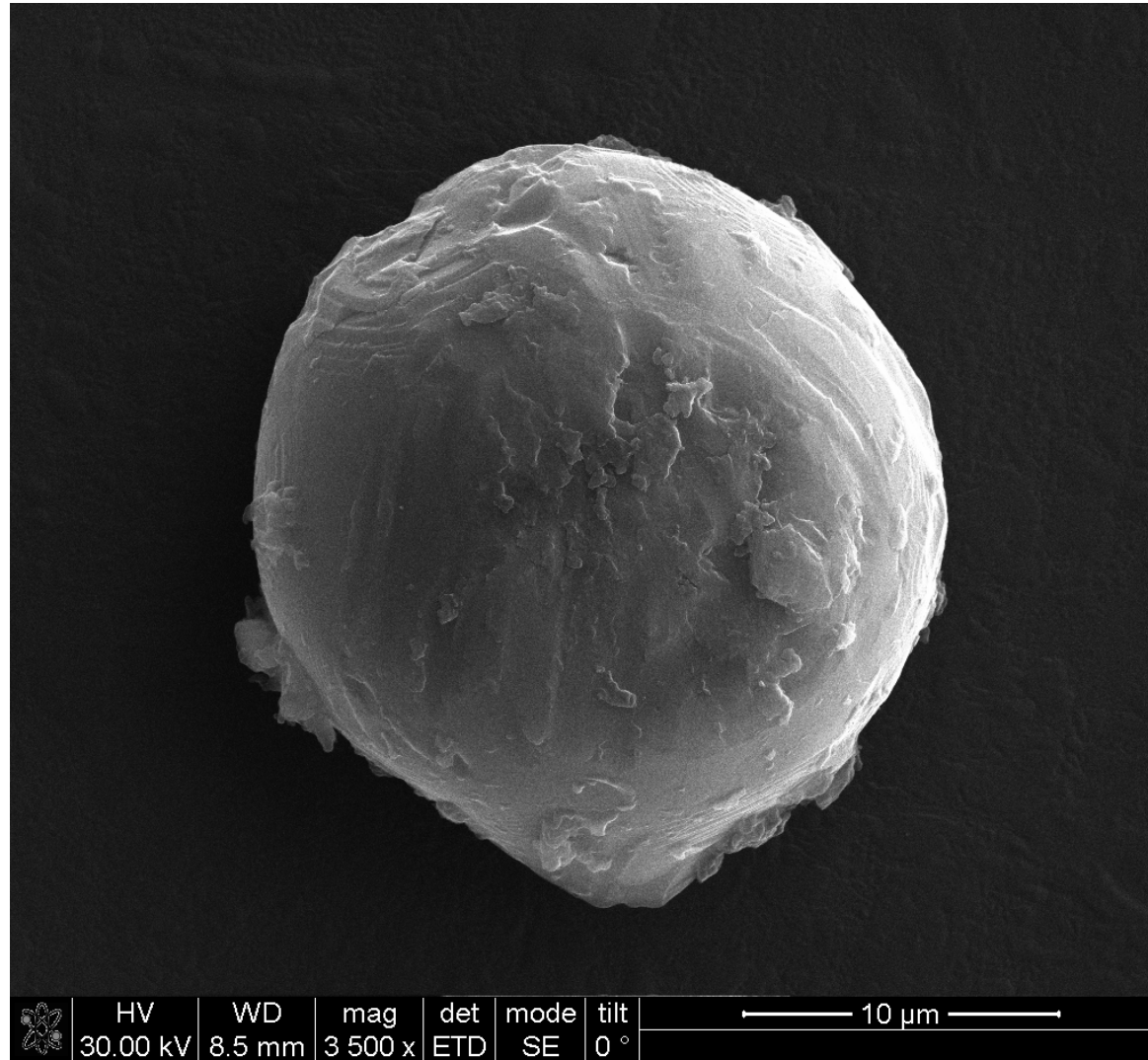




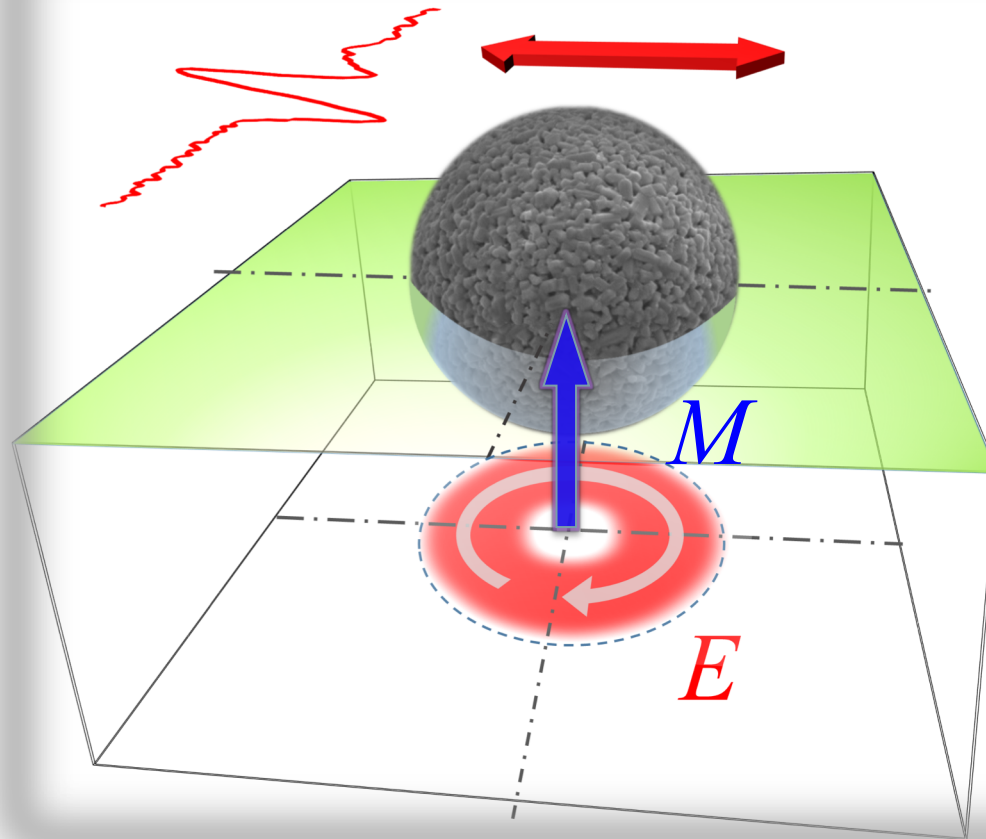


# *Sub-wavelength dielectric resonators*

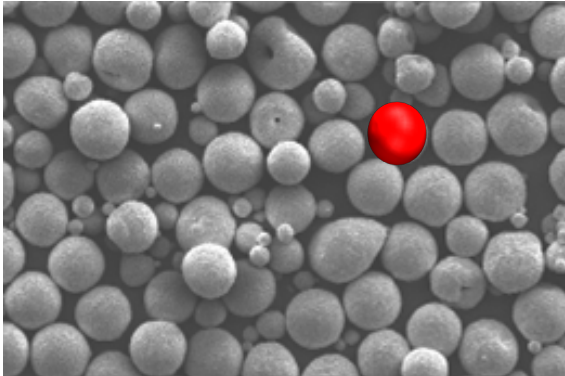
*Spectroscopy and mode imaging*



*$\text{TiO}_2$  microsphere:  
~20  $\mu\text{m}$  diameter*

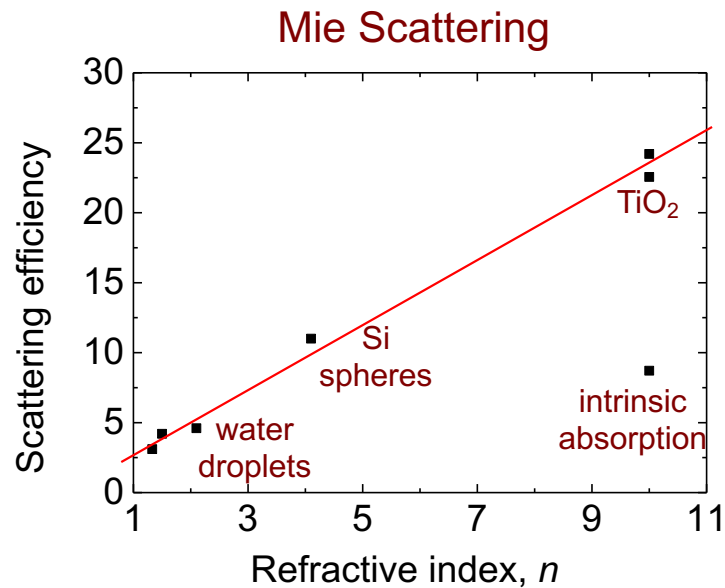
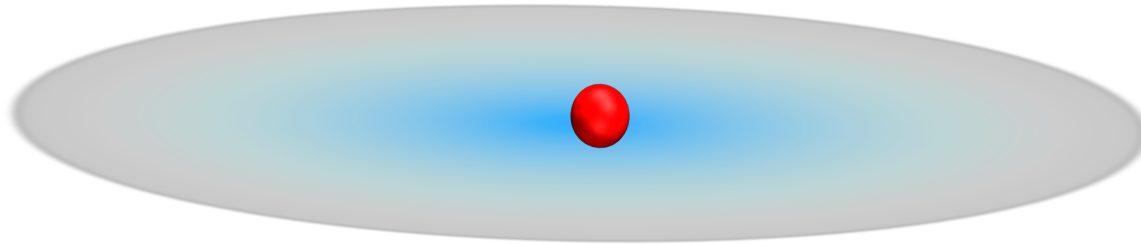


Sub-wavelength size  
of  $\text{TiO}_2$  resonators  $d \sim \lambda/n$



What is the effect of a single sphere?

Sub-wavelength size  
of  $\text{TiO}_2$  resonators  $d \sim \lambda/n$

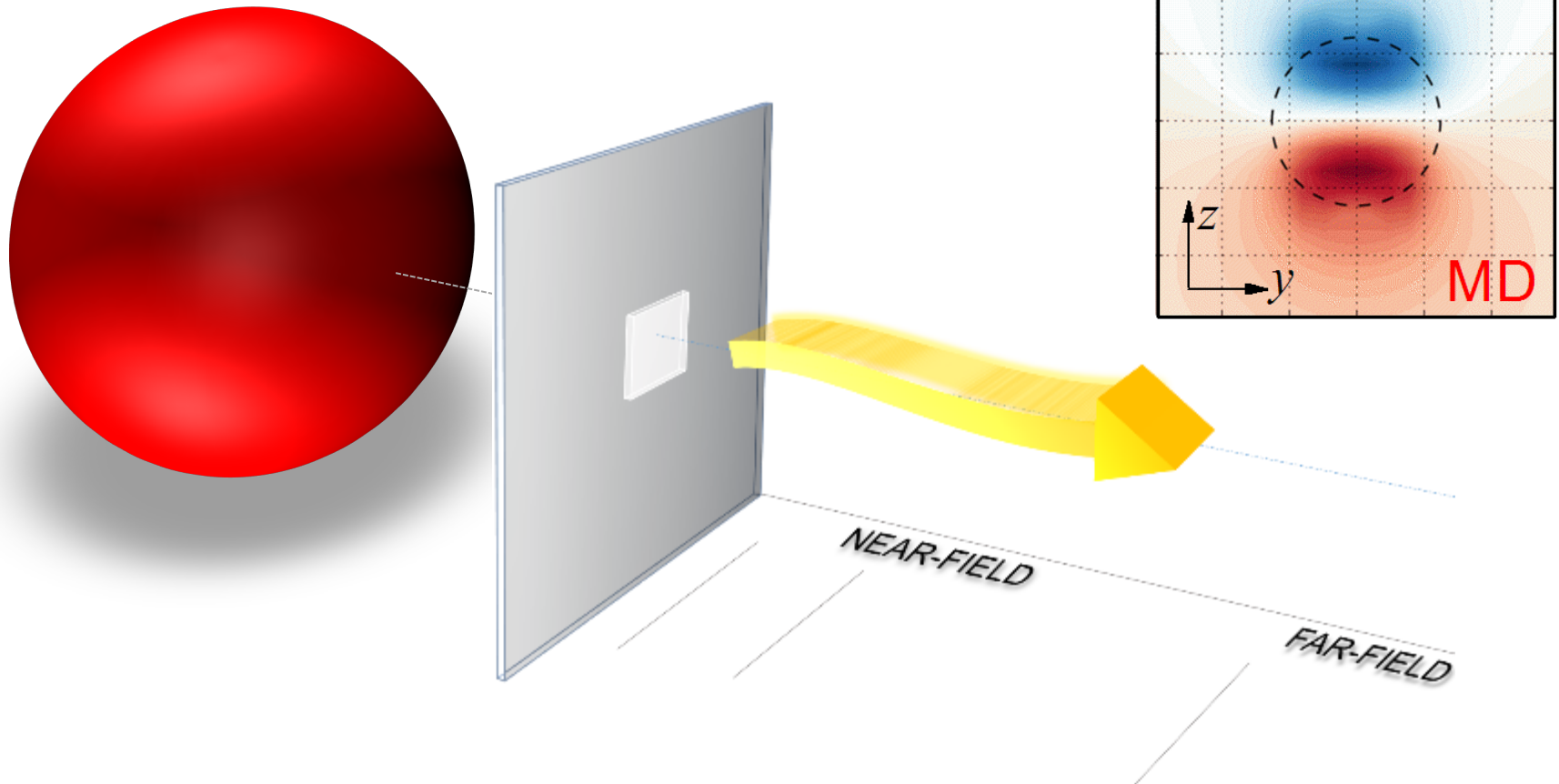


*Total scattered power reduces due to the physical cross-section scaling with  $n^{-2}$*

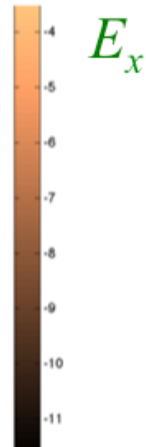
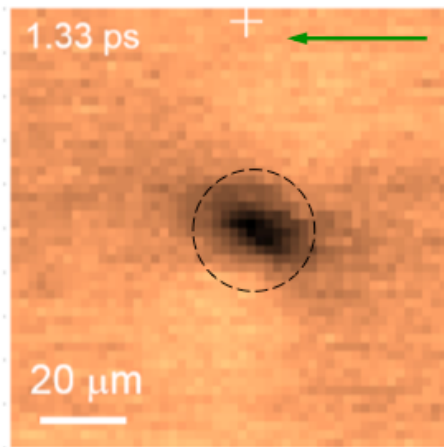
Far-field total extinction by a single  $\text{TiO}_2$  sphere  
(est. for typical THz-TDS) :

0.1-1.0%

*High EM field confinement by a dielectric object*



*Enhanced transmission through aperture can be used to probe high- $\epsilon$  resonators*



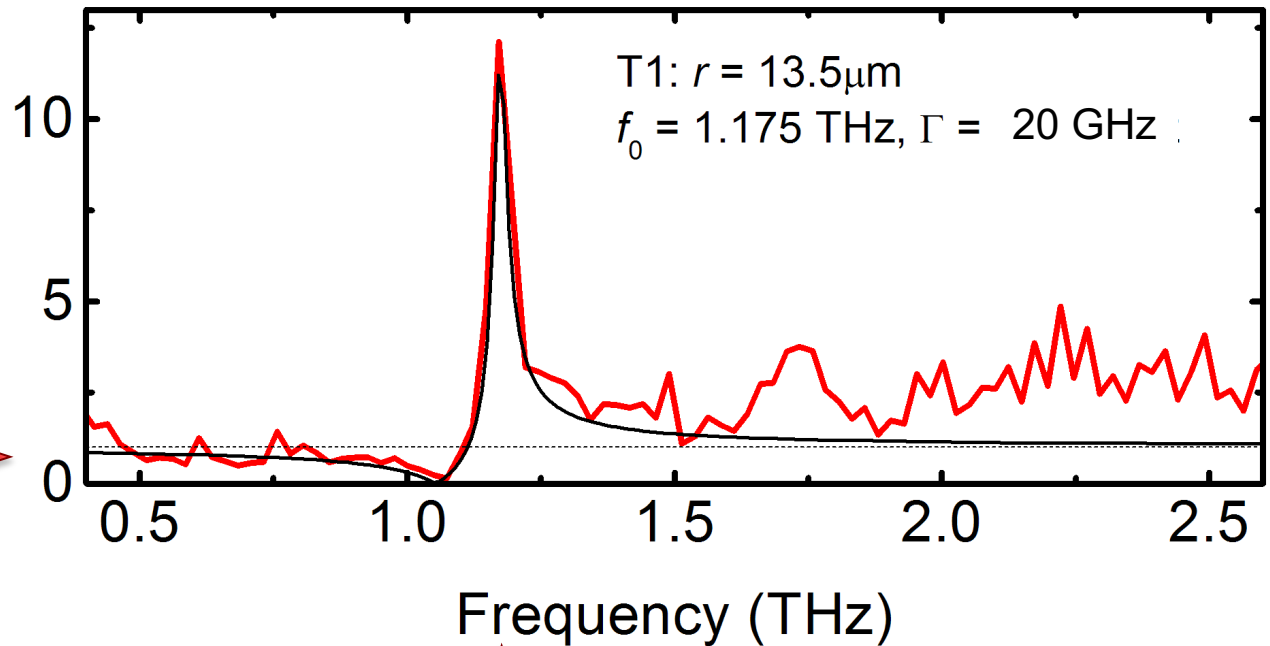
*Enhancement factor and width to quantify the resonator*



*No effect on transmission away from resonance*

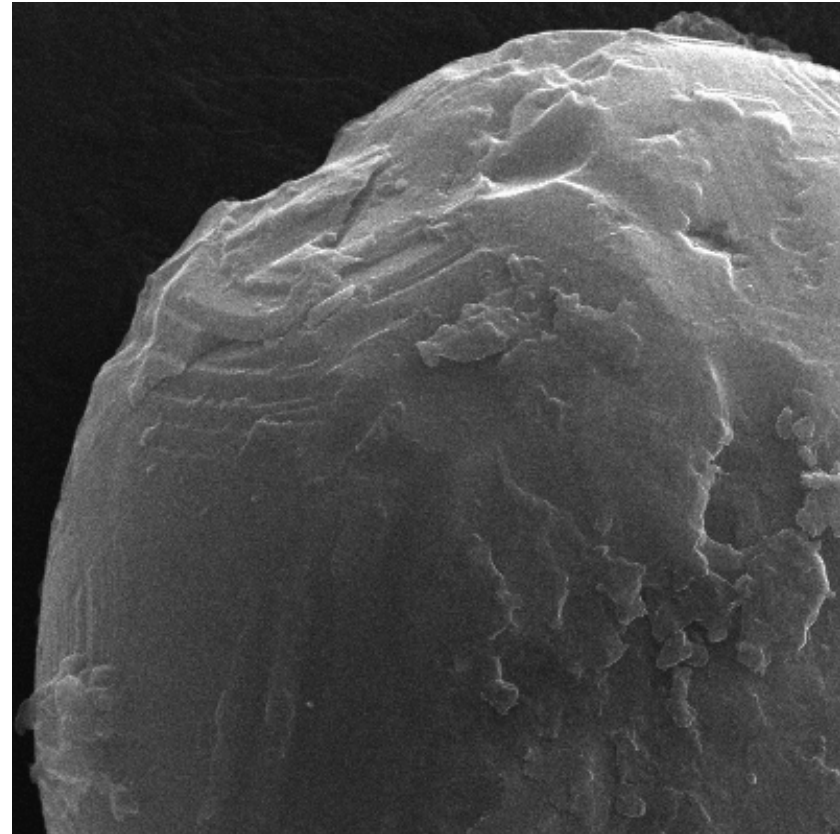
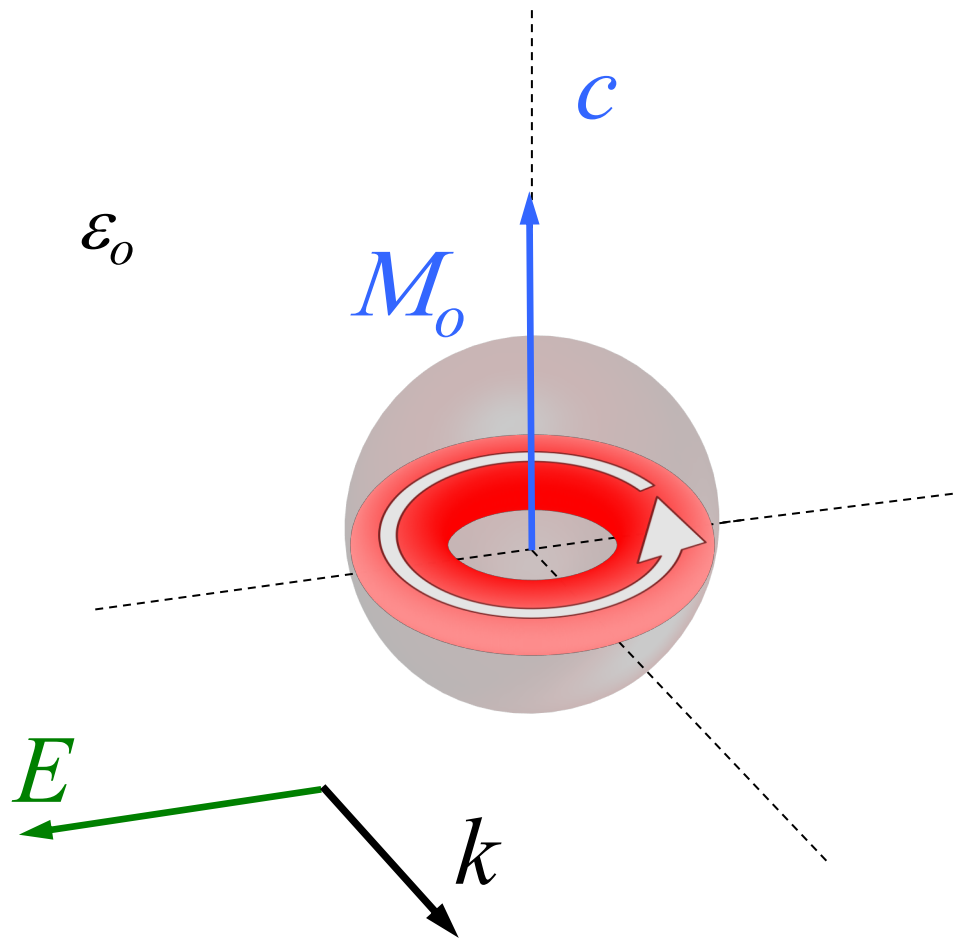


$E_{\text{ms}}/E_0$



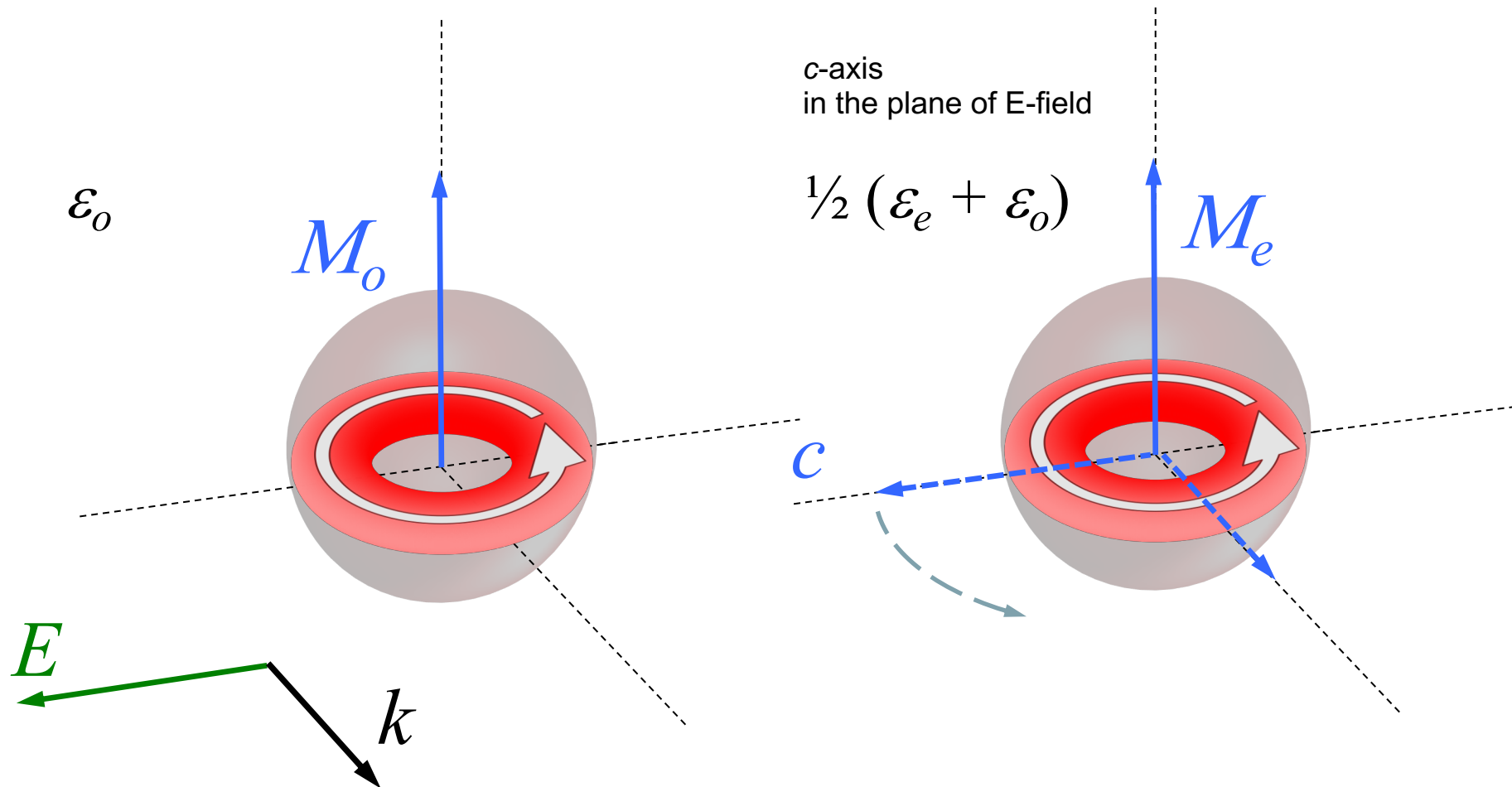
*Interference minimum (Fano line-shape)*

$\text{TiO}_2$ :  $\epsilon_e = \sim 150$ ;  $\epsilon_o = \sim 70$

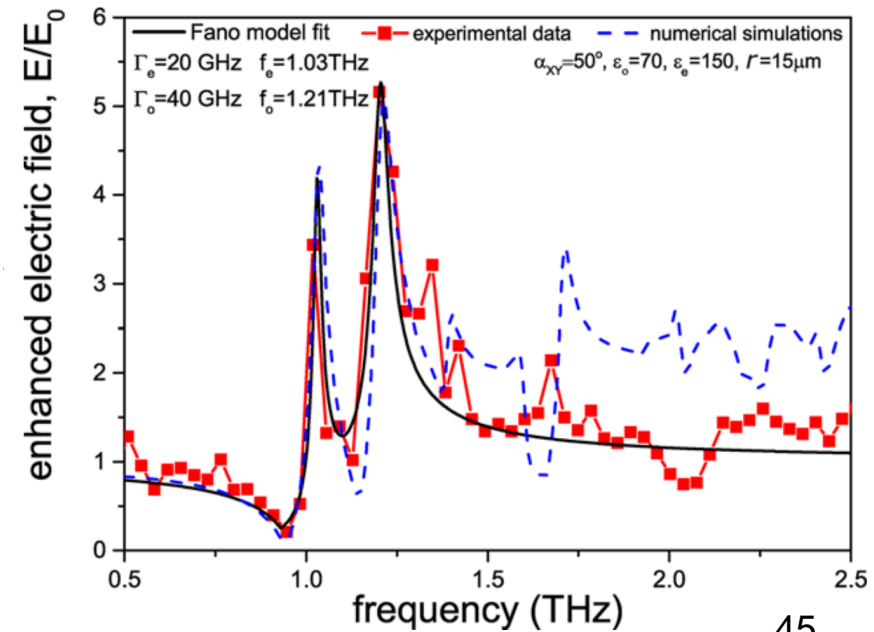
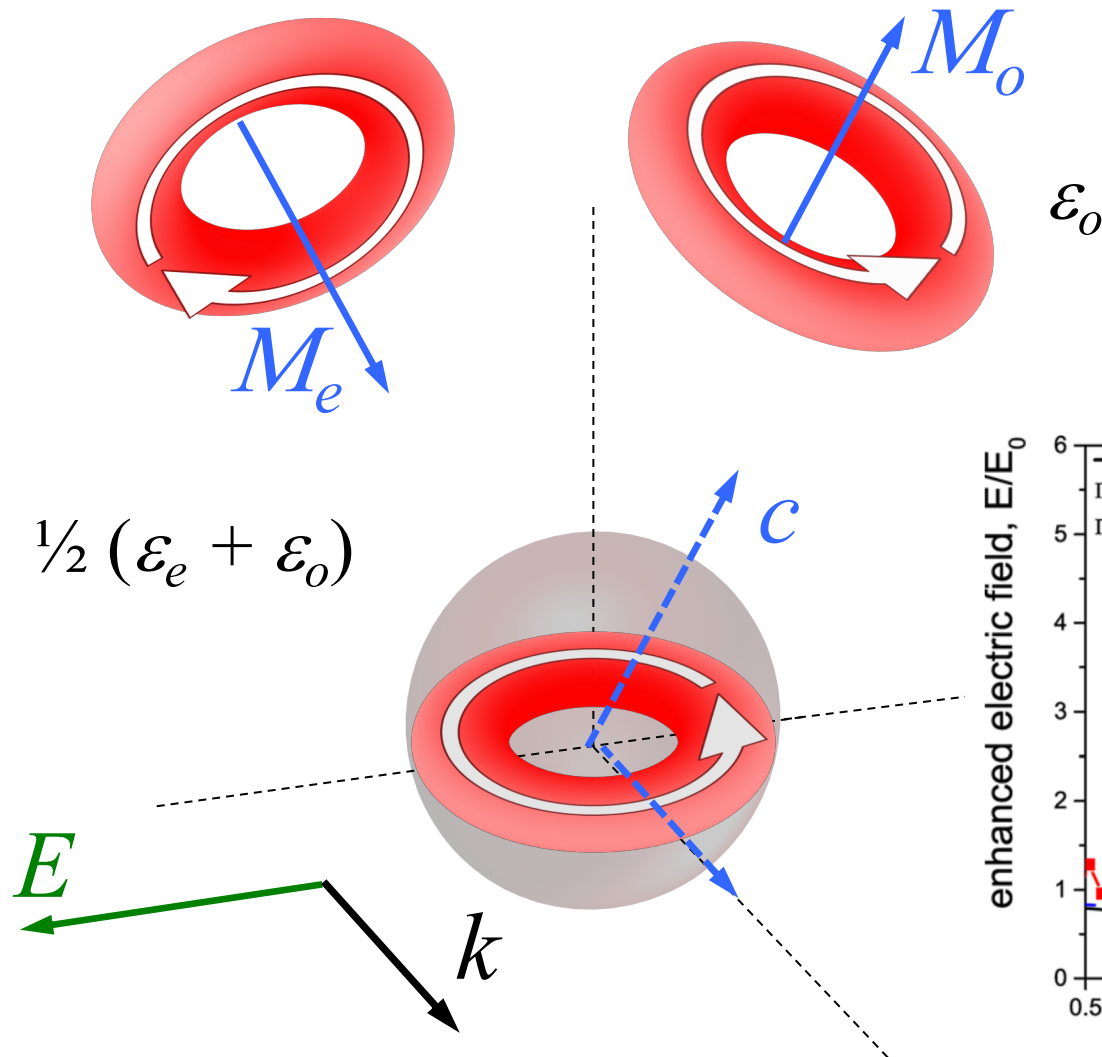


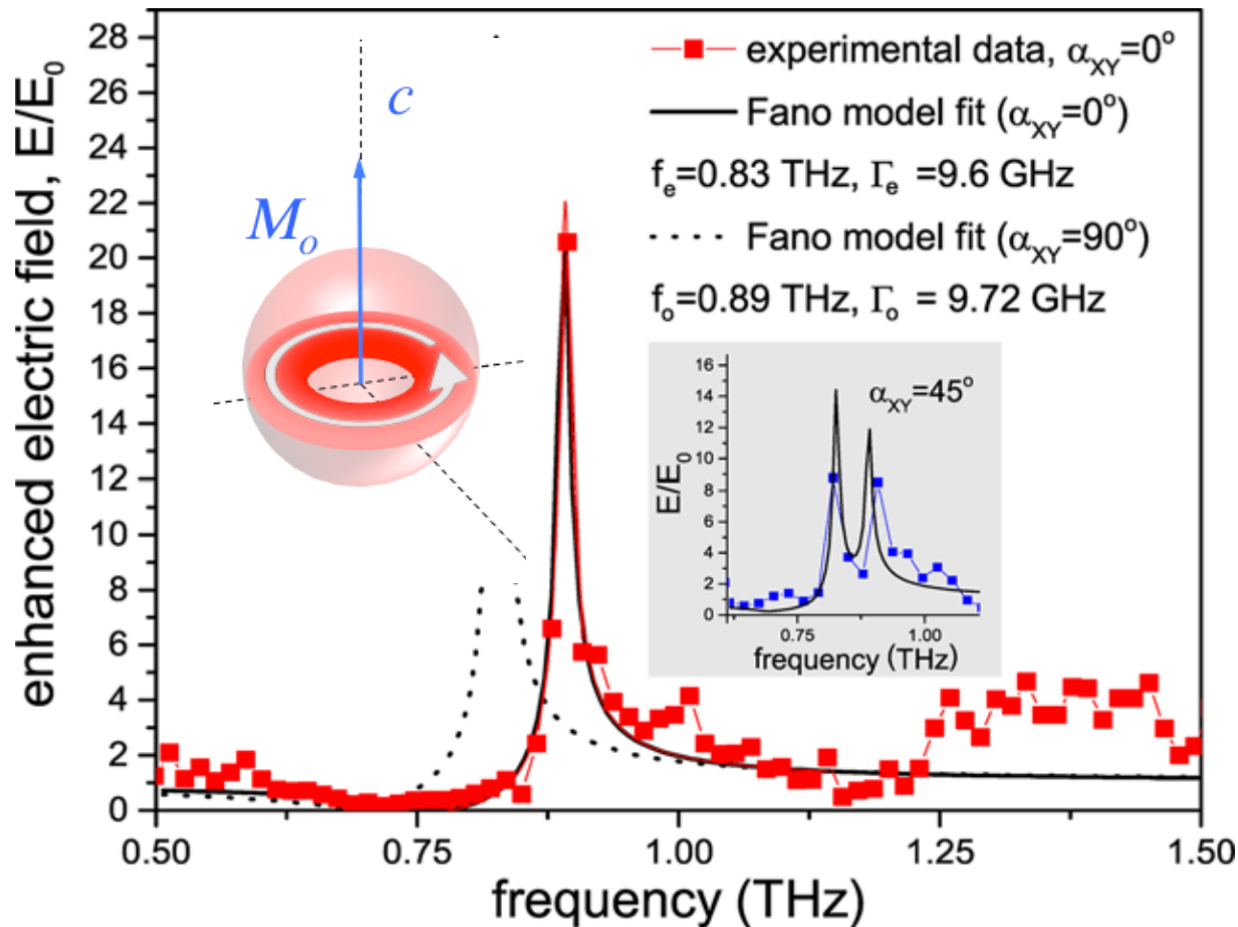


TiO<sub>2</sub>:  $\epsilon_e = \sim 150$ ;  $\epsilon_o = \sim 70$



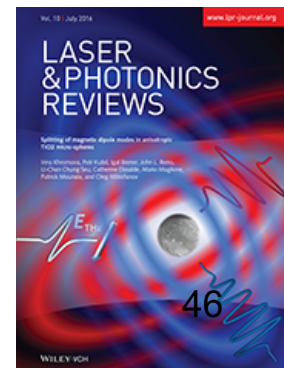
TiO<sub>2</sub>:  $\epsilon_e = \sim 150$ ;  $\epsilon_o = \sim 70$

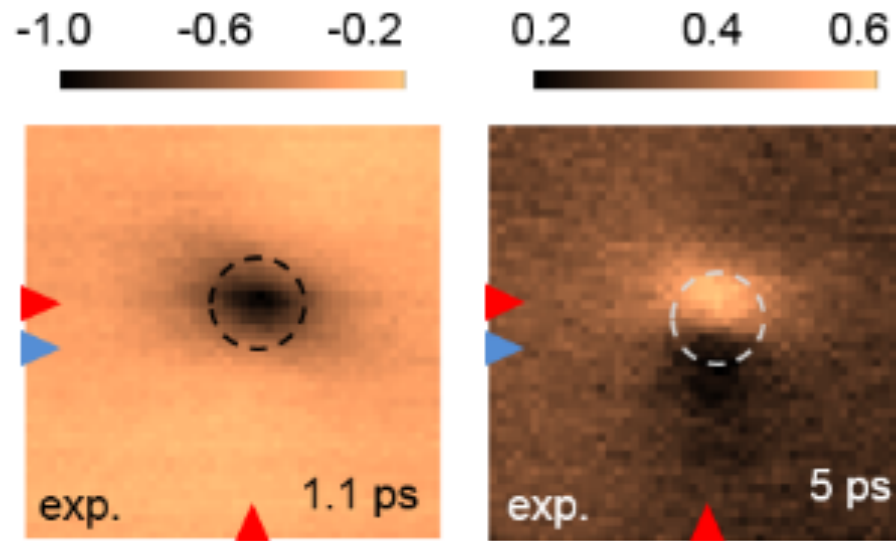


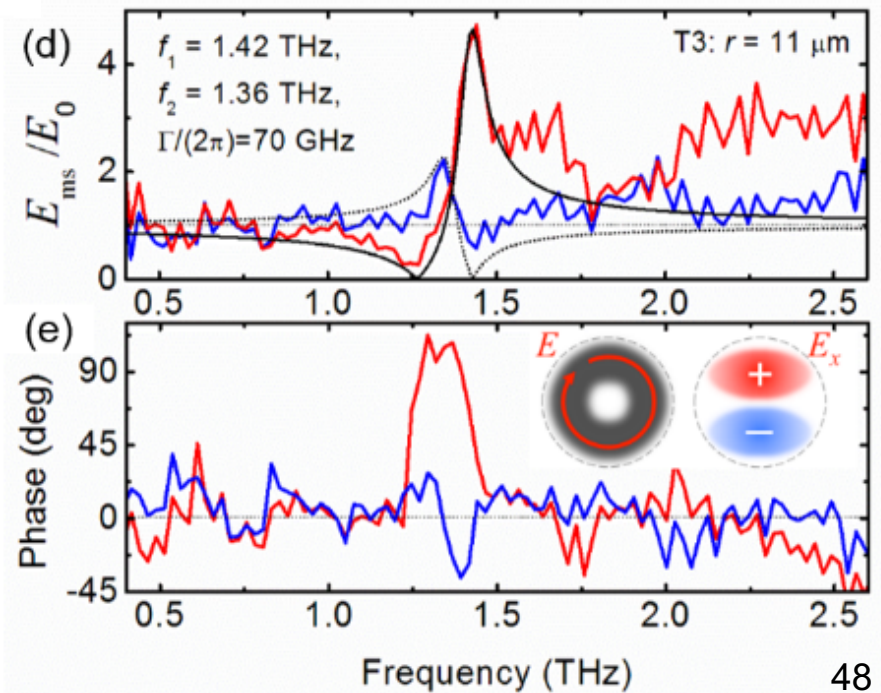
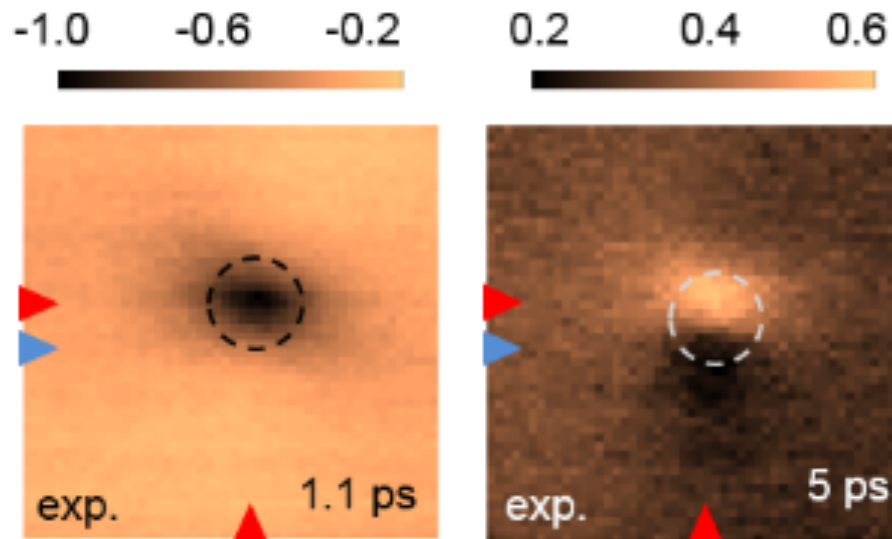


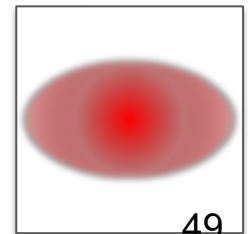
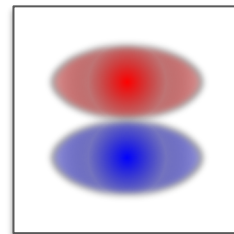
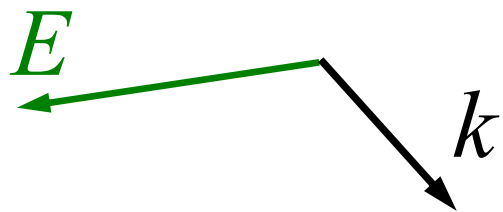
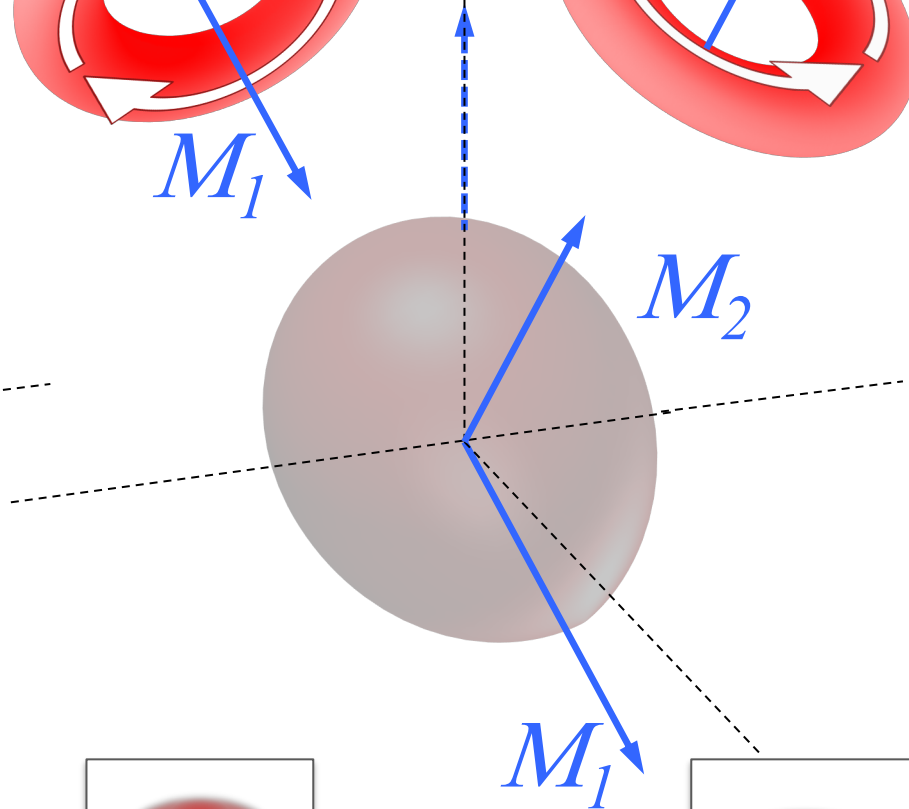
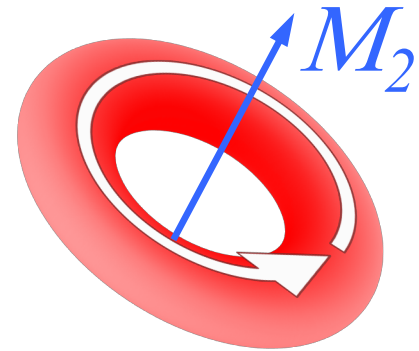
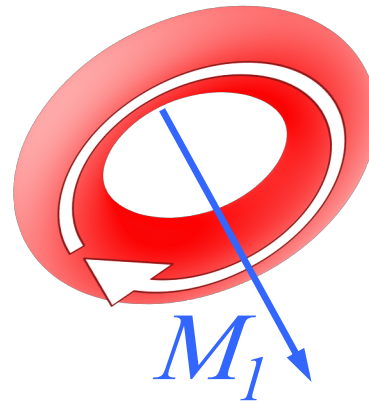
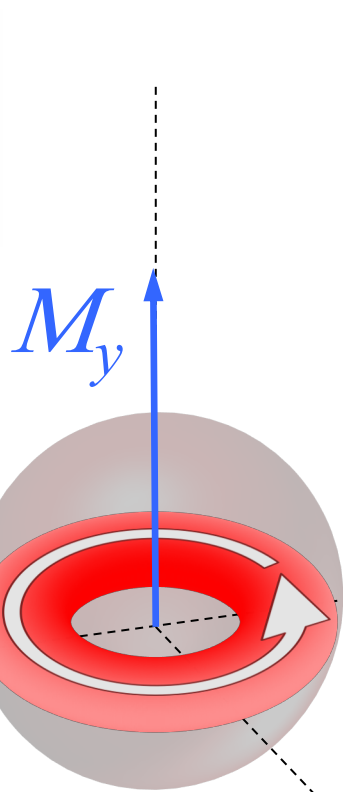
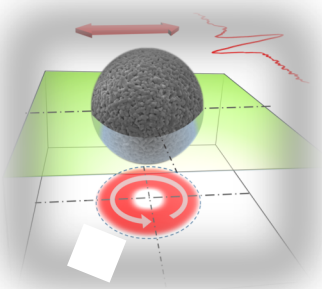
*Near-field spectroscopy allows ‘seeing’ both modes*

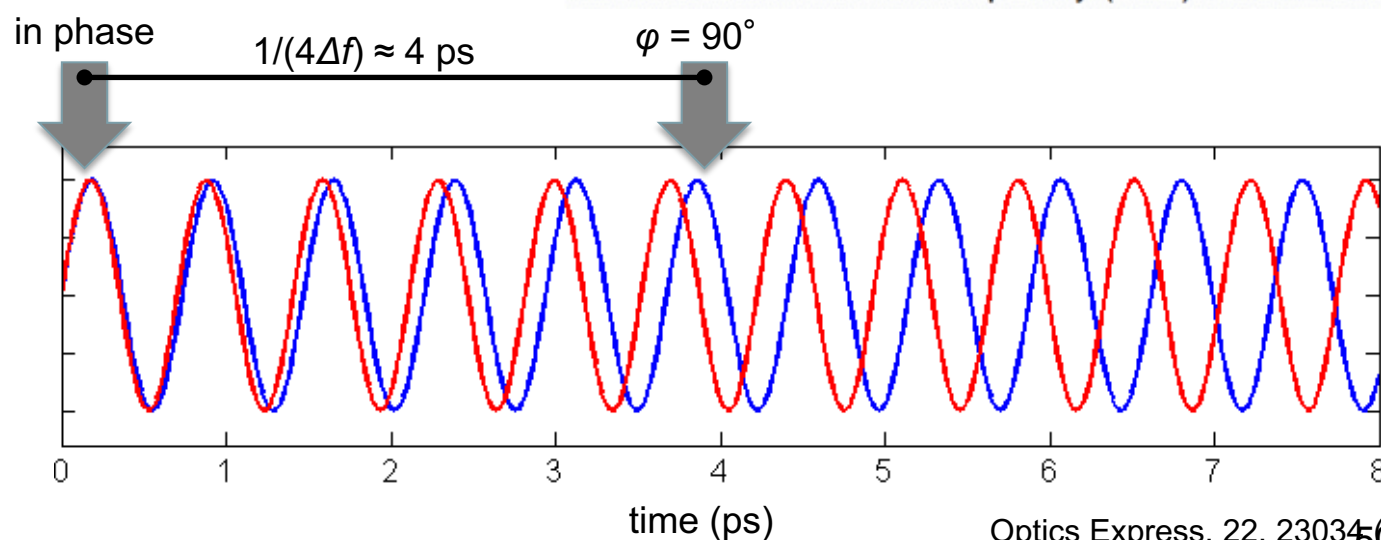
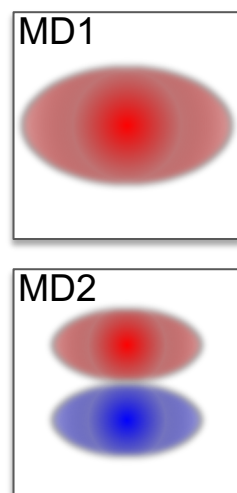
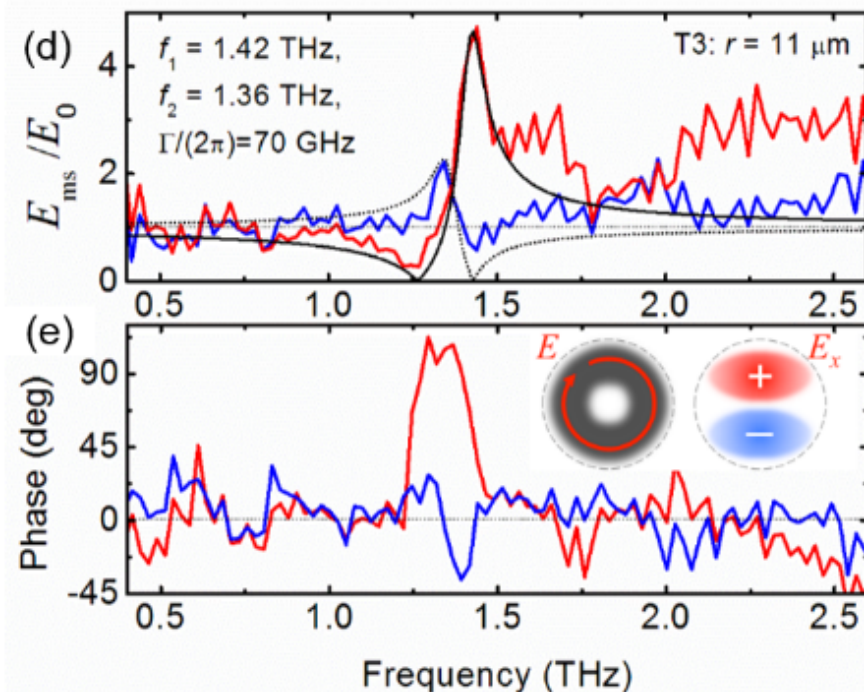
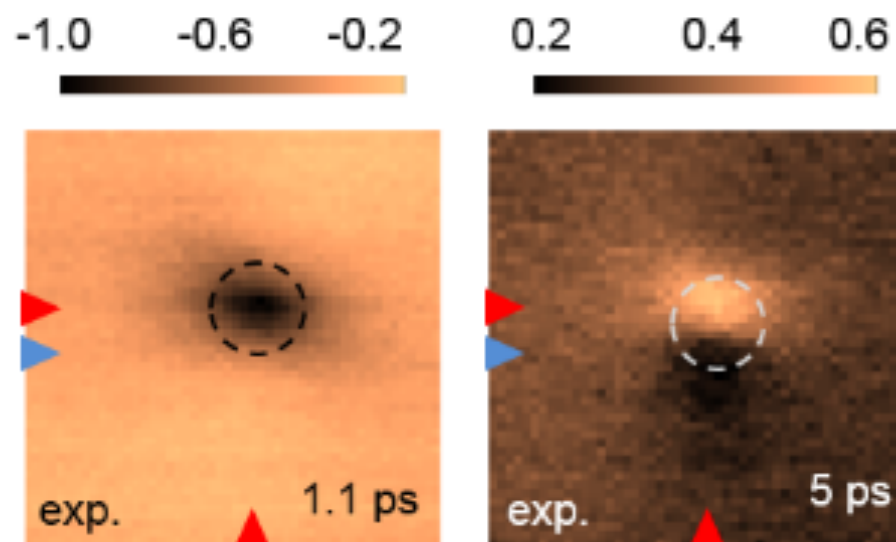
I. Khromova et al., *Laser and Photon. Reviews* (2016)







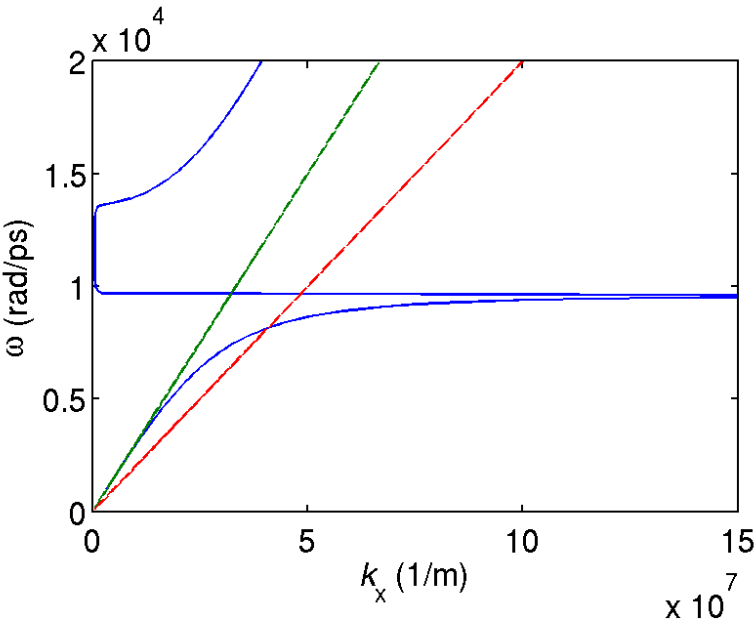




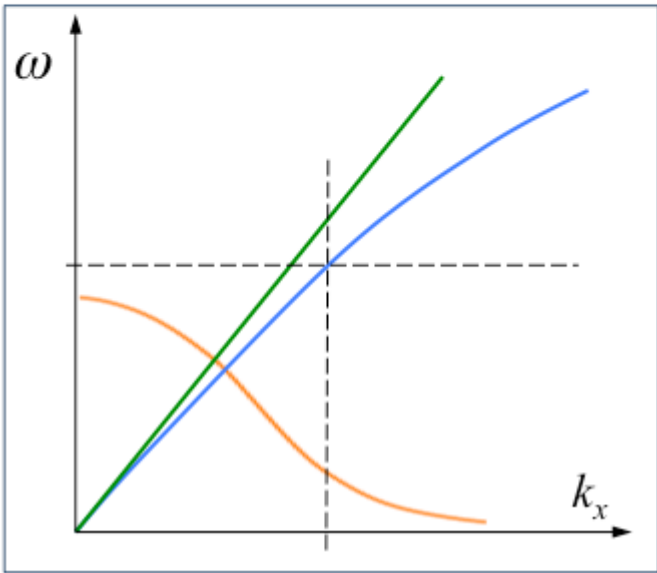
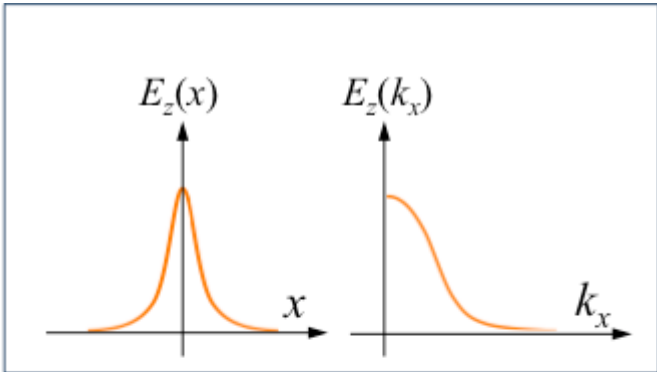
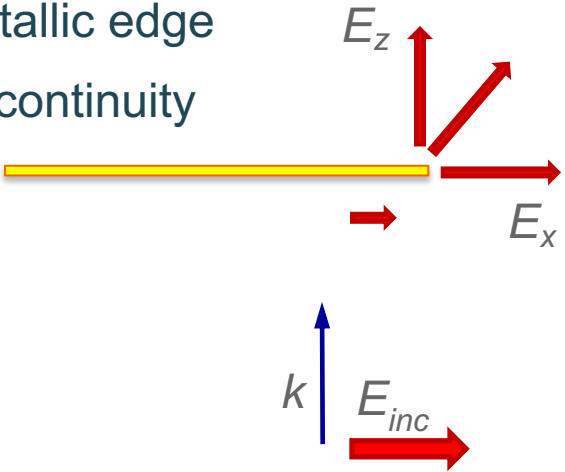
# *Plasmonic resonators and surface waves*

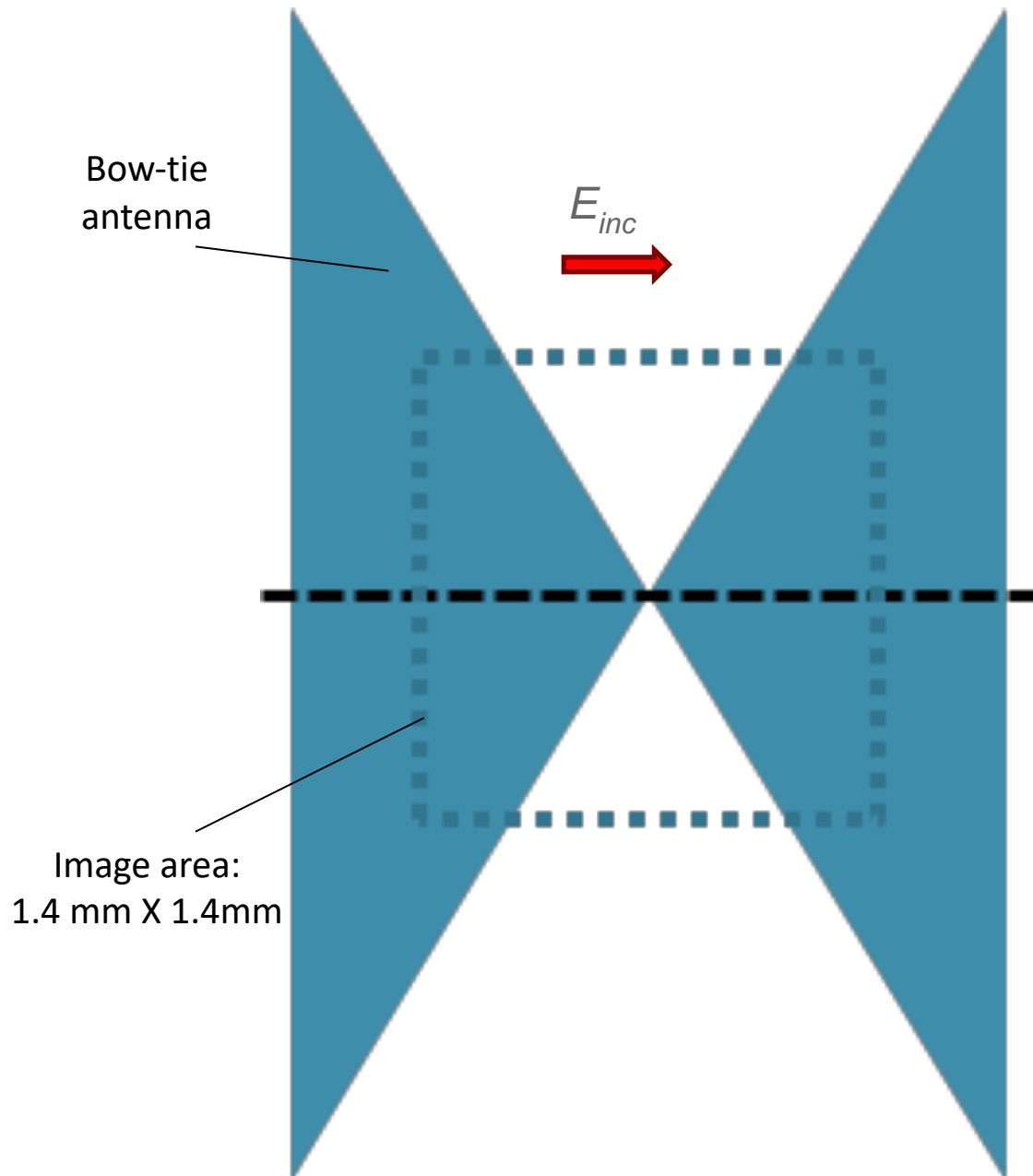
*Spectroscopy and mode imaging*

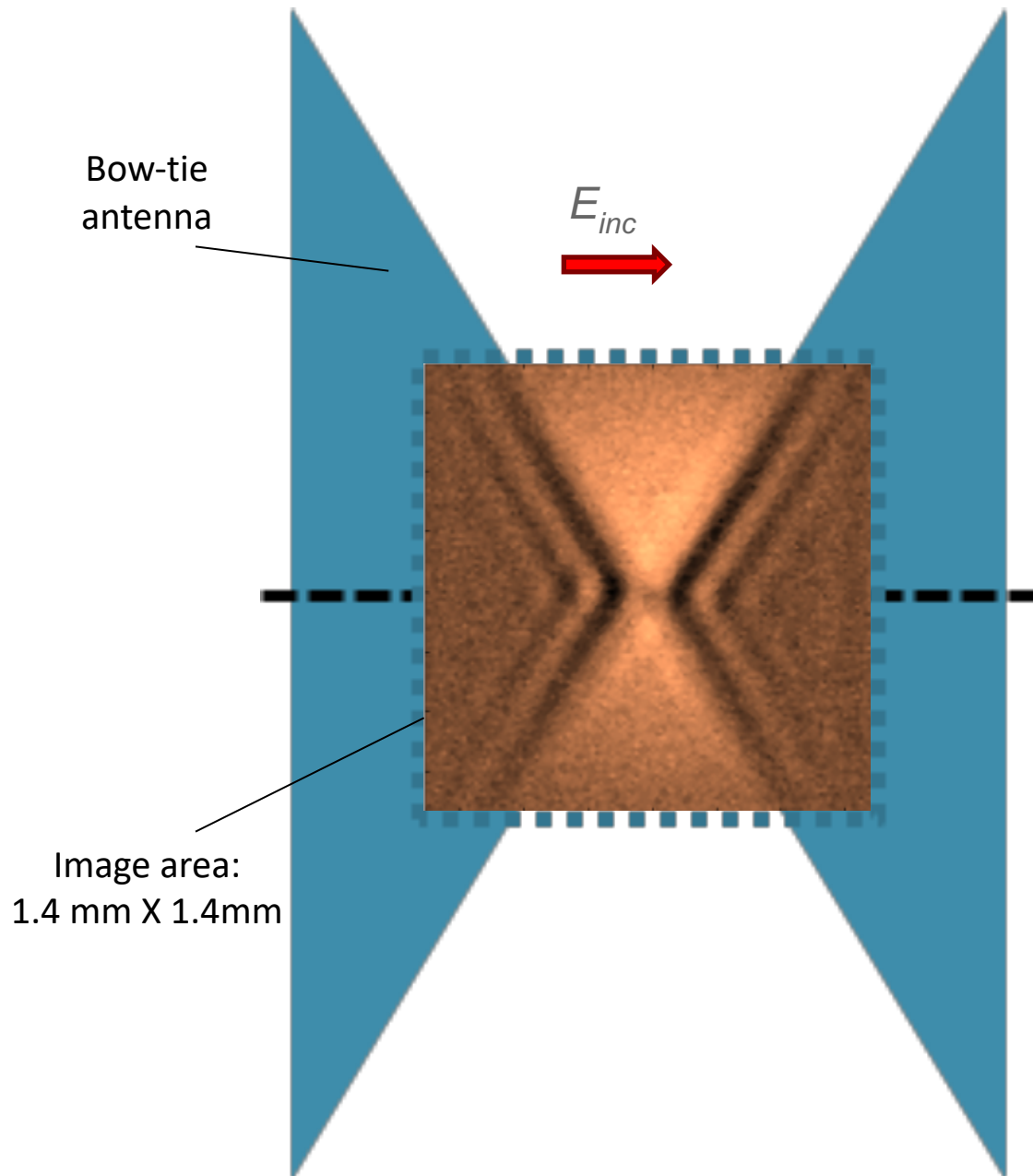


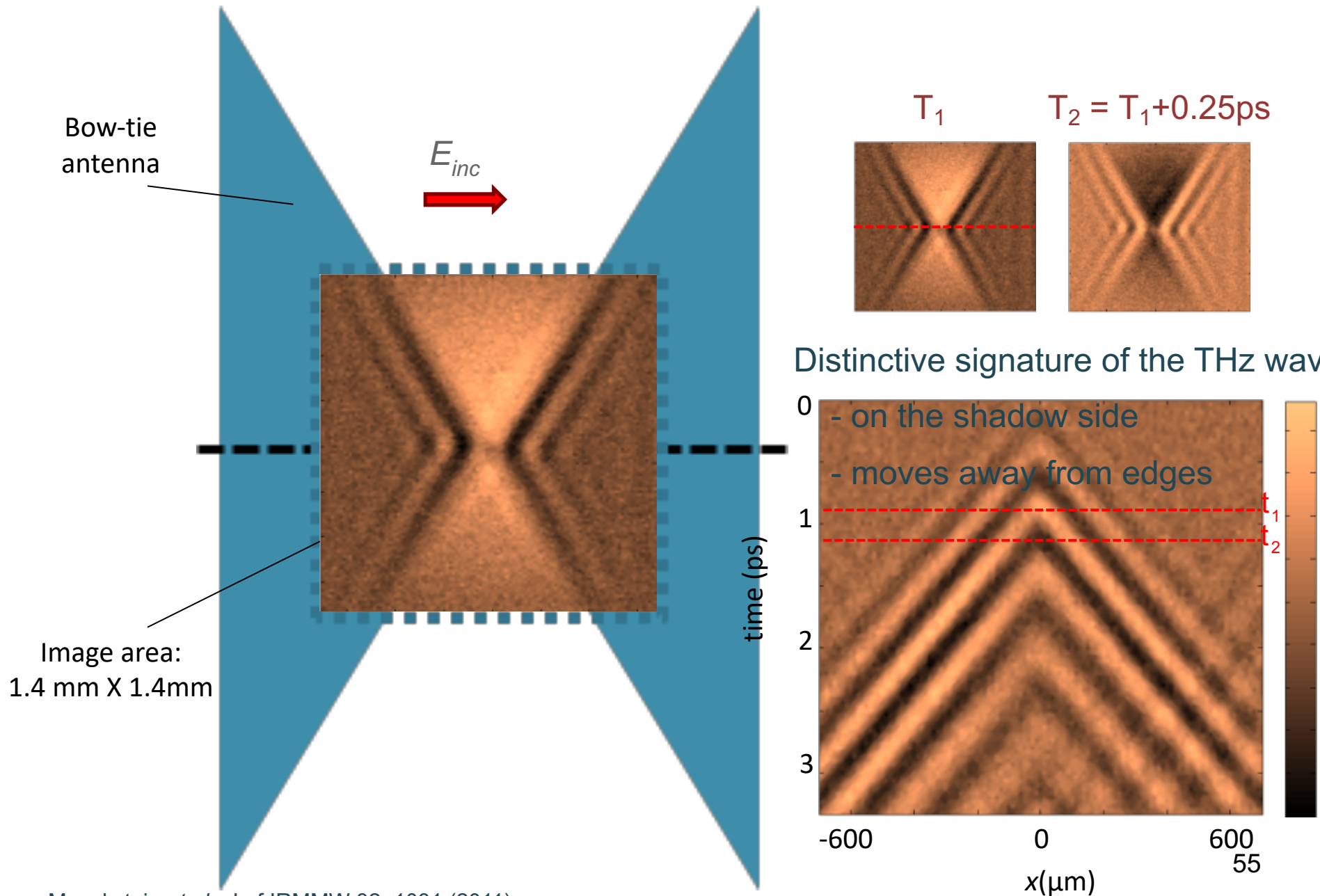


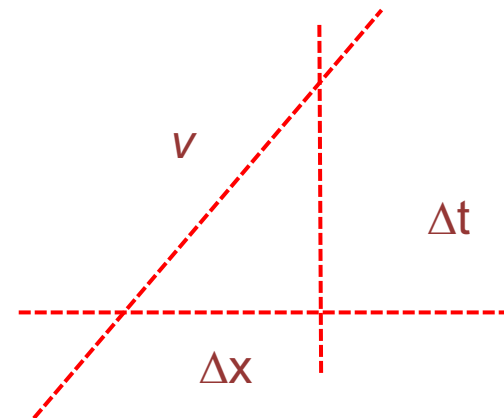
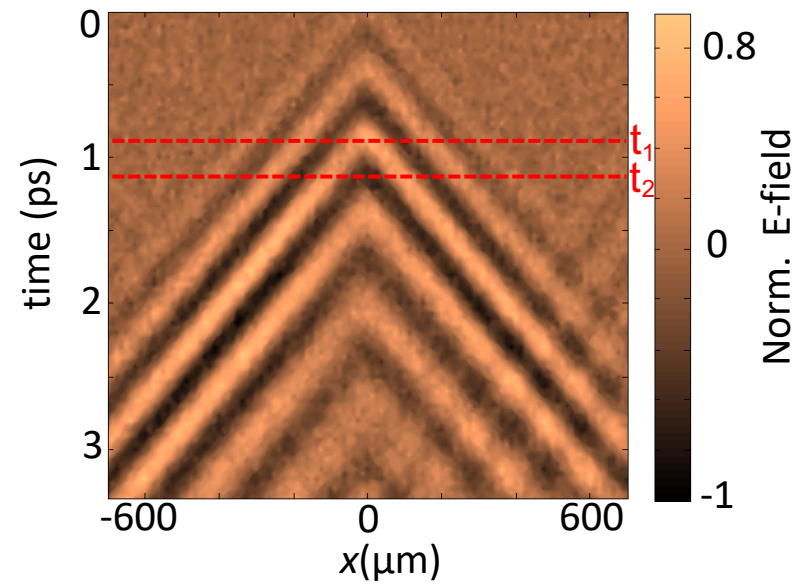
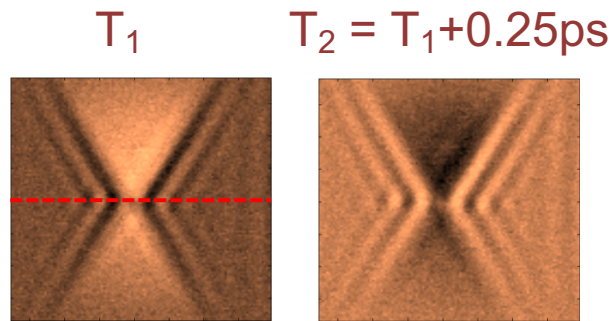
Metallic edge  
discontinuity

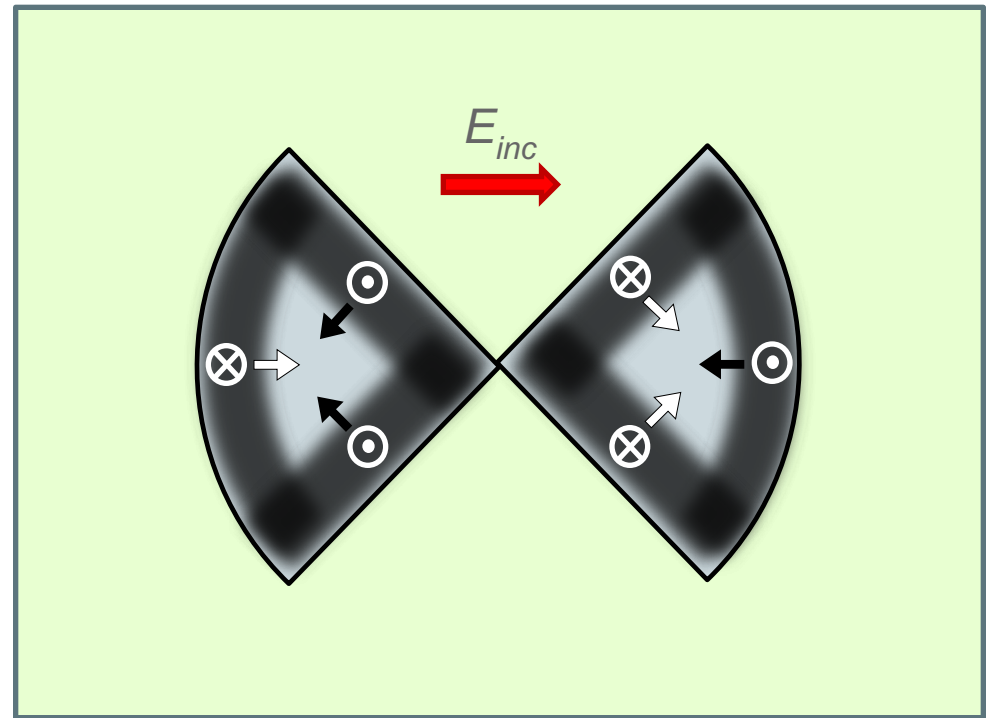
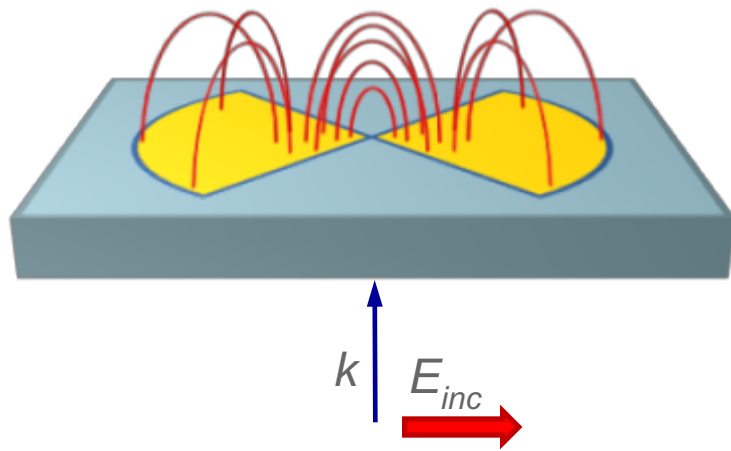


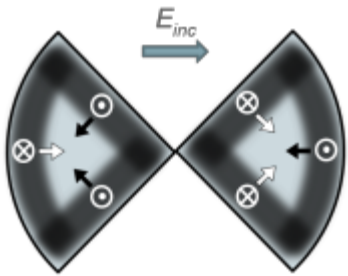




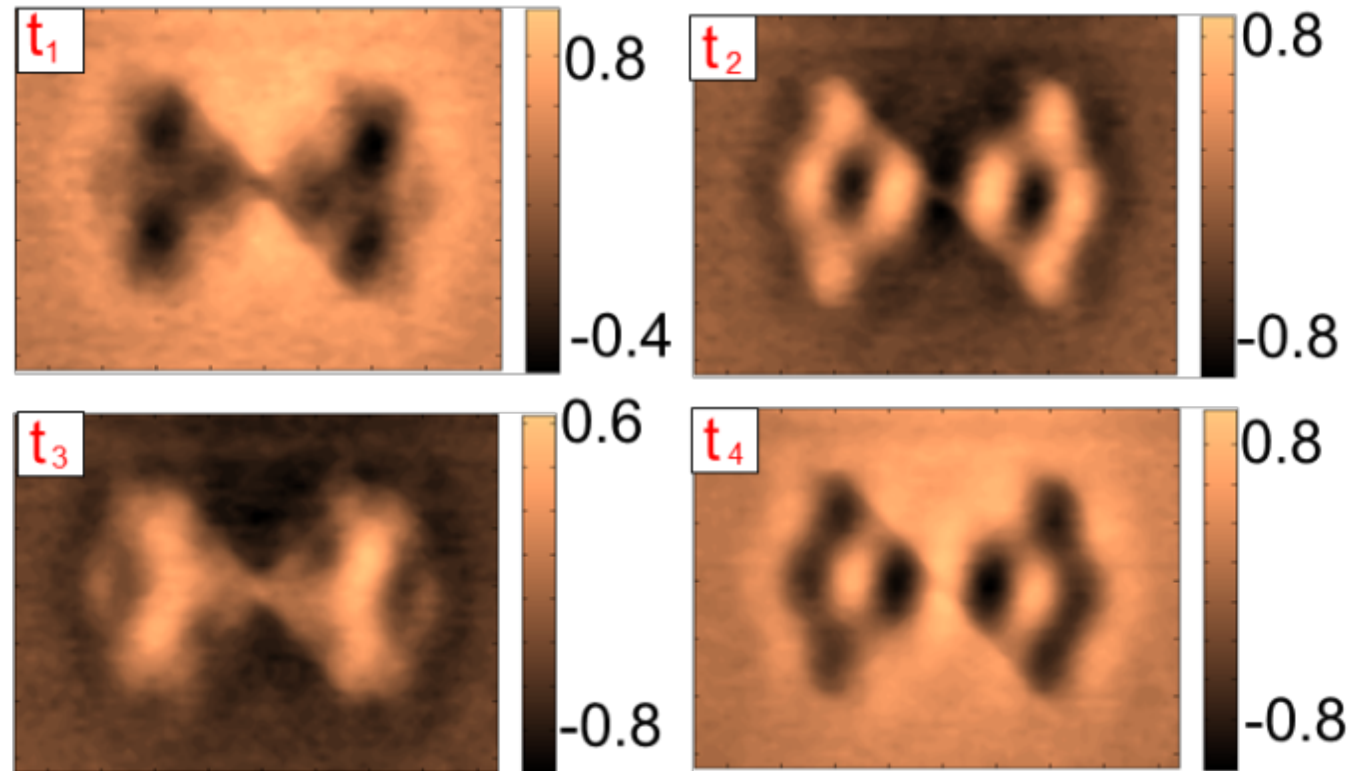




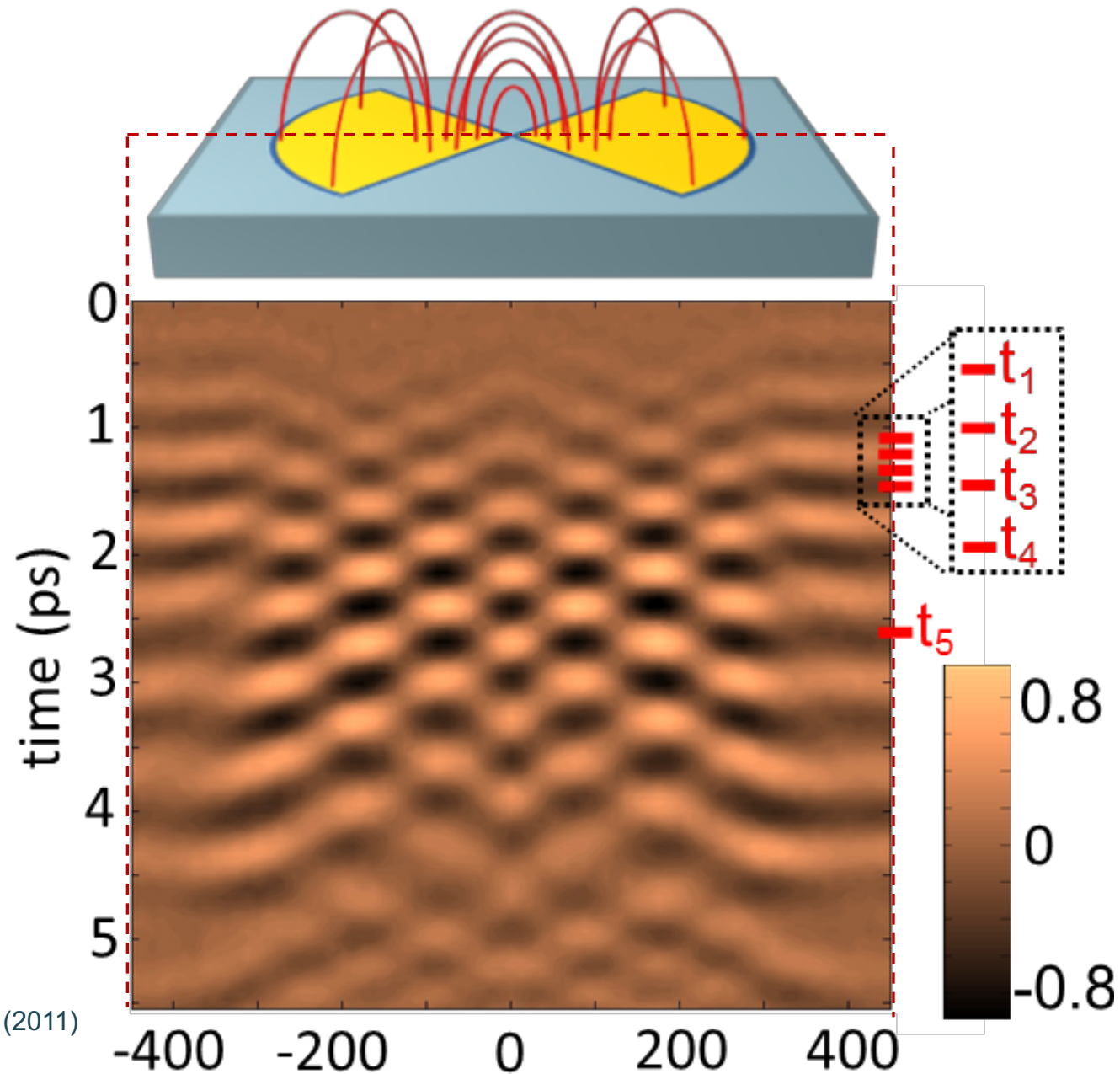




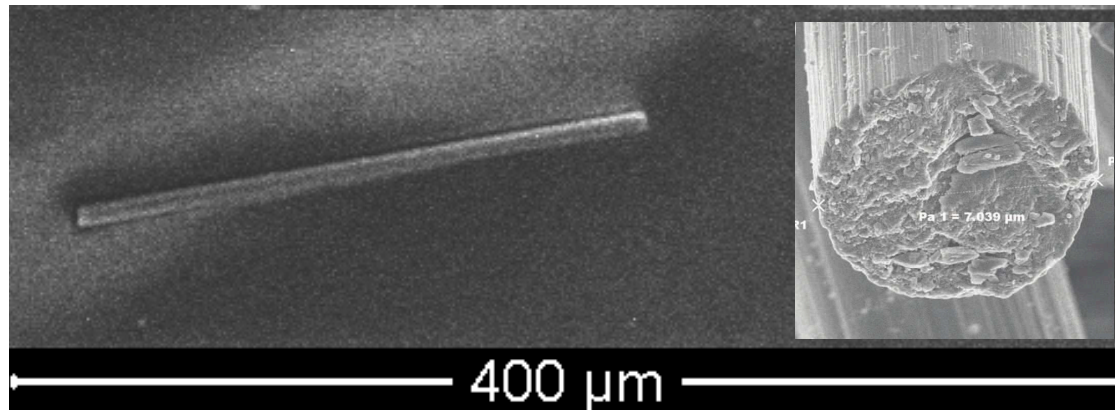
Consecutive images (“frames”) ~0.13 picosecond apart



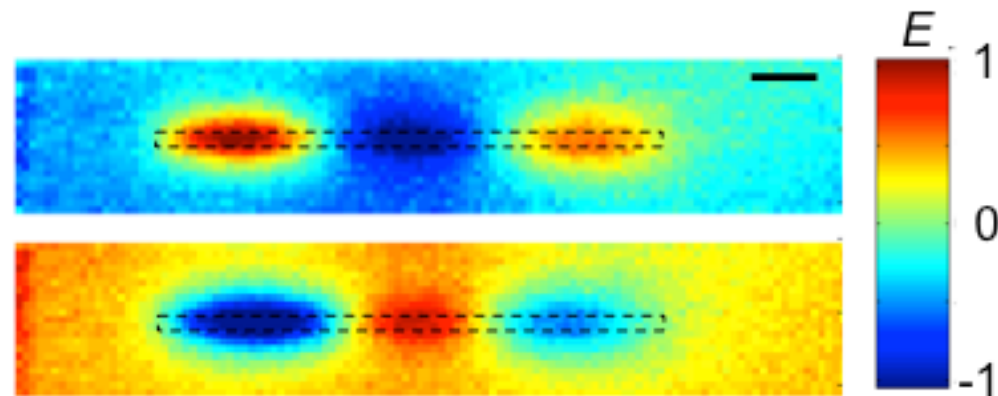
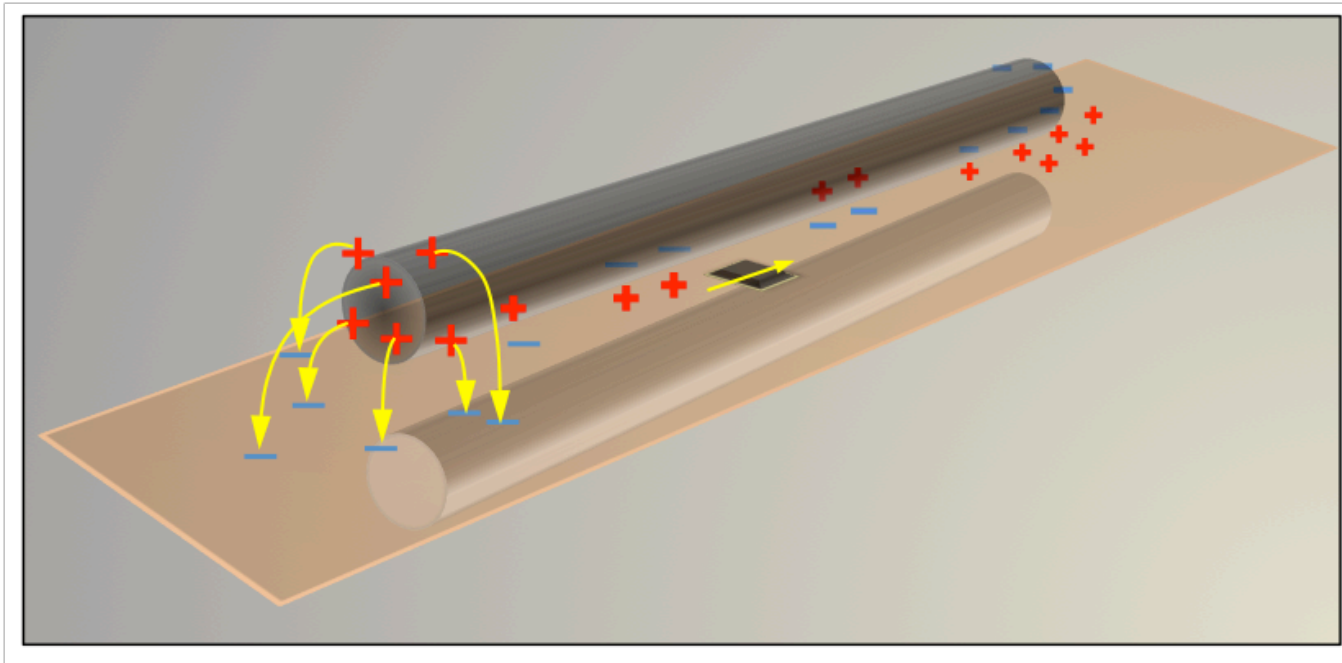
Mueckstein *et al.*, J.of IRMMW 32, 1031 (2011)

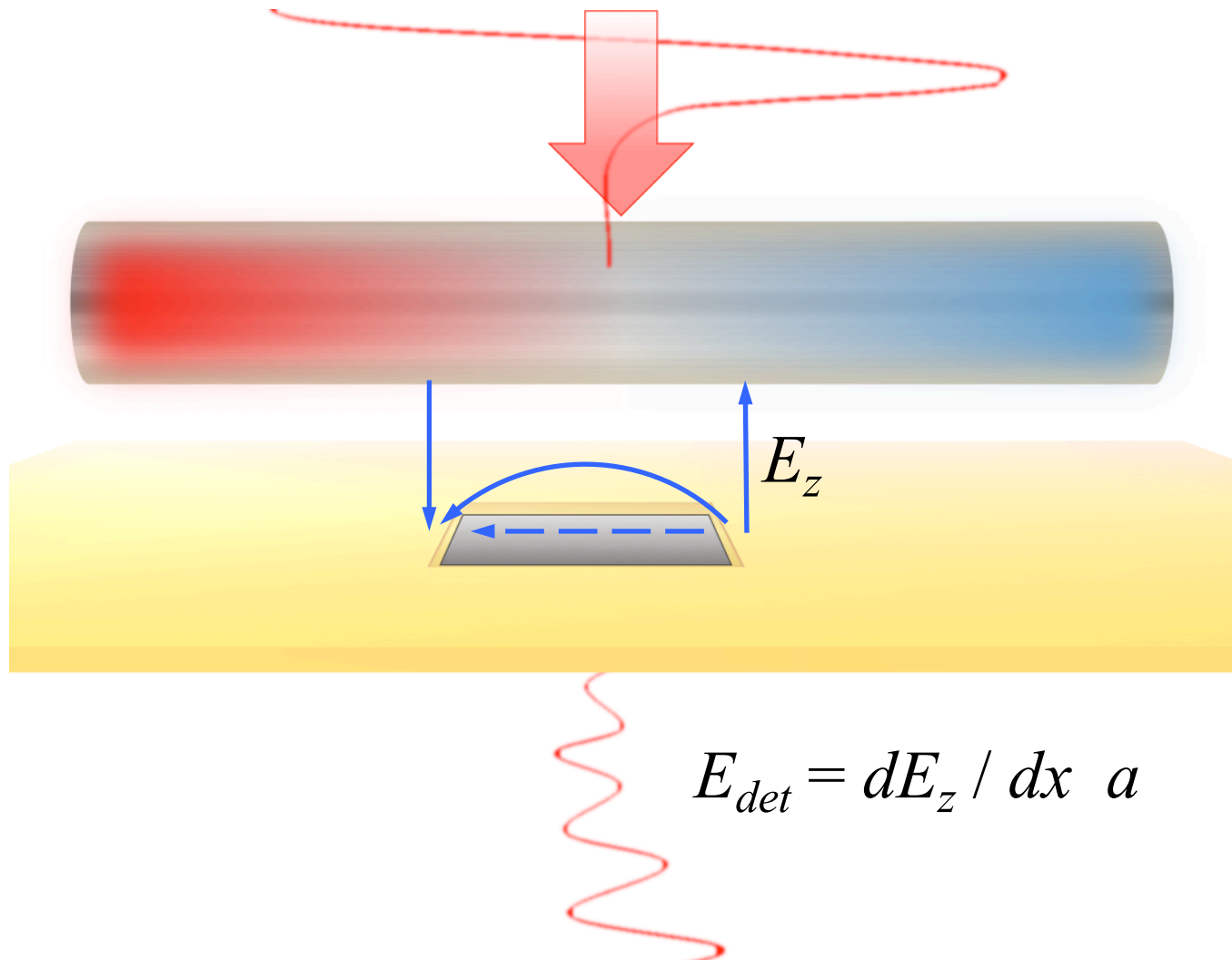




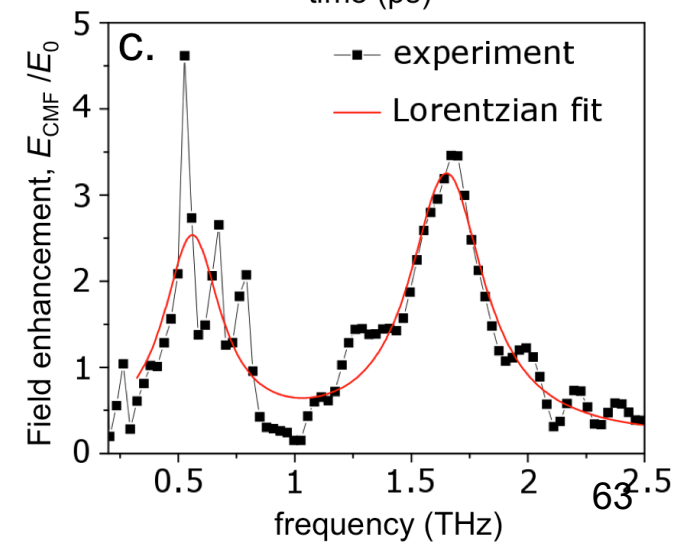
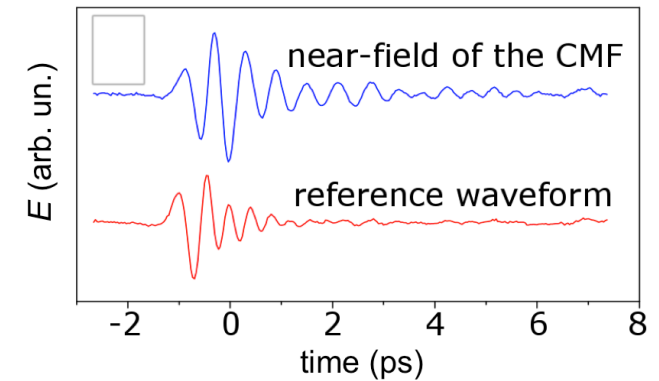
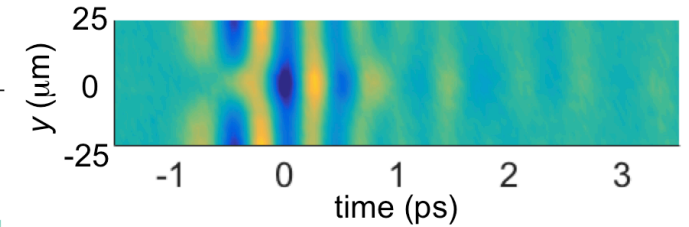
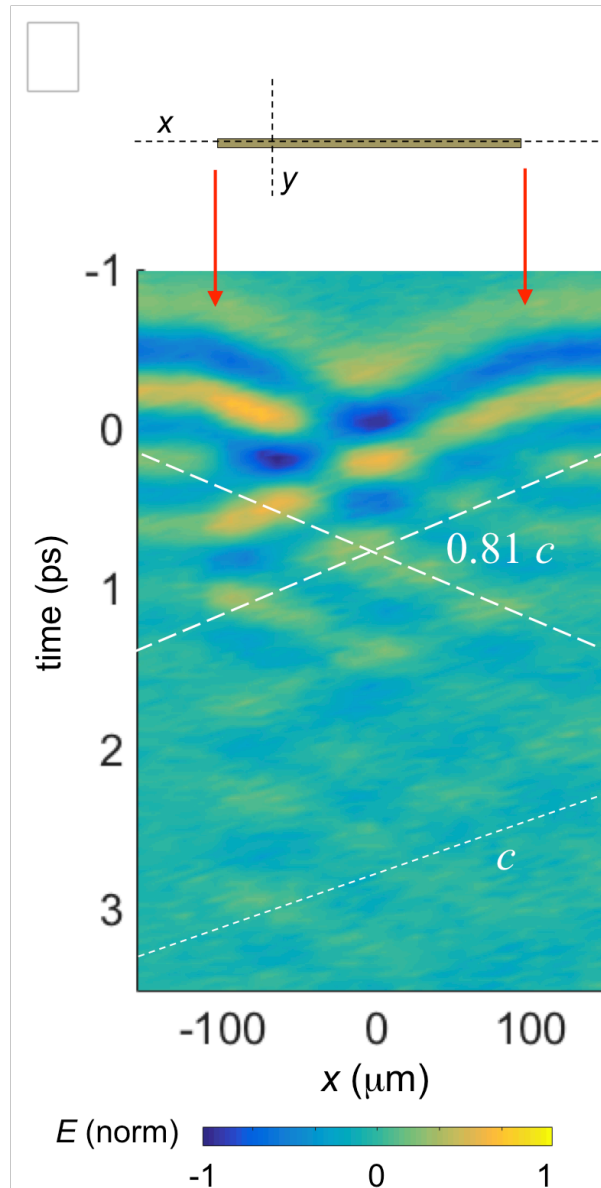
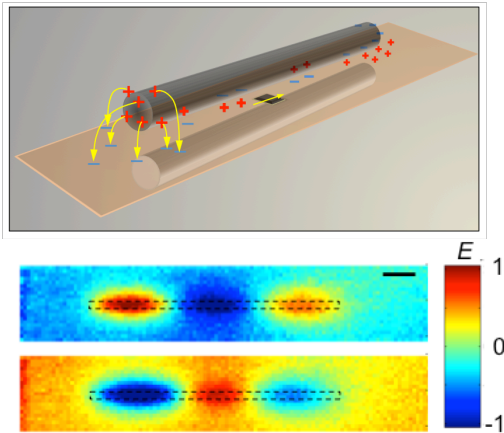


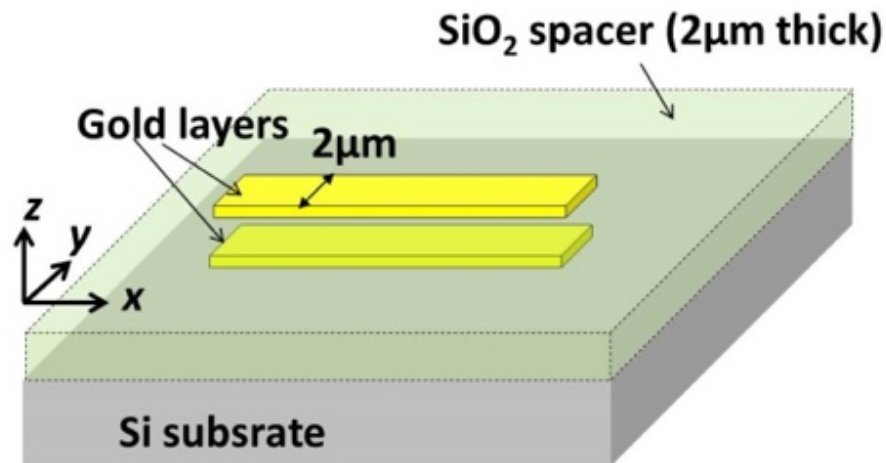
*Conductive carbon fibres:  
6.5 µm diameter, 50-250 µm long*



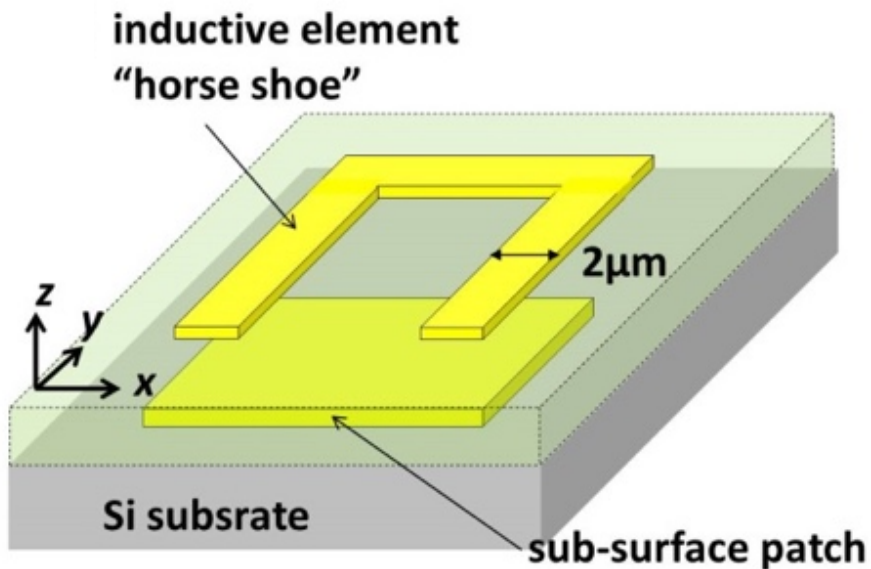
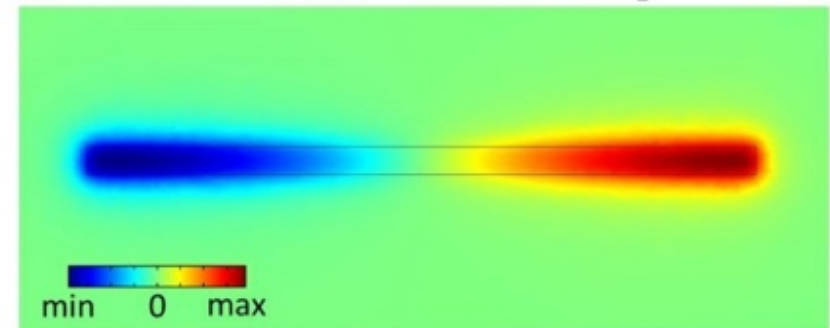


*Symmetry changes from anti-symmetric to symmetric*

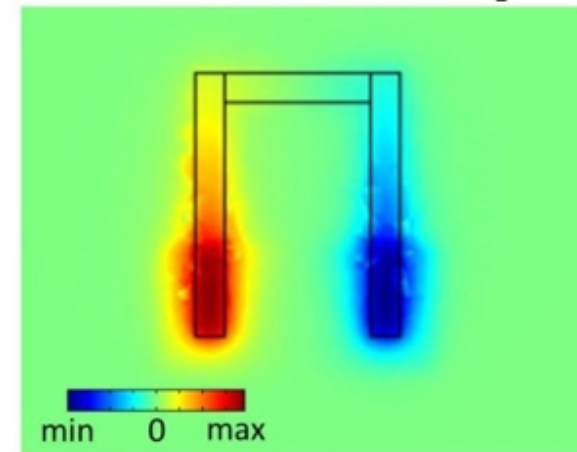


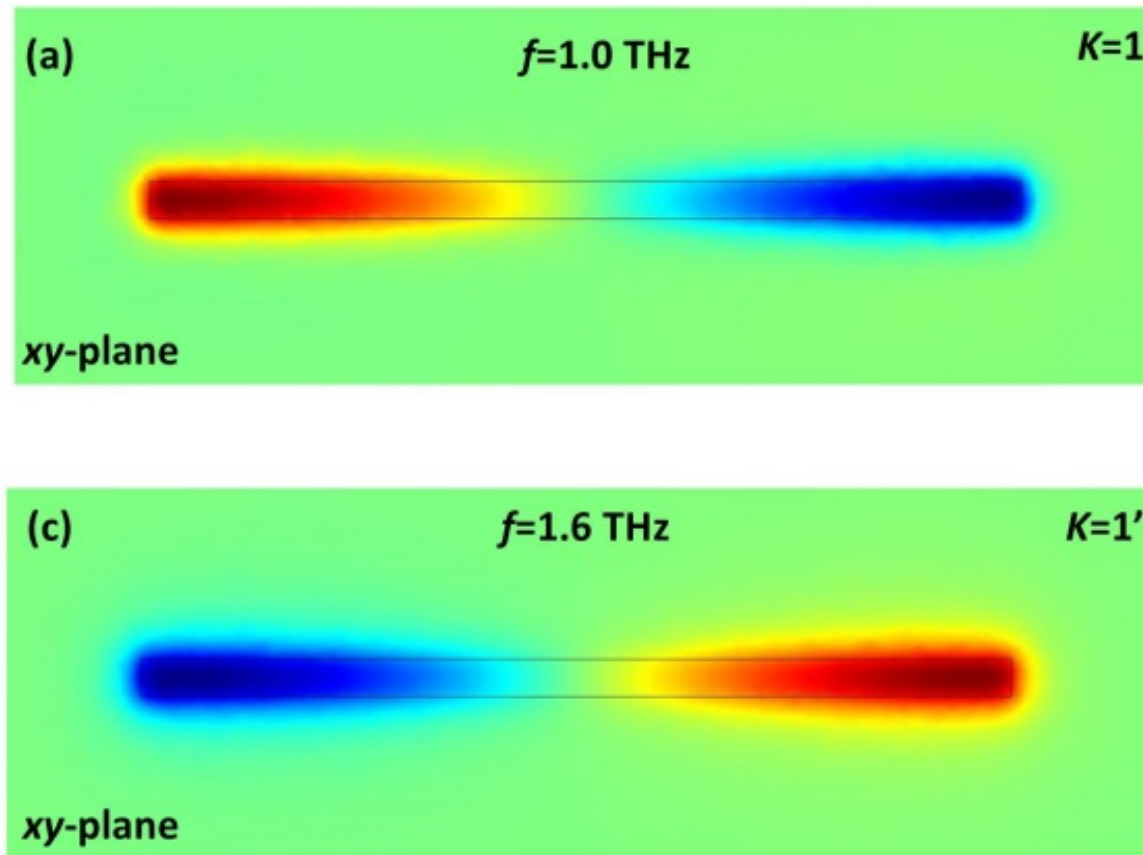


Vertical electric field  $E_z$



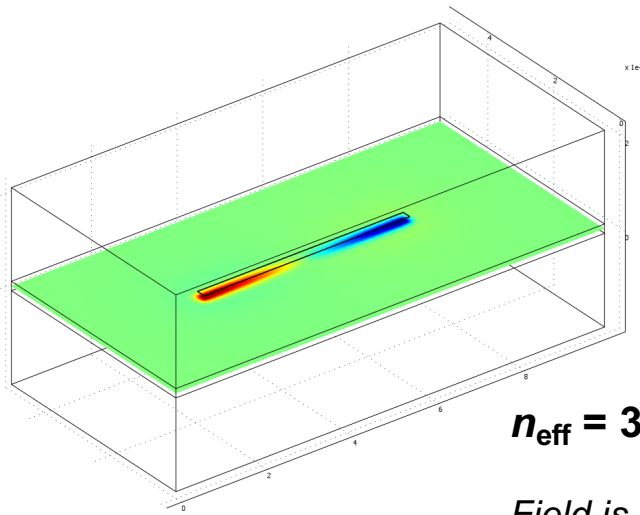
Vertical electric field  $E_z$





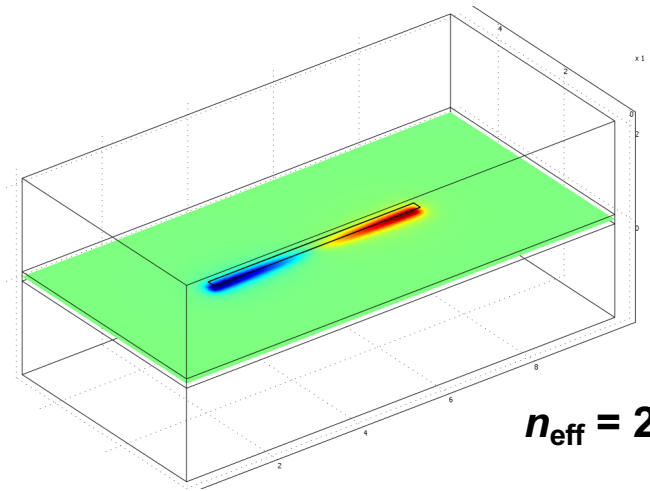
EM simulations reveal two modes with identical field distribution but at different frequencies...

# Two modes with the same spatial distribution



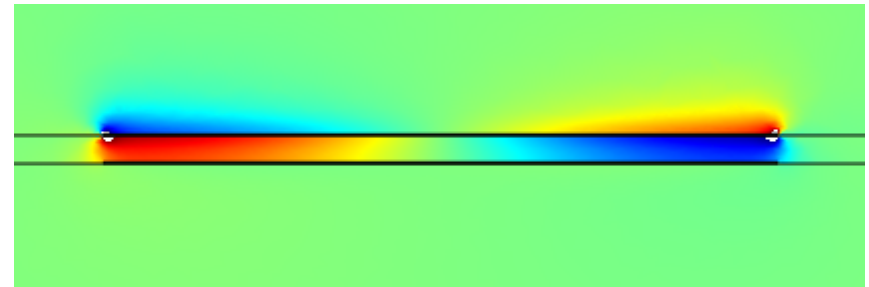
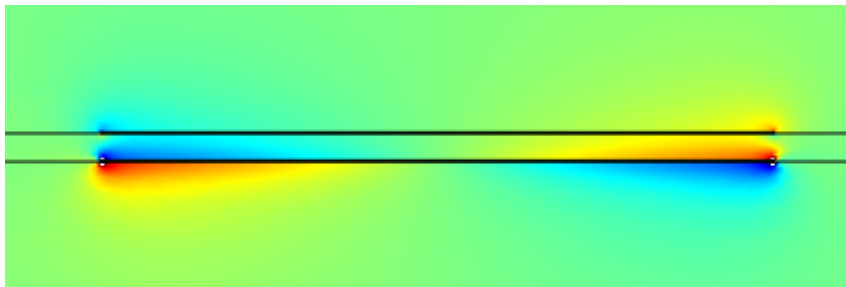
$$n_{\text{eff}} = 3.1$$

*Field is in  
high index  
substrate*

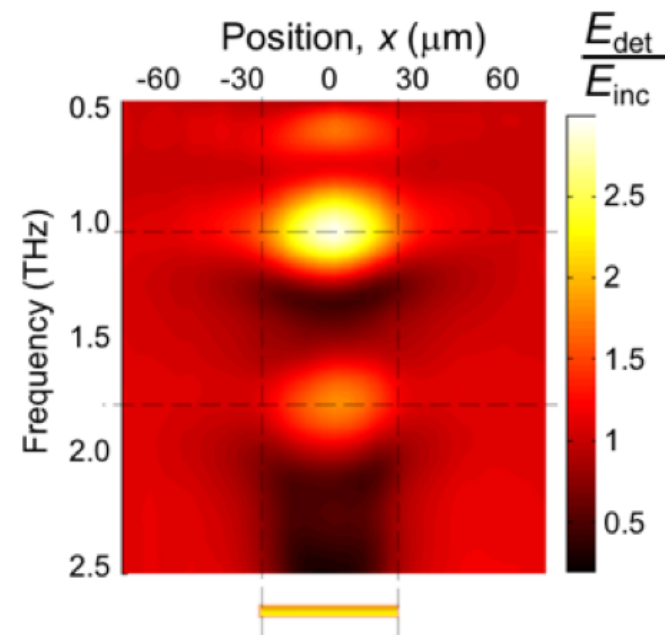
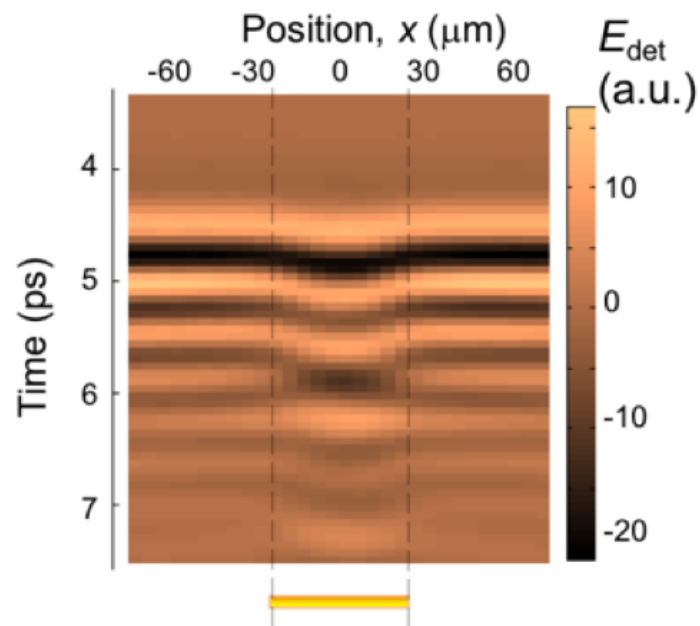
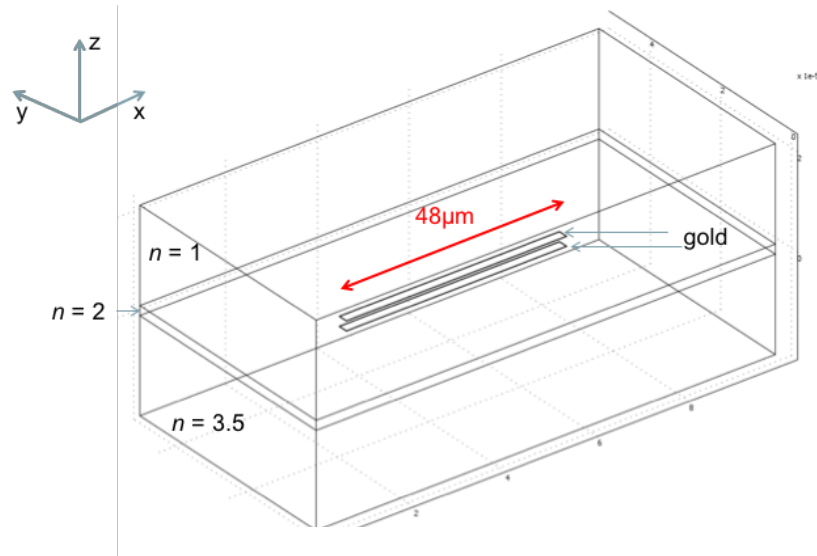


$$n_{\text{eff}} = 2.1$$

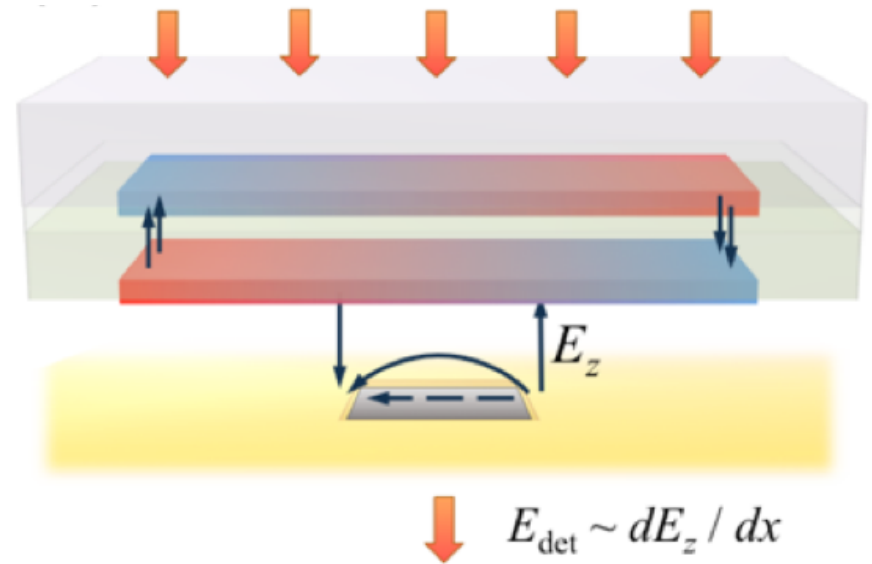
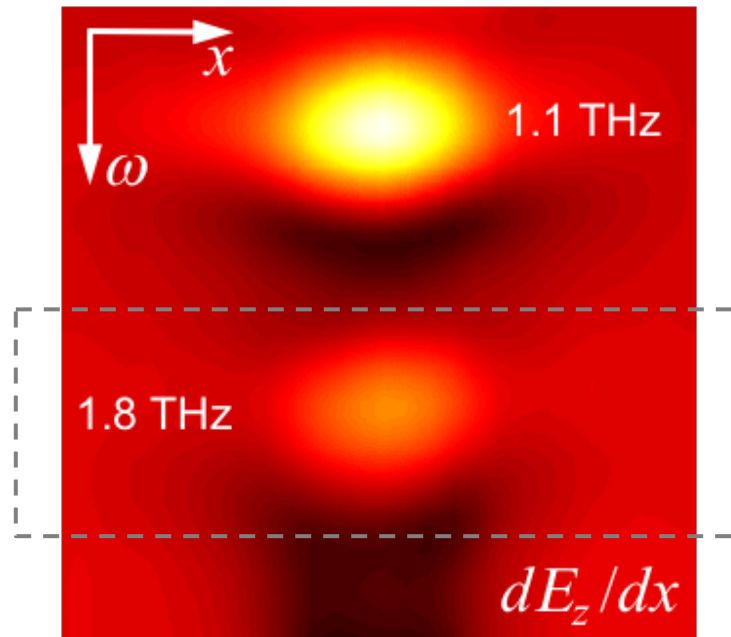
*Field is in  
low index  
spacer*



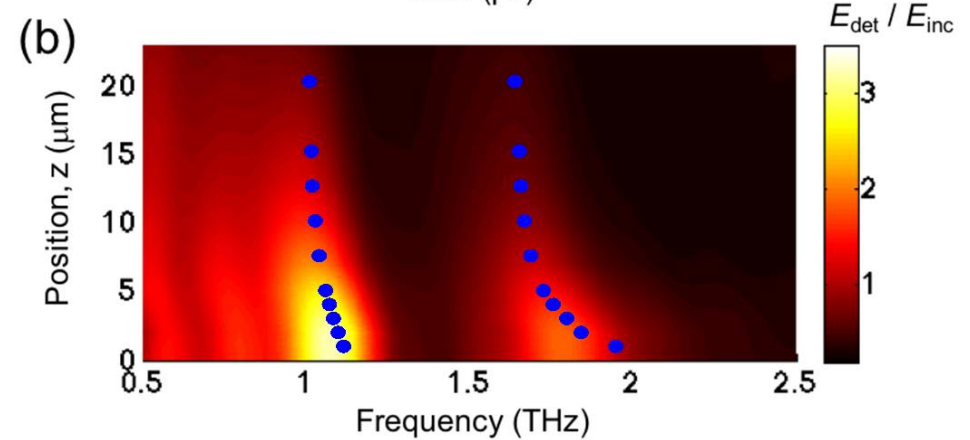
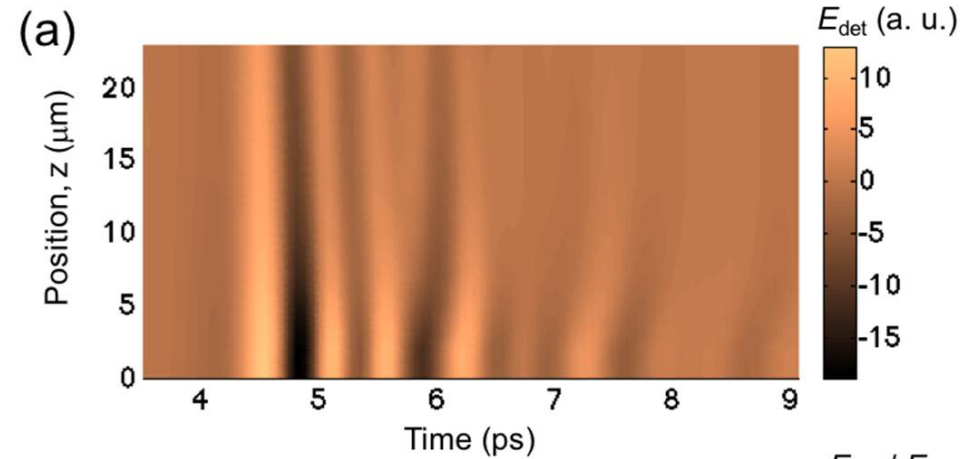
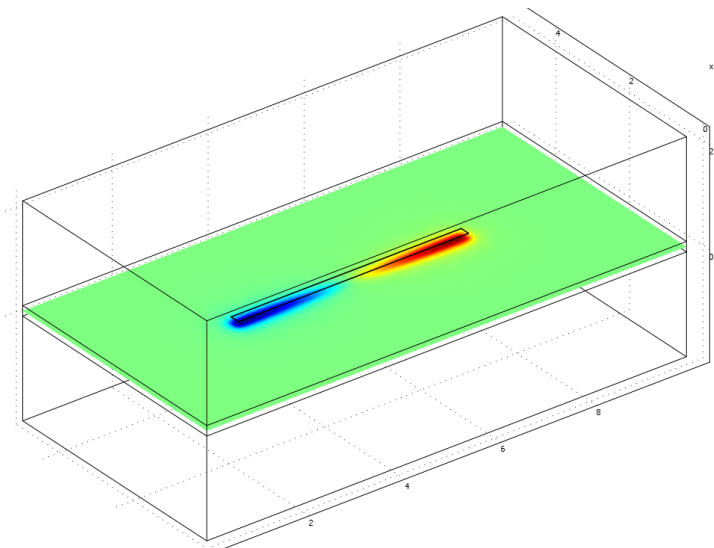
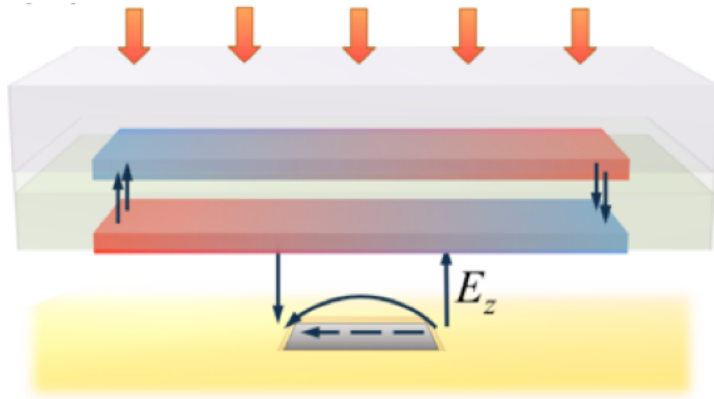


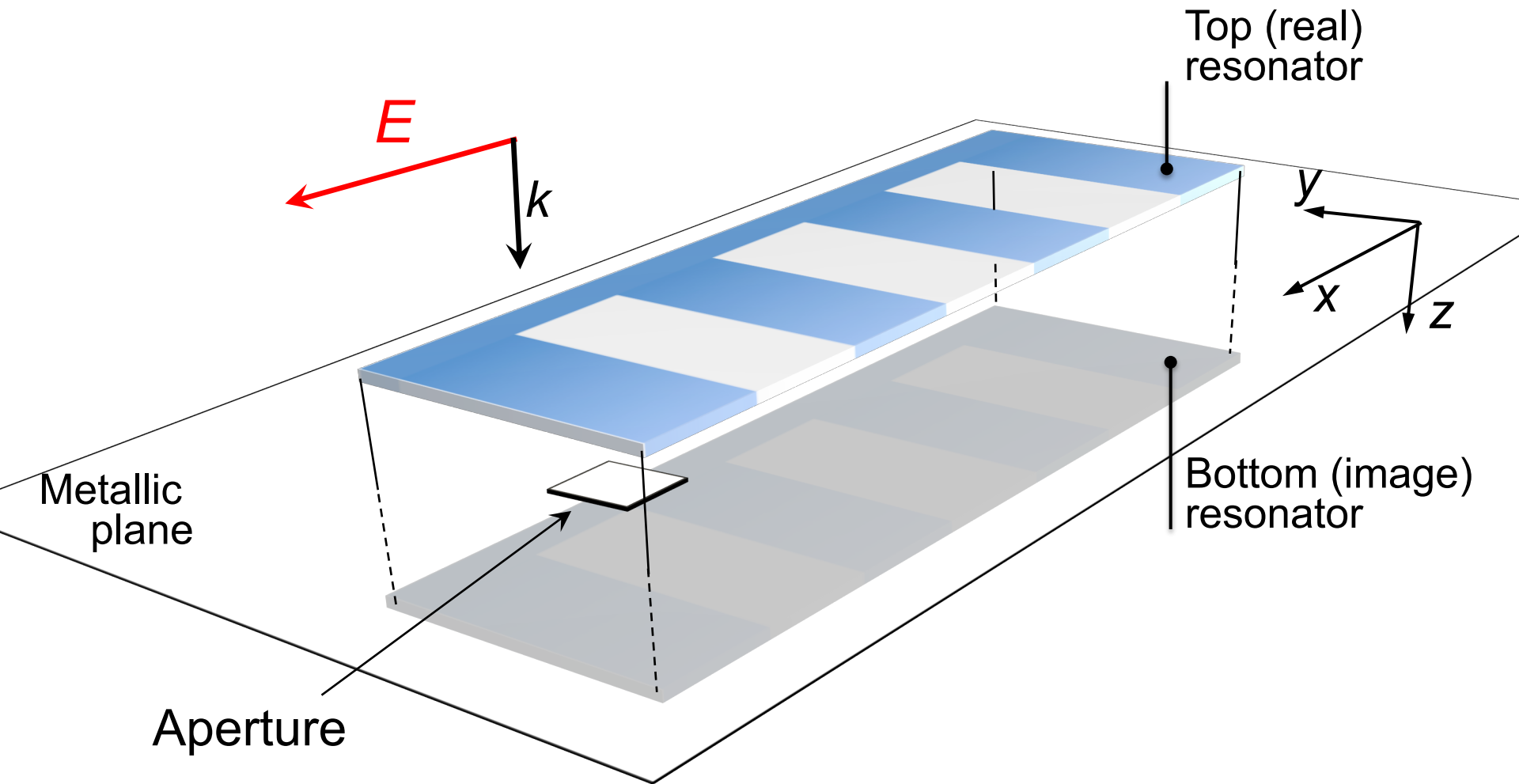


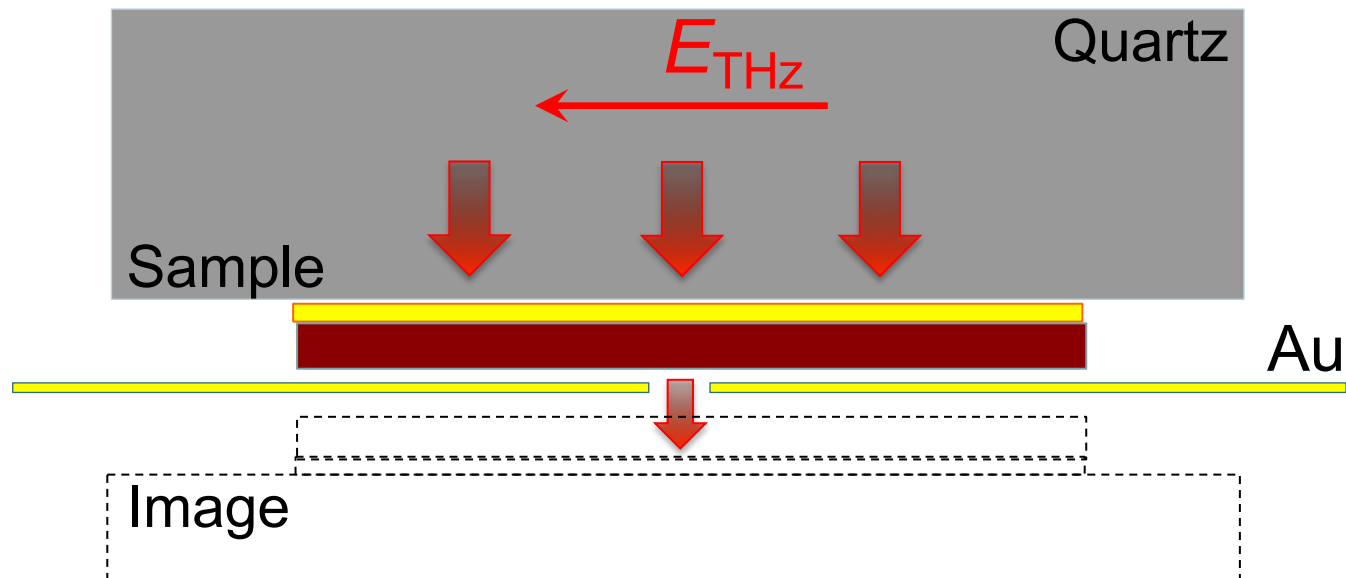


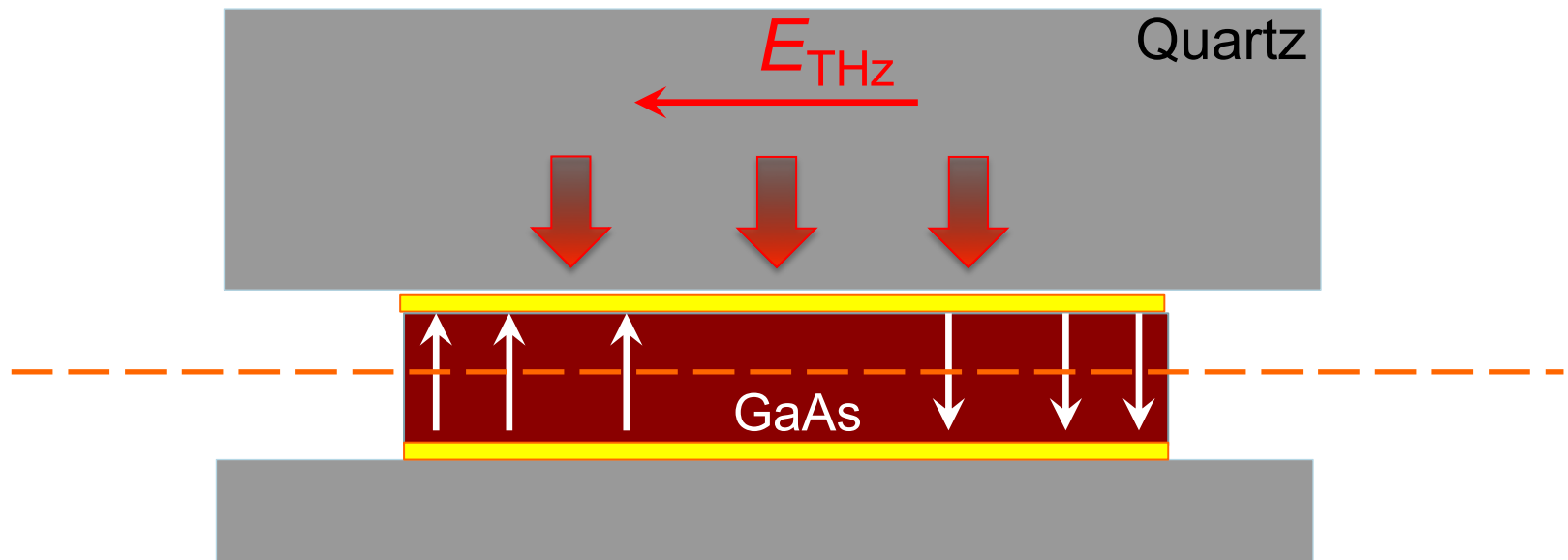


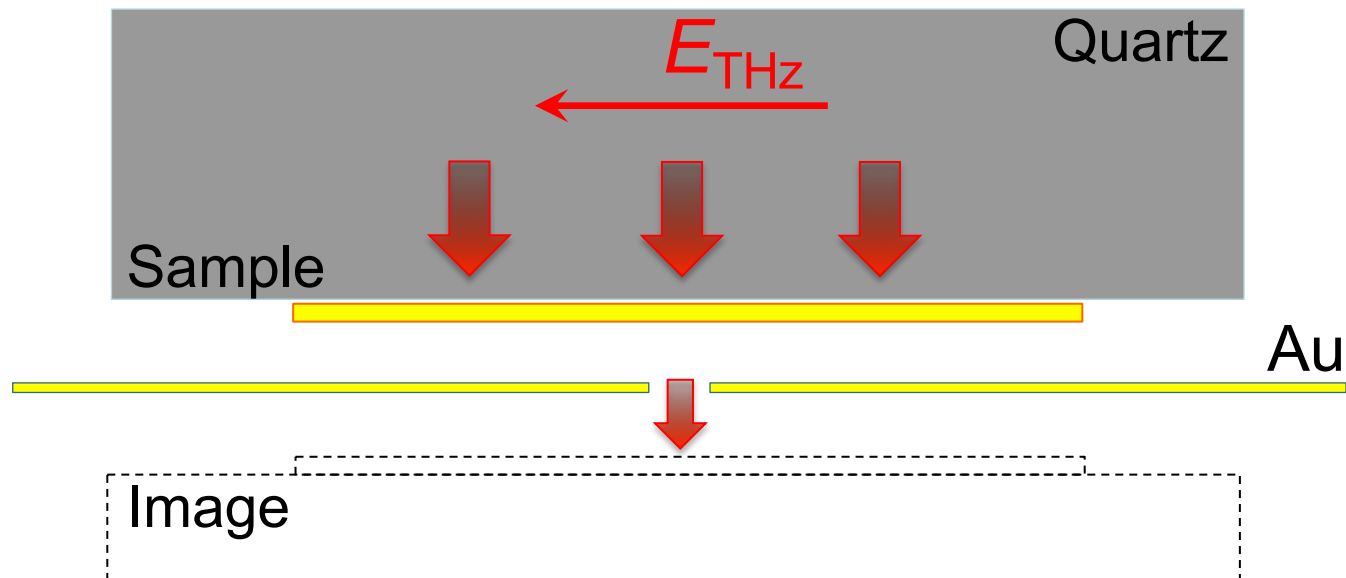
## *Near-field probe functionality*

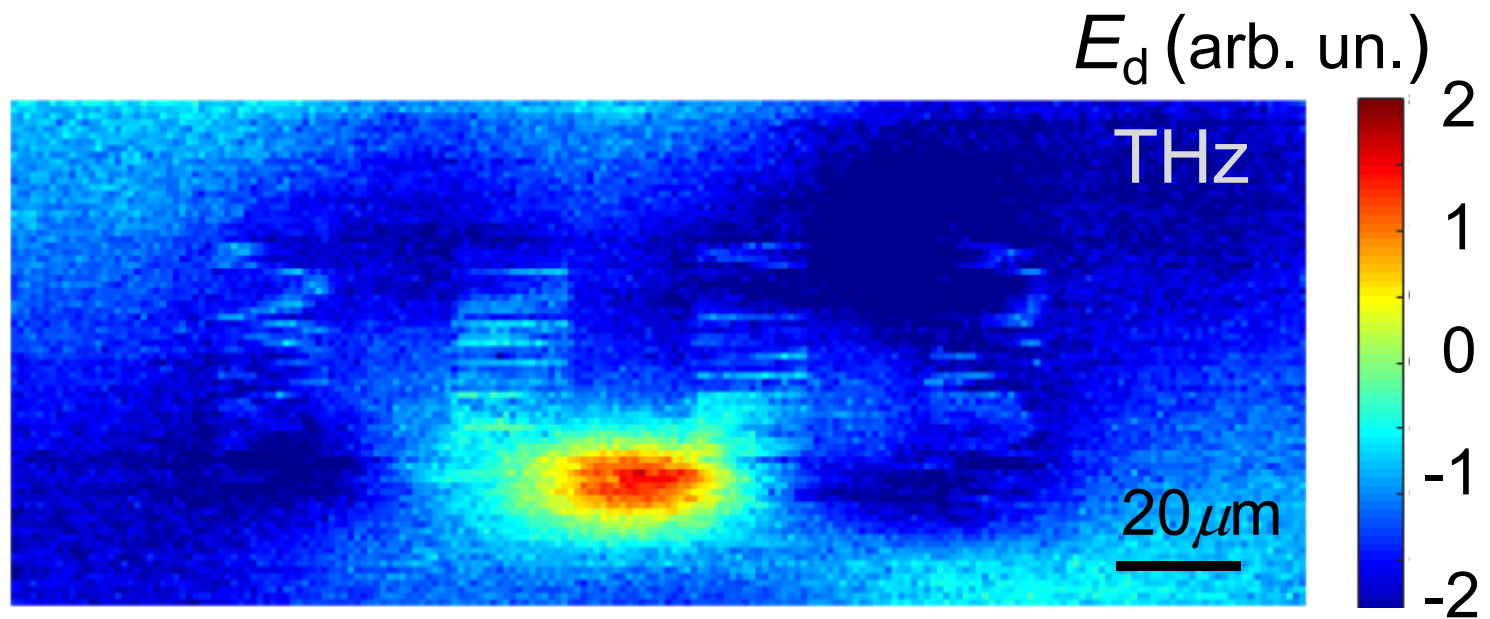




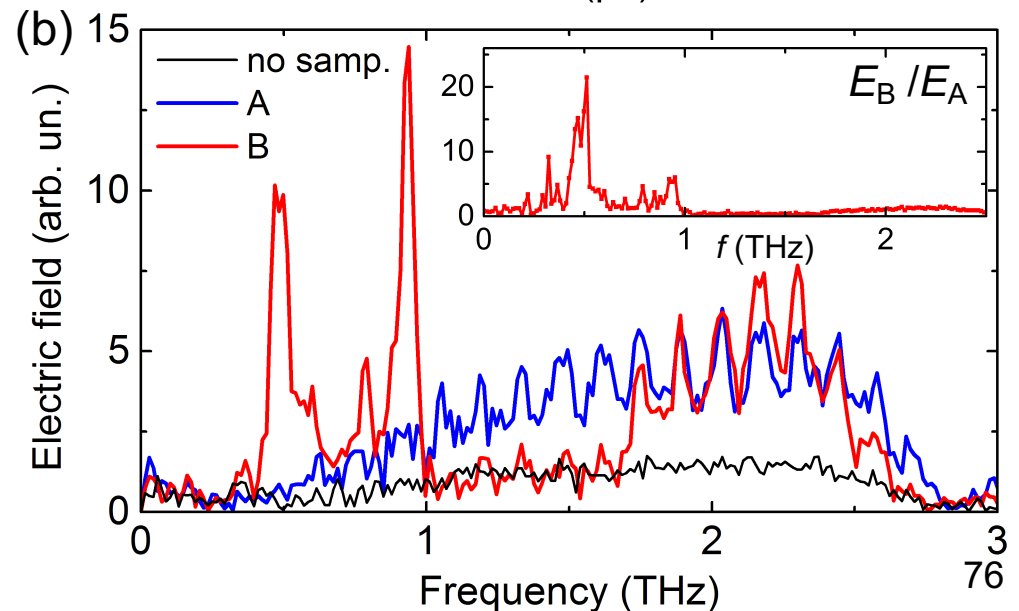
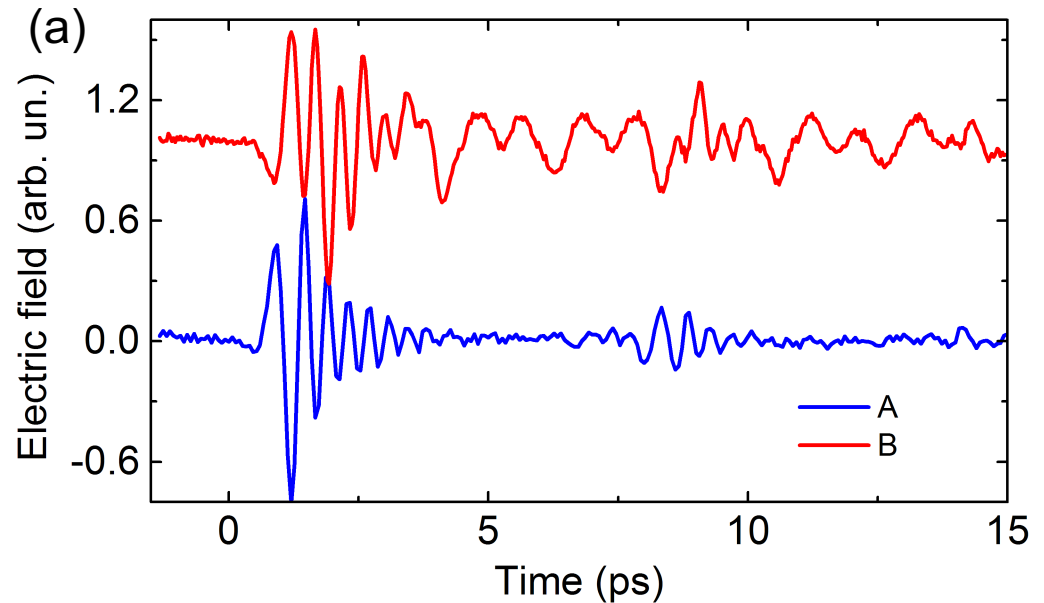
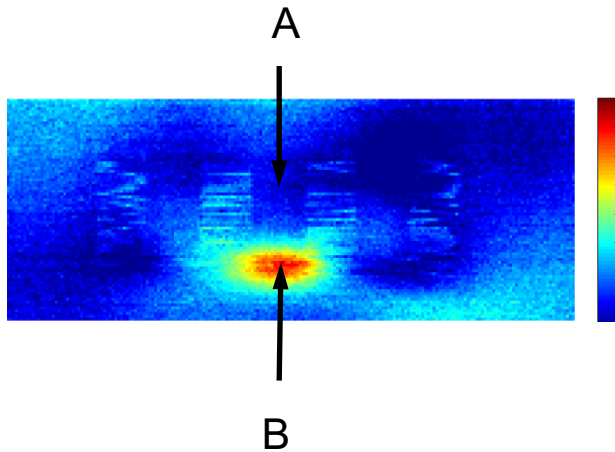






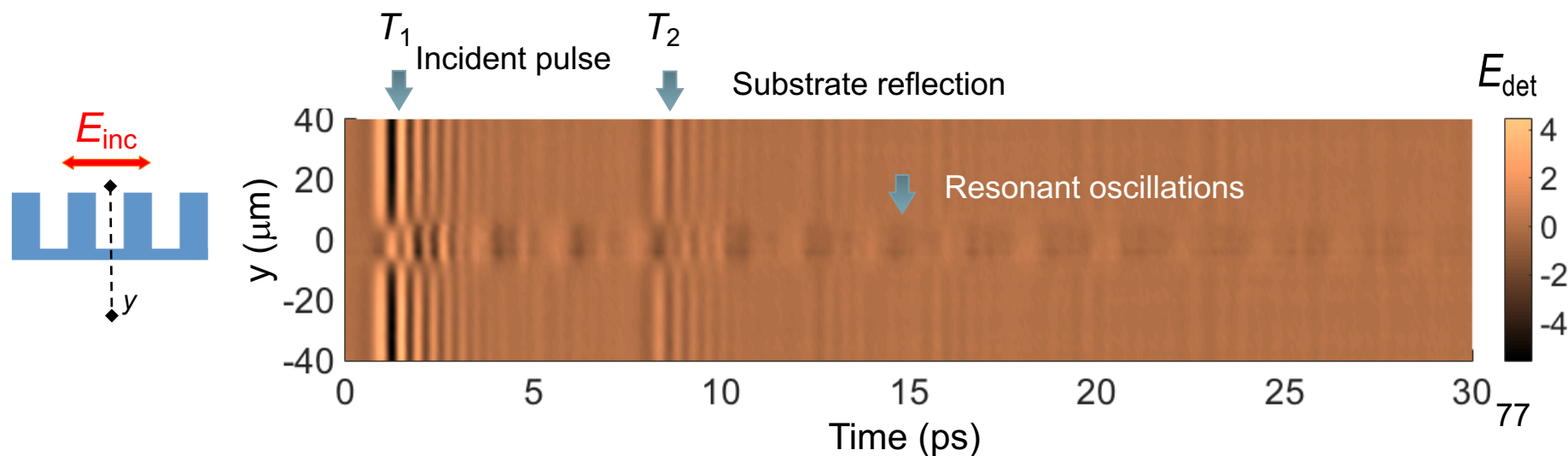
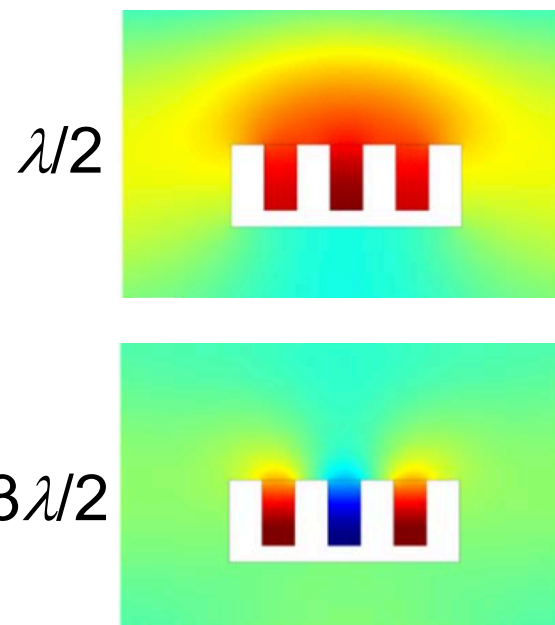
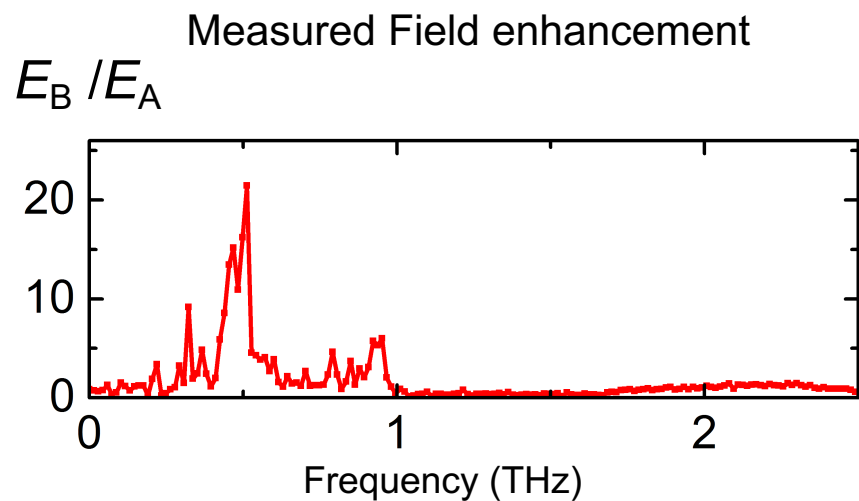


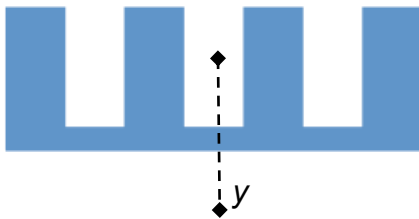




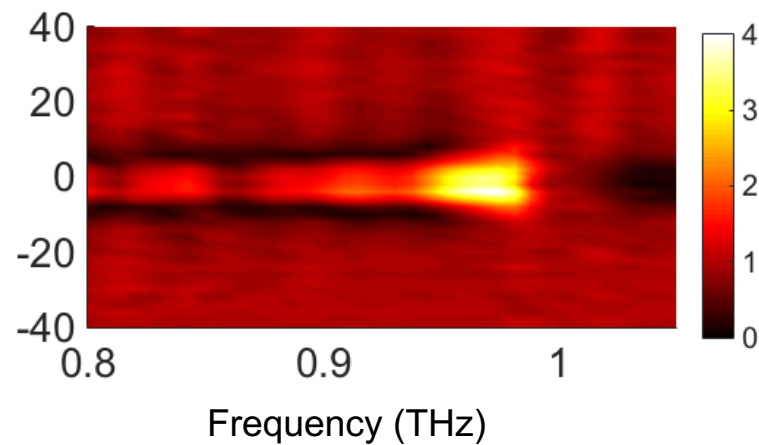
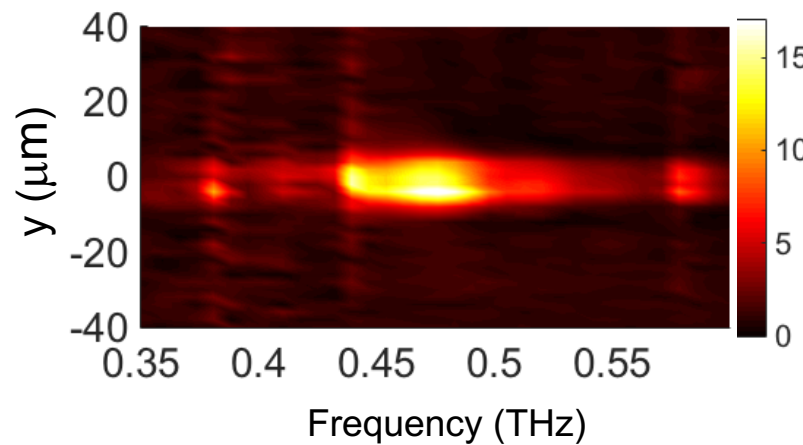
Spectral features:

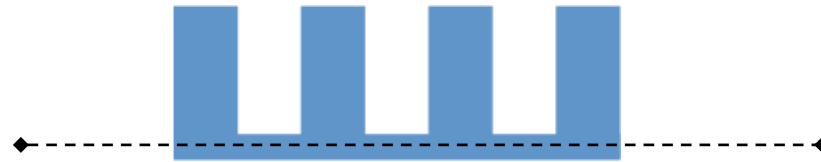
1. Resonance at ~0.5 THz  
(the enhancement is quite high ~20)
2. Resonance at ~1 THz  
(the field enhancement ~5)
3. Field suppression 1-1.7 THz
4. Broad band at > 2 THz



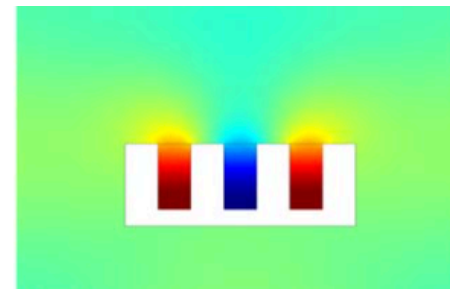
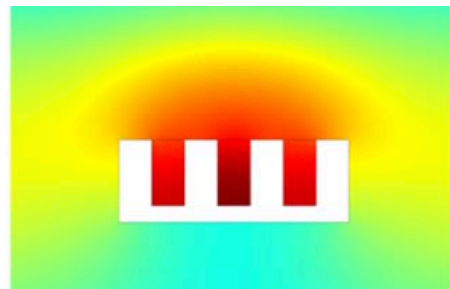
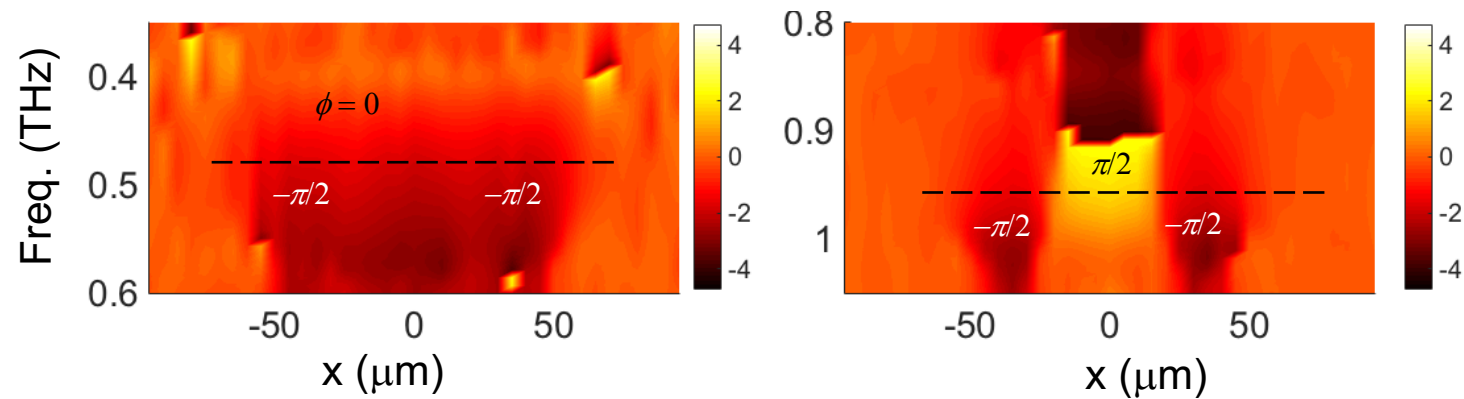


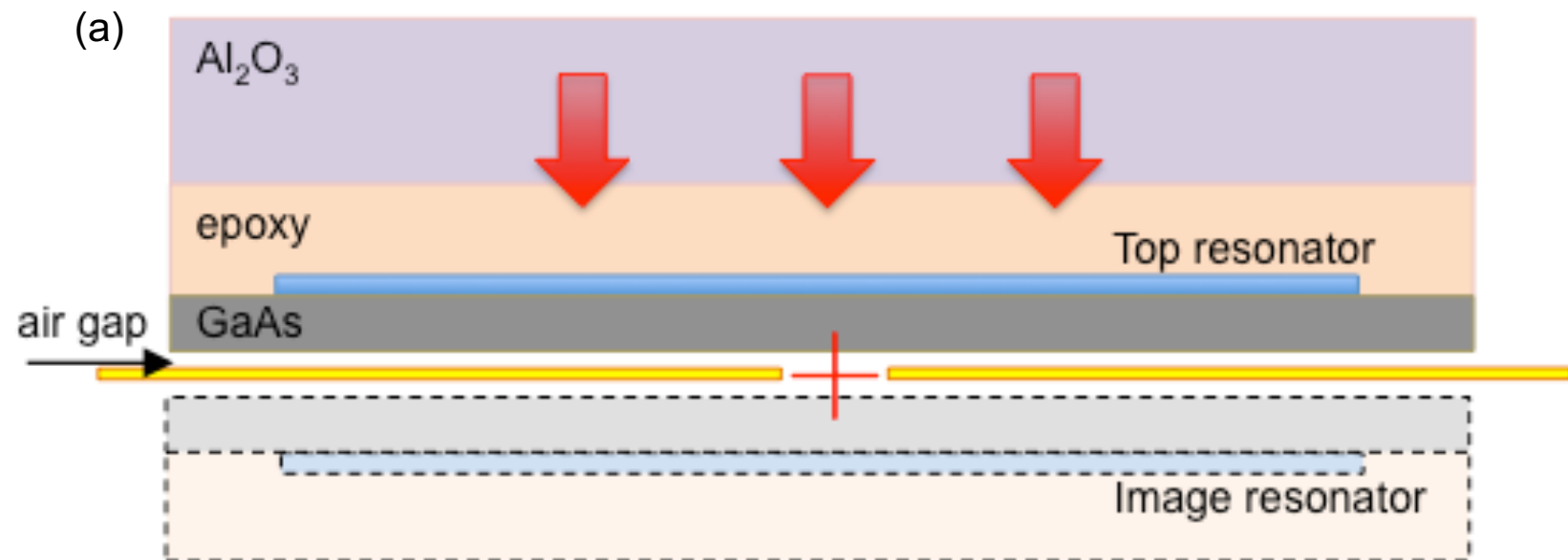
Field Amplitude (norm.)

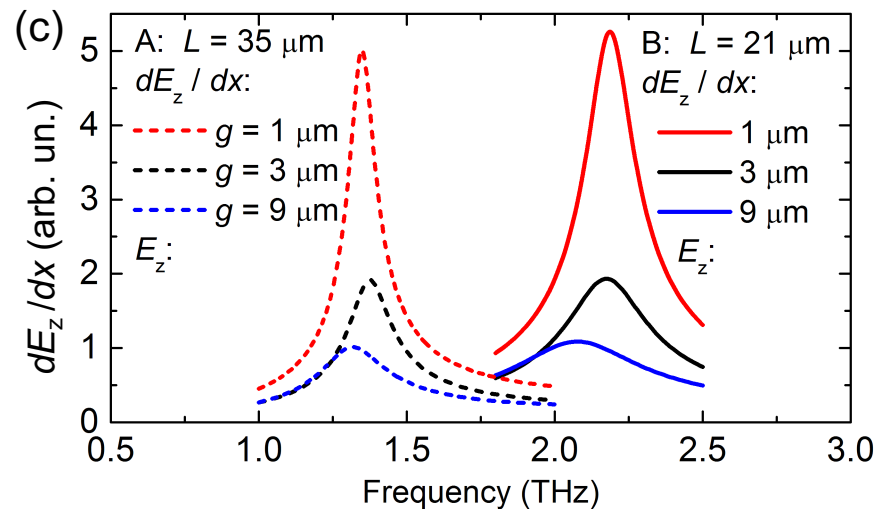
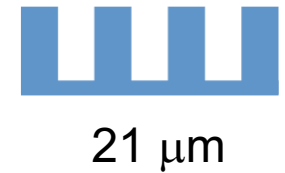
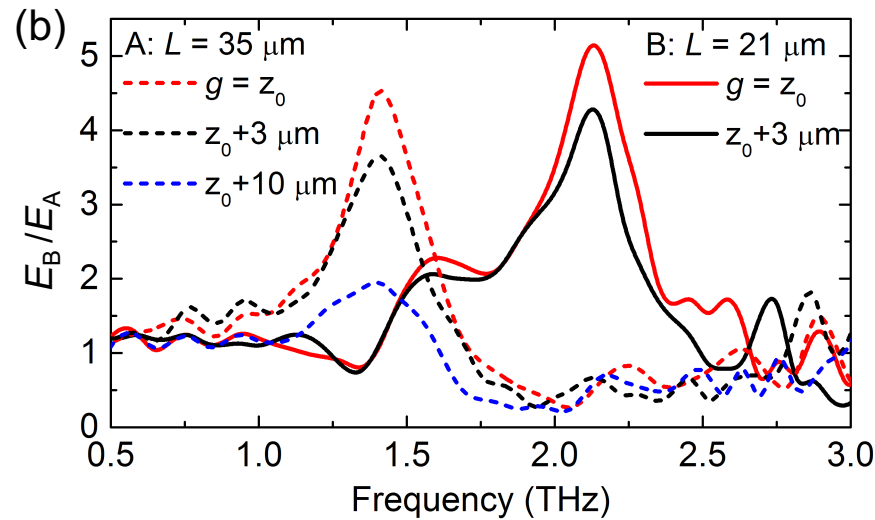
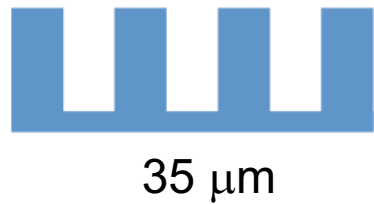




Phase,  $\phi$  (radians)

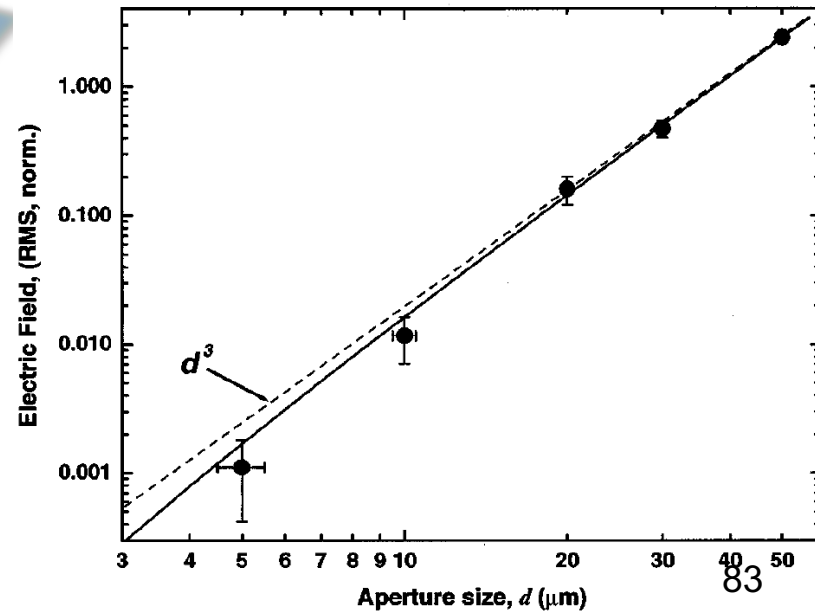
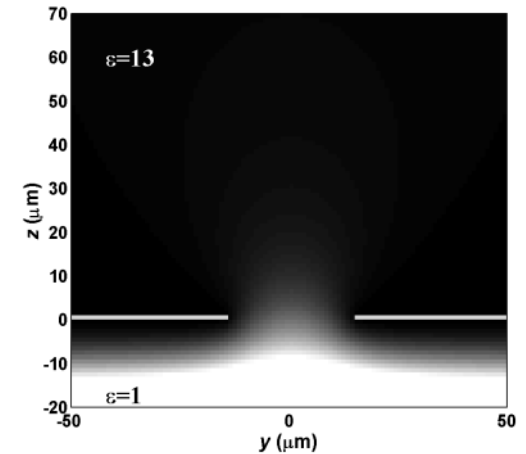
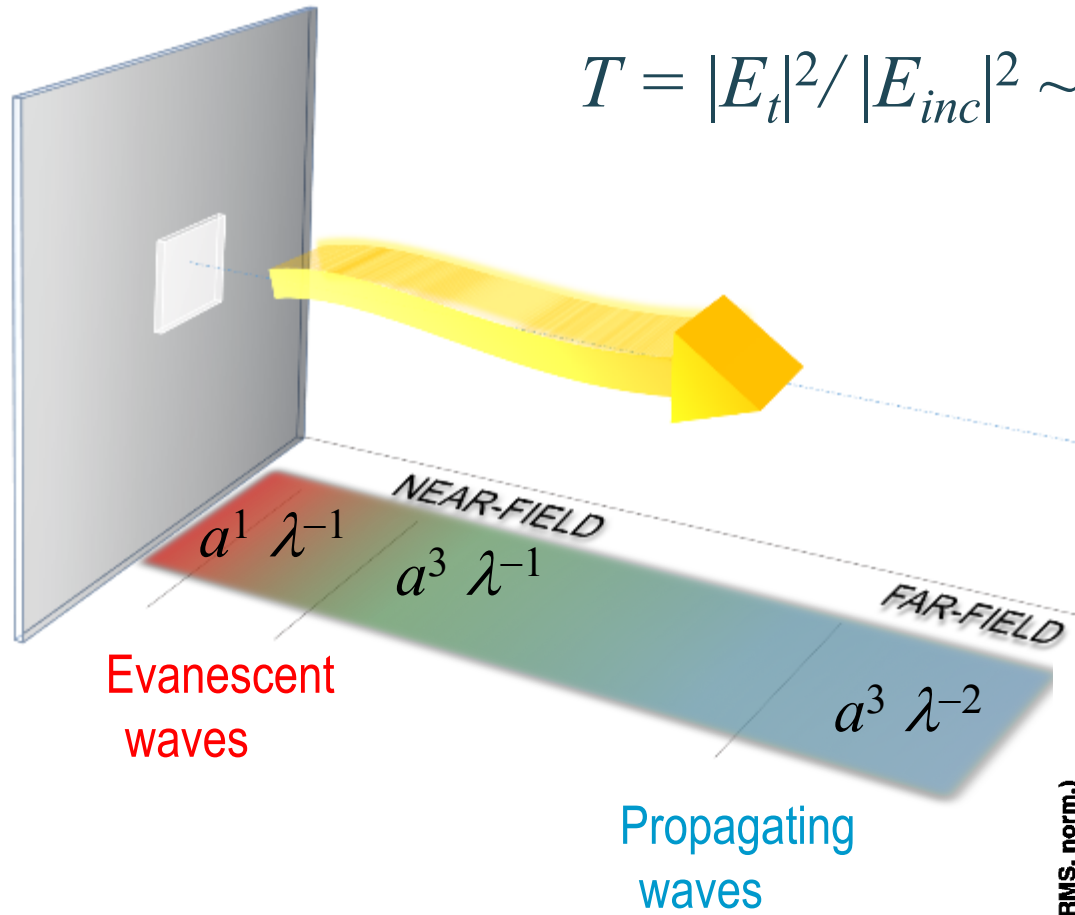






# *NF Probes with embedded detectors*

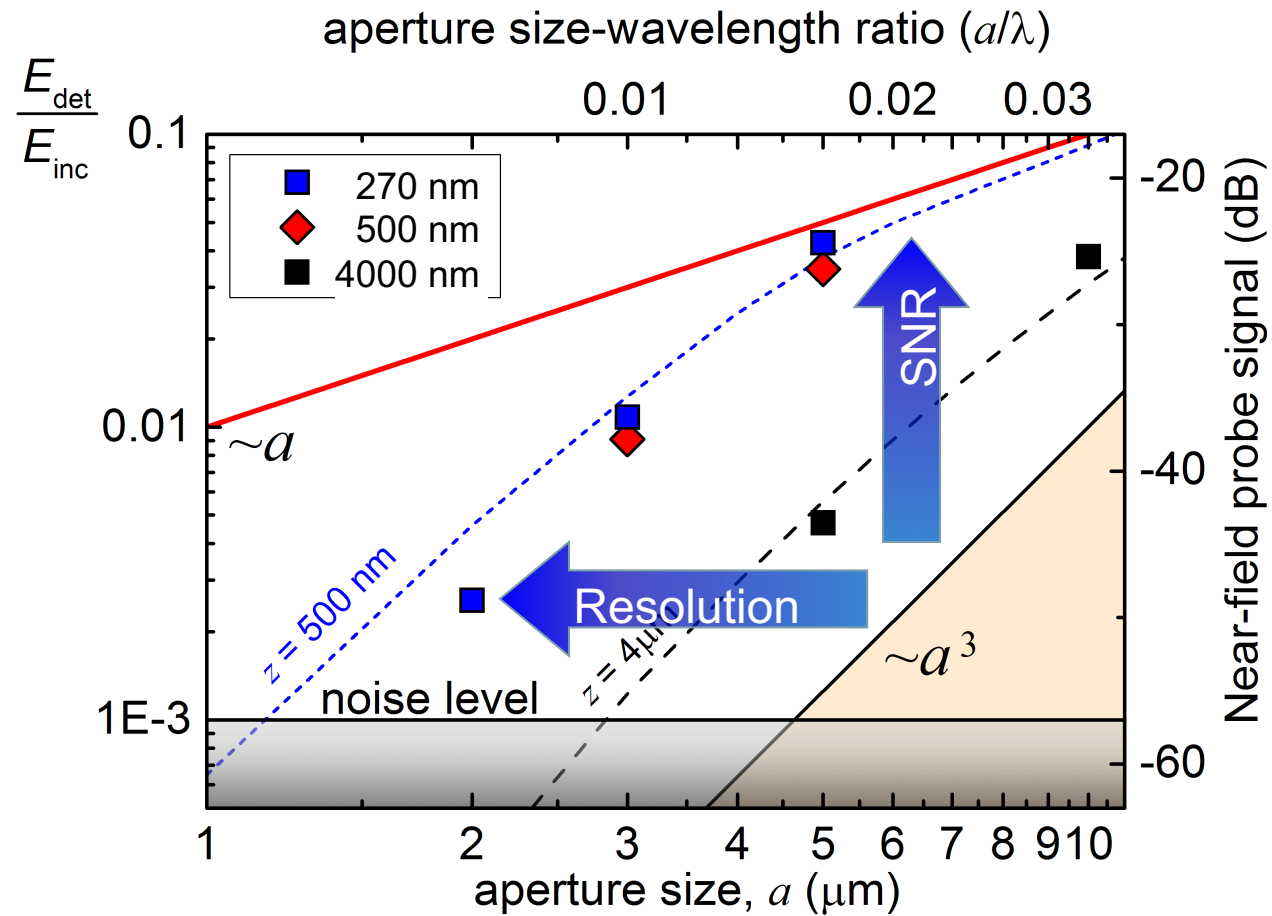
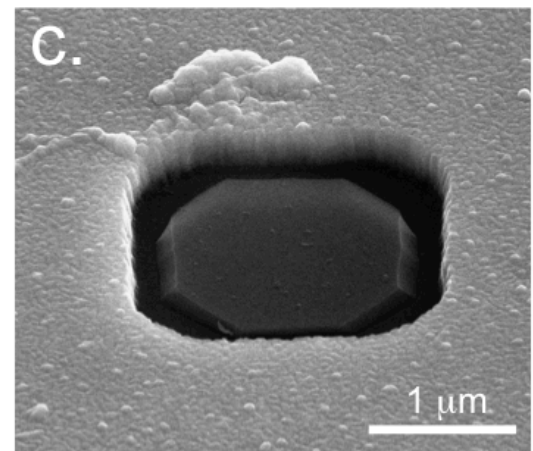
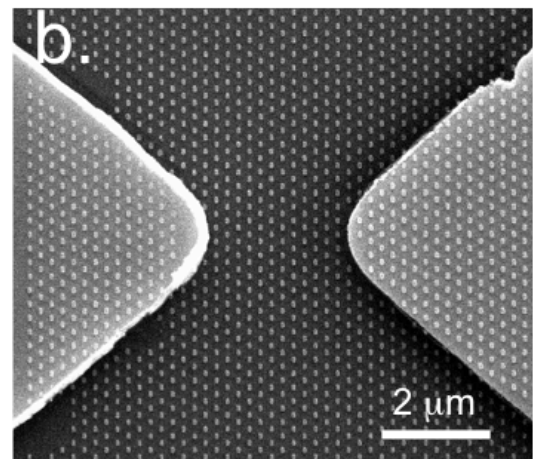
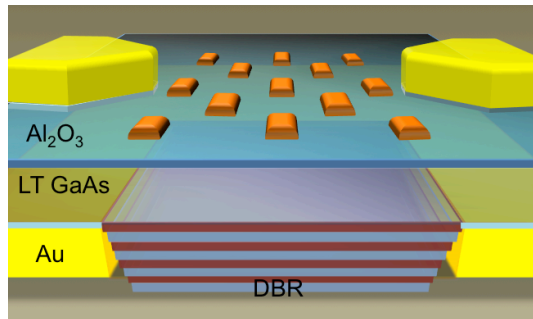
*Spatial resolution and sensitivity*



Mitrofanov et al., APL **77**, 3496 (2000)  
APL **79**, 907 (2001)

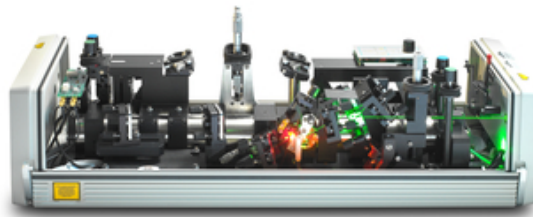
Adam, J IRMTW **32**, 976 (2011)





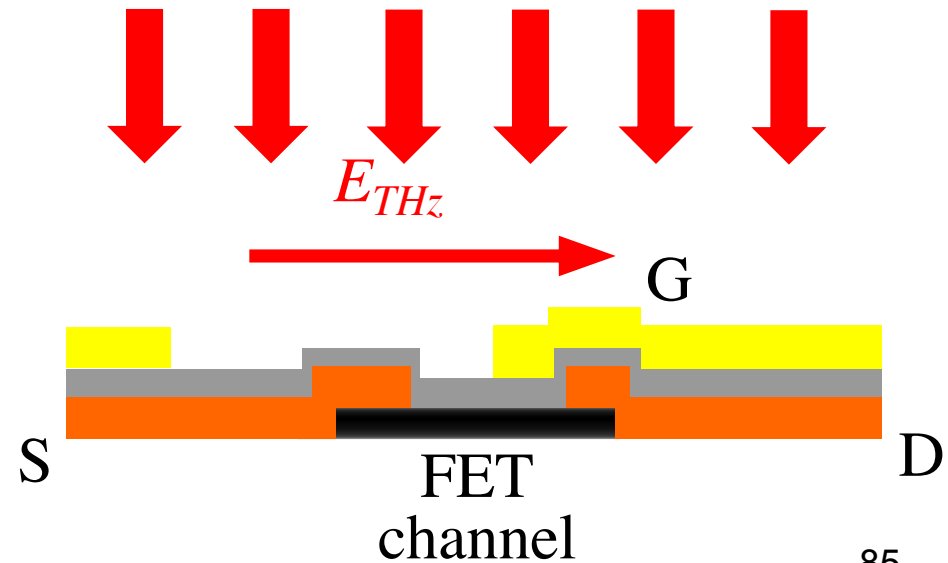
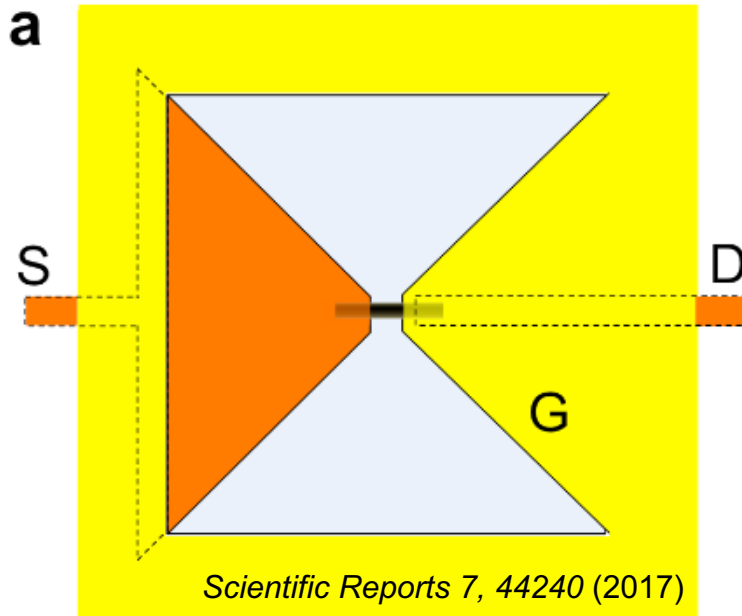
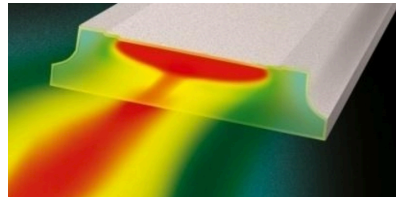
Mitrofanov et al.,  
IEEE Trans. THz S&T (2016)

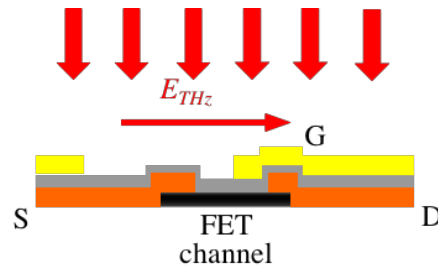
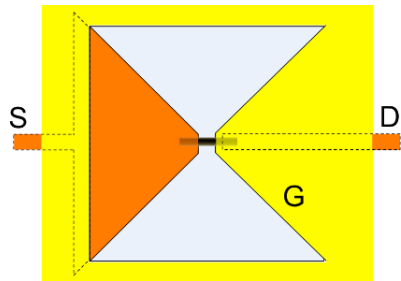
Mitrofanov et al.,  
ACS Photonics (2015)



Femtosecond pulse laser  
Spectra-Physics

THz Quantum Cascade Laser  
CNR





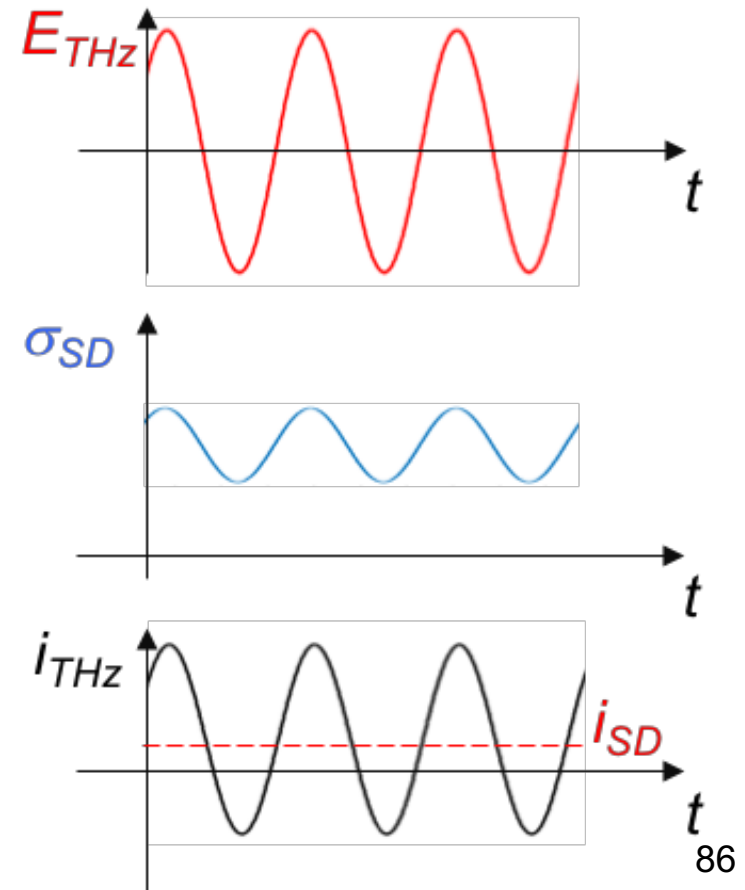
*THz modulation  
of charge carrier density*

Asymmetric coupling of THz wave to the channel leads to a DC current or voltage across the channel

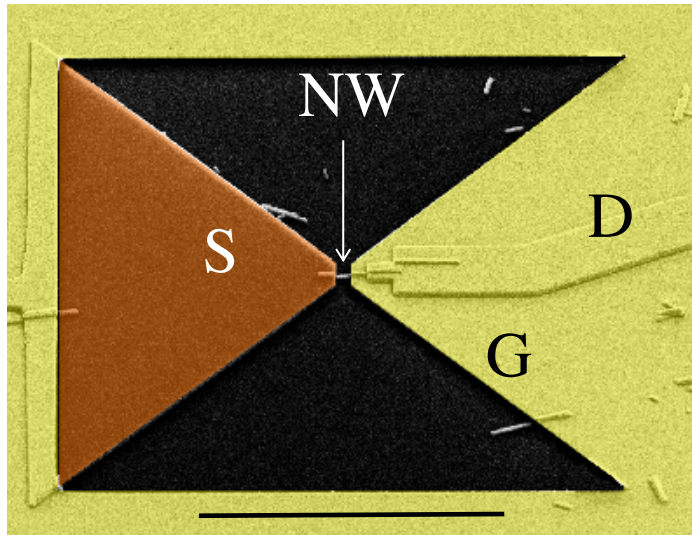
Detection mechanisms:

FET Rectification effect

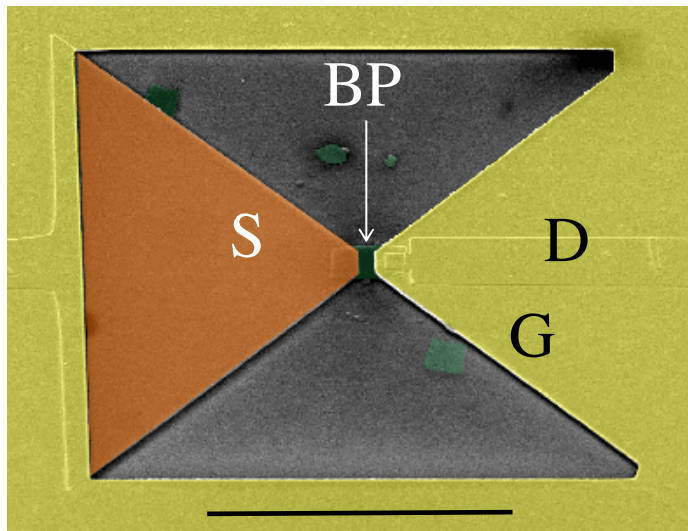
Thermoelectric effect



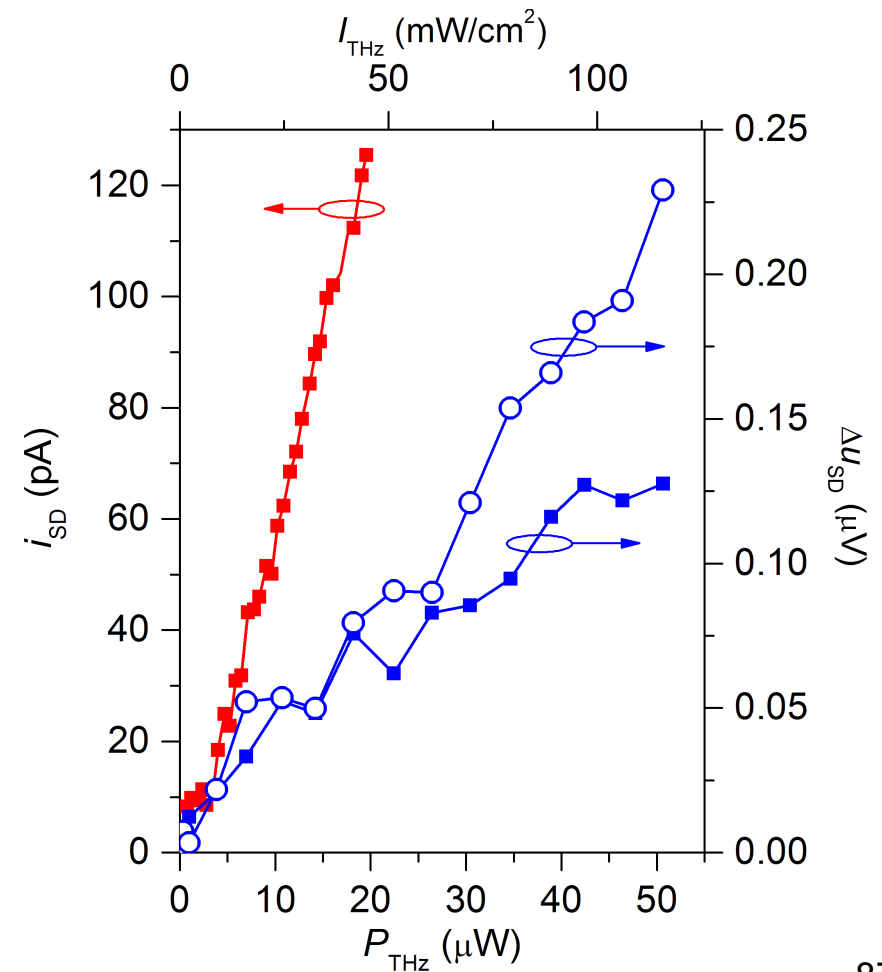
InAs nanowires and Black Phosphorus flakes operate as integrated FET THz detectors in PV and rectified current modes

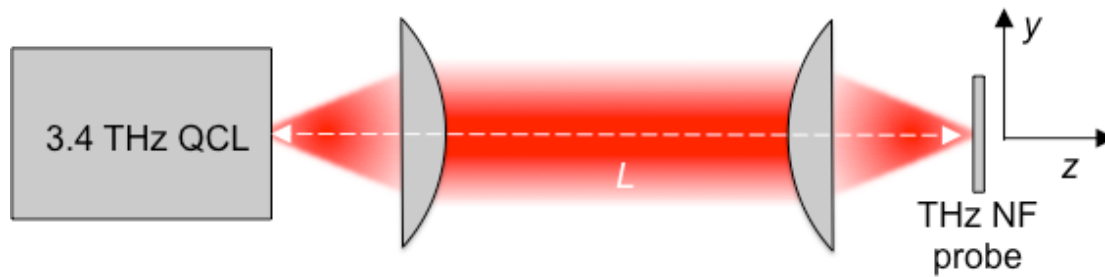


10  $\mu\text{m}$

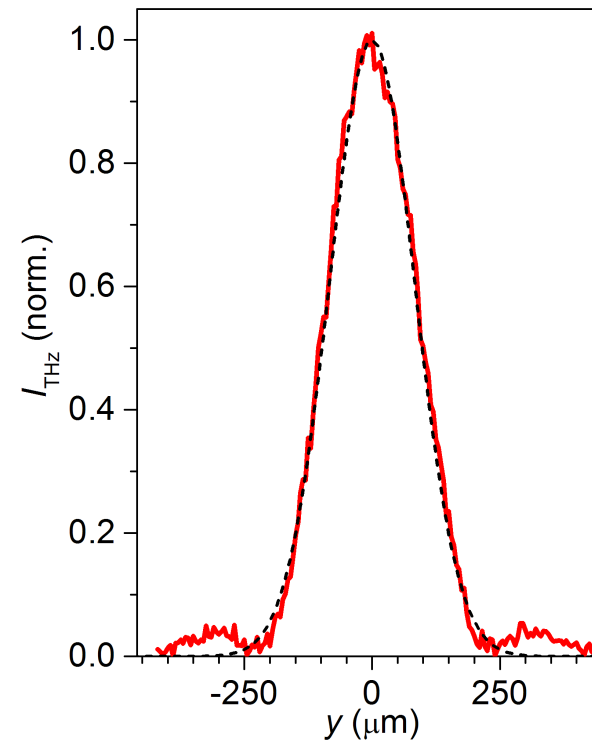
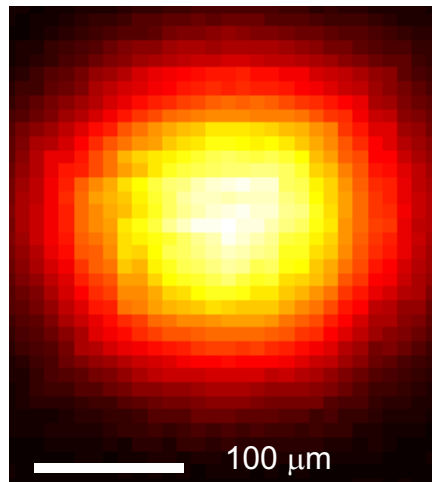


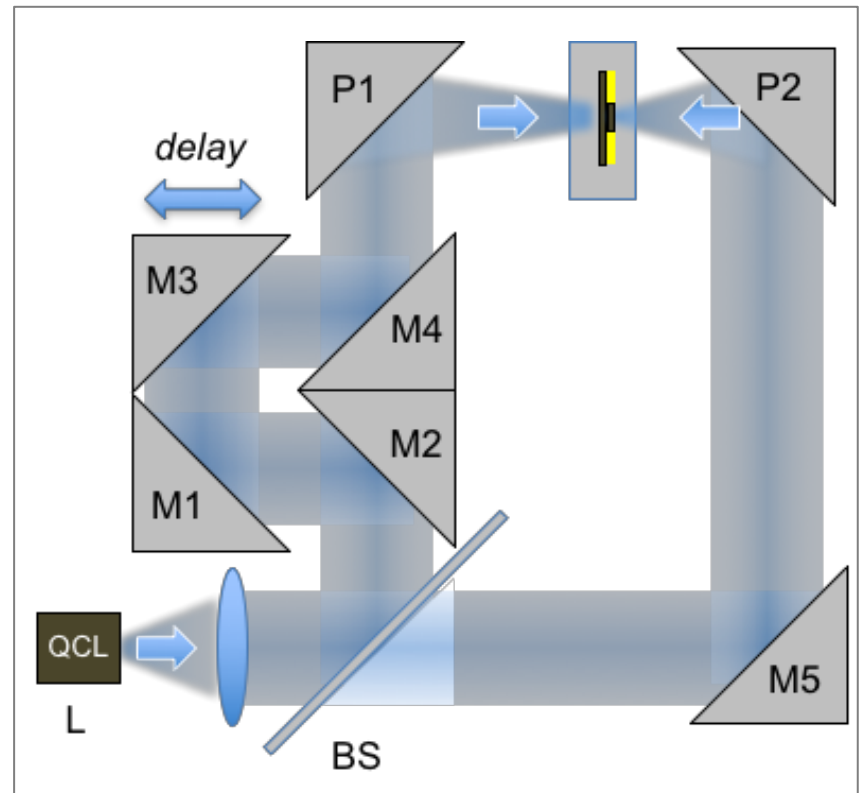
10  $\mu\text{m}$



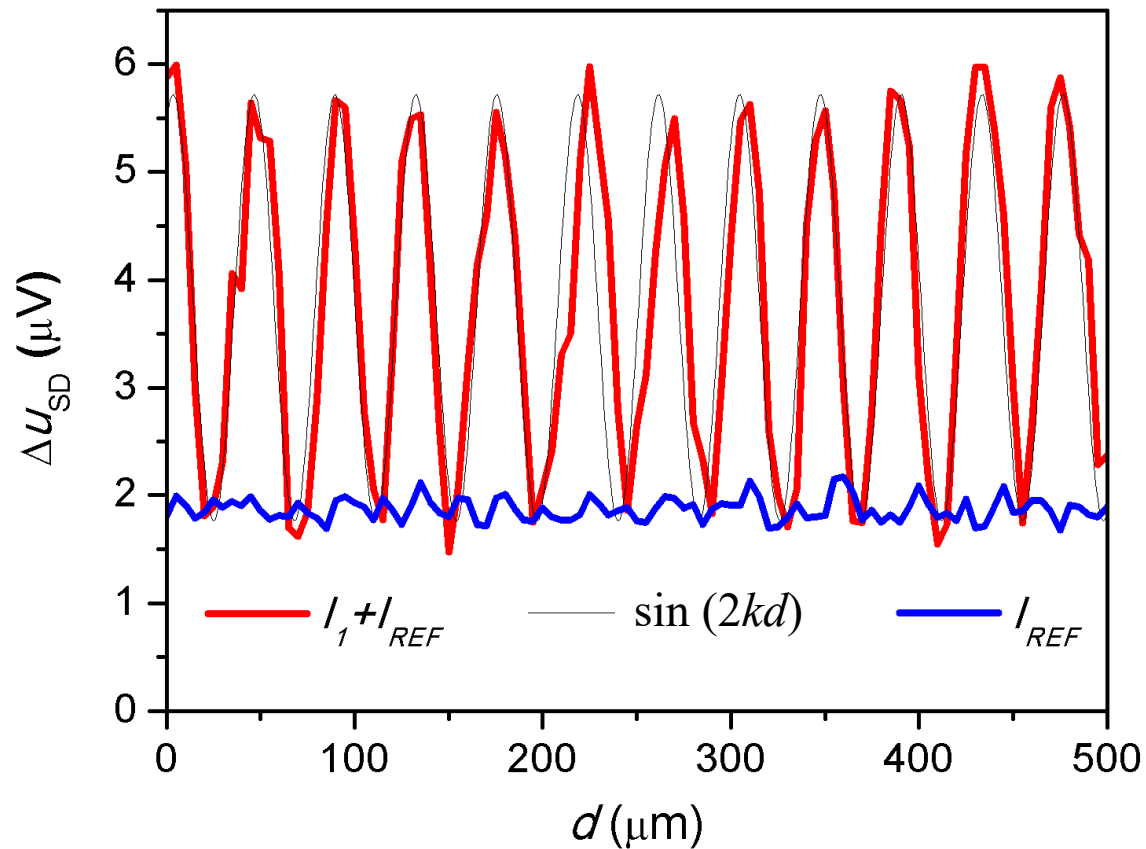


Beam profile in the focal plane



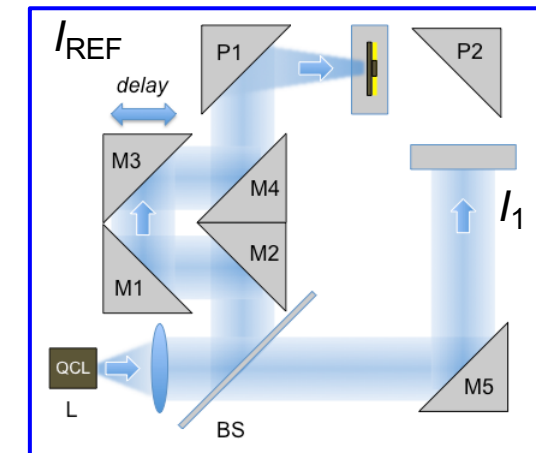
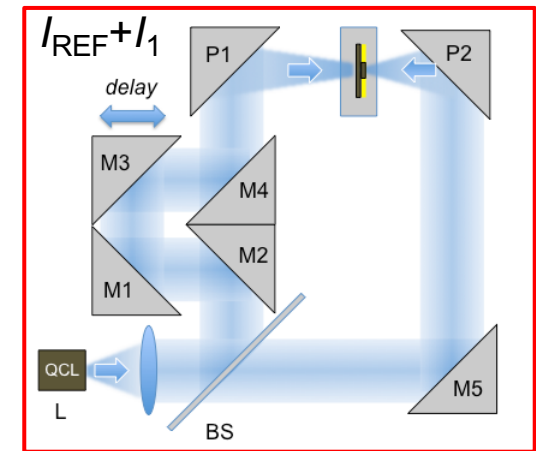


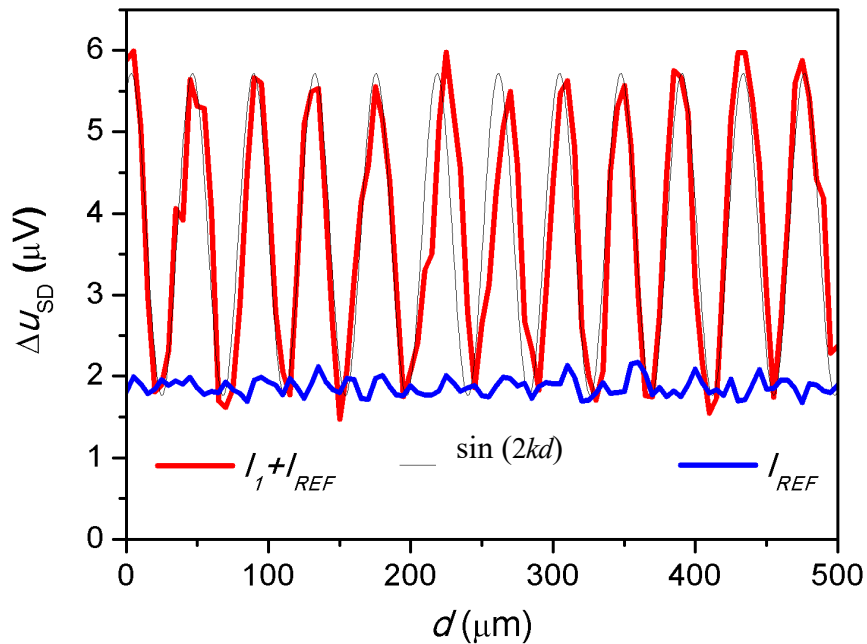
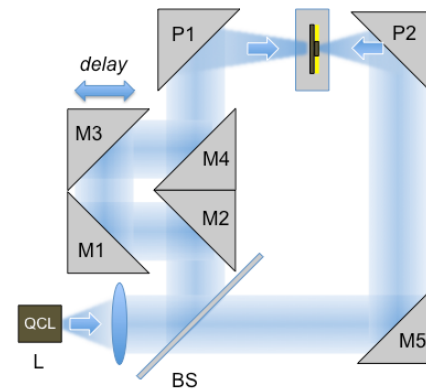
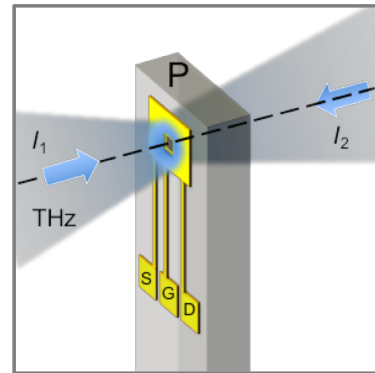
89



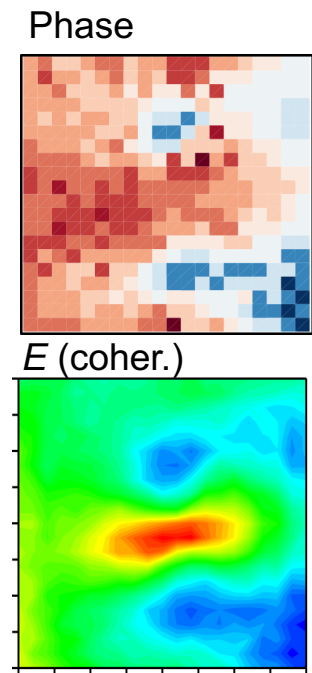
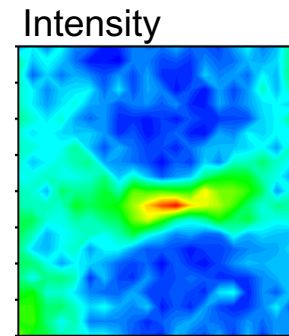
The detected signal depends on the phase

Coherent gain is observed:  $(E_1 + E_{REF})^2 > E_1^2 + E_{REF}^2$





Scientific Reports 7, 44240 (2017)



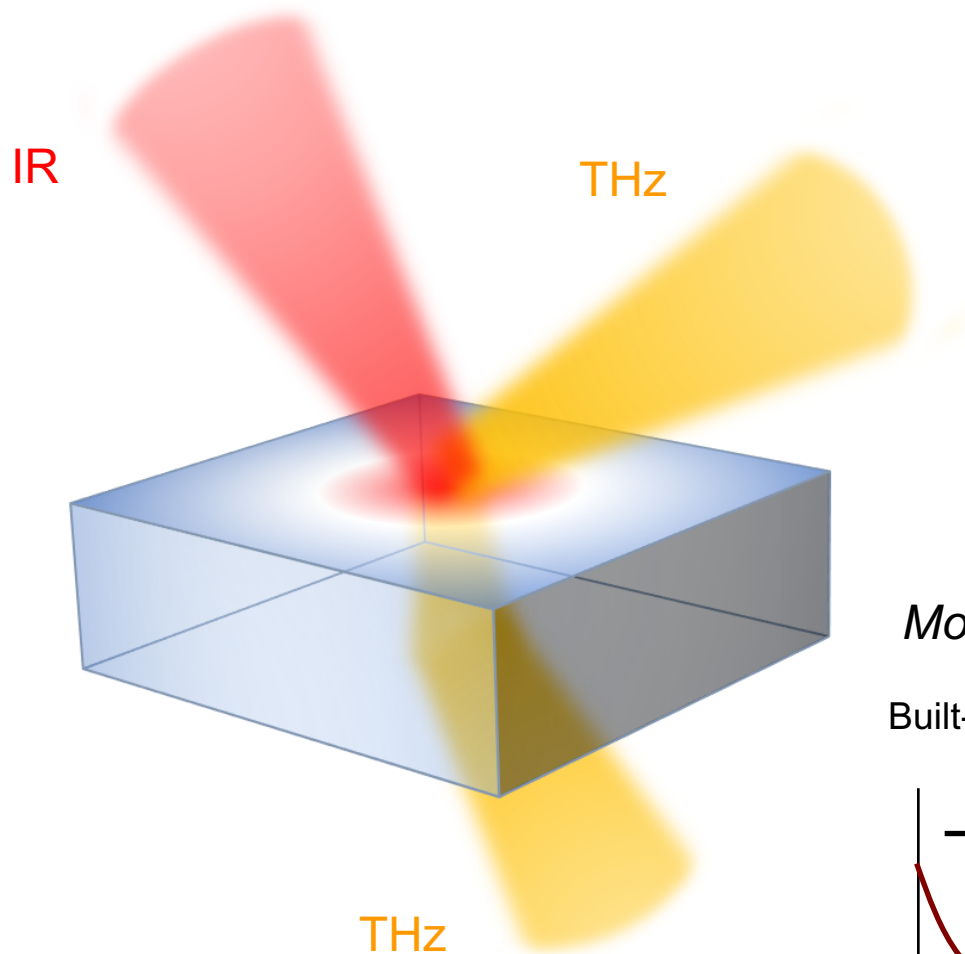
Optica 5, 651-657 (2018)



1. *Near-field light-matter interactions are rich in physics and observation provide invaluable insight.*
2. *There are several near-field techniques available it is ideal to develop a technique with the application in mind*
3. *Resonators can produce THz magnetic fields*
4. *Surfaces waves / surface plasmon resonances can be probed by the aperture type THz microscopy*

## *Extra: NF Imaging Application - THz pulse generation*

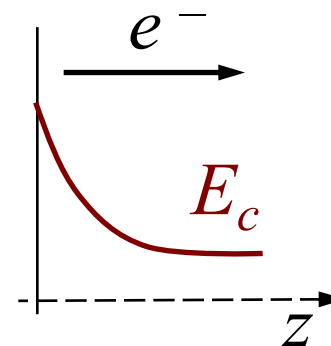
*Transient current mapping*



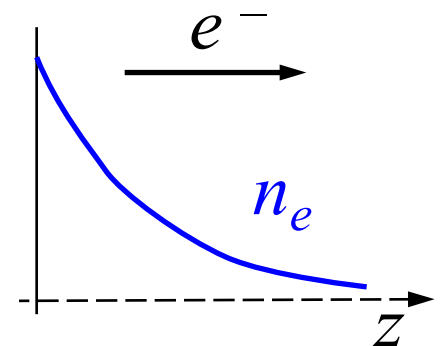
$$E \propto \frac{\partial I}{\partial t}$$

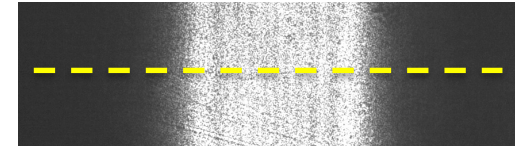
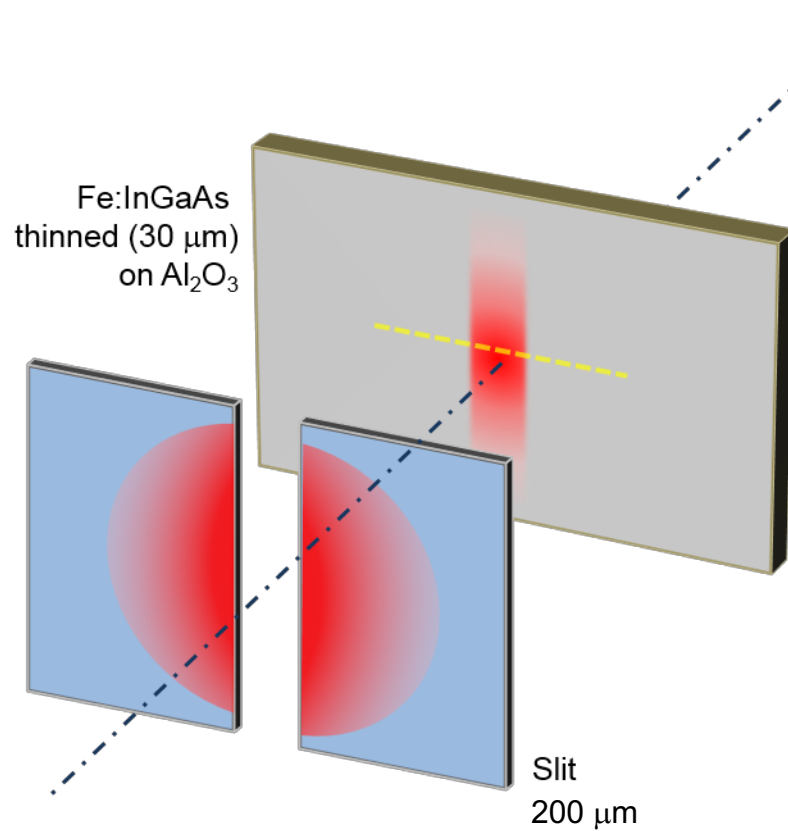
*Models:*

Built-in surface field

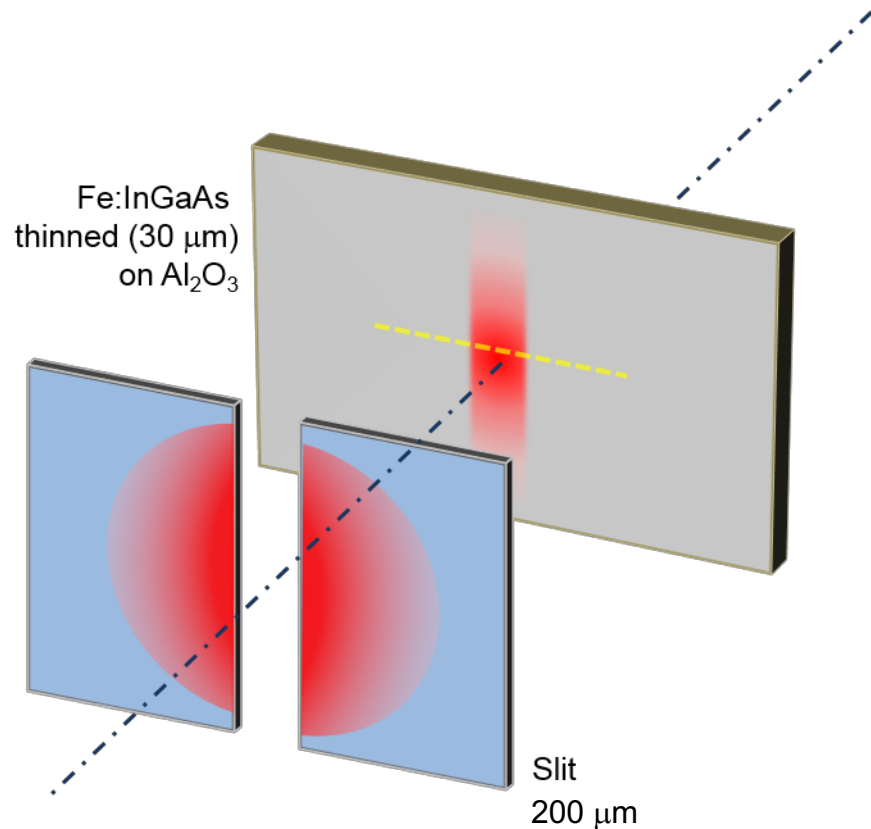


Carrier density gradient



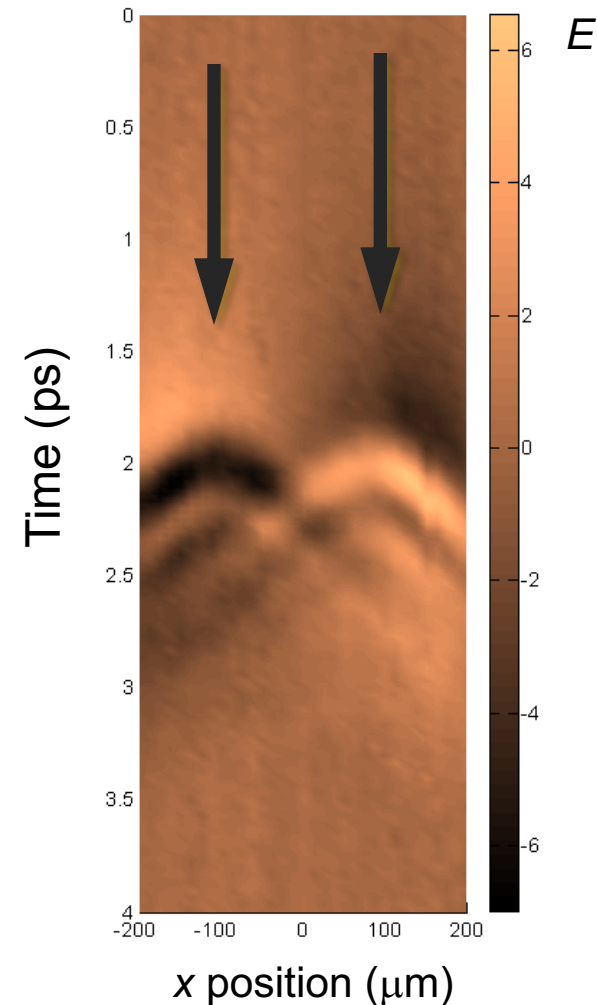
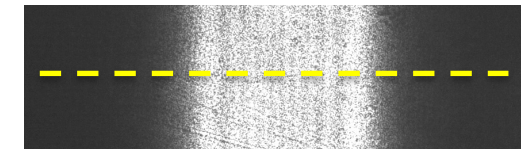


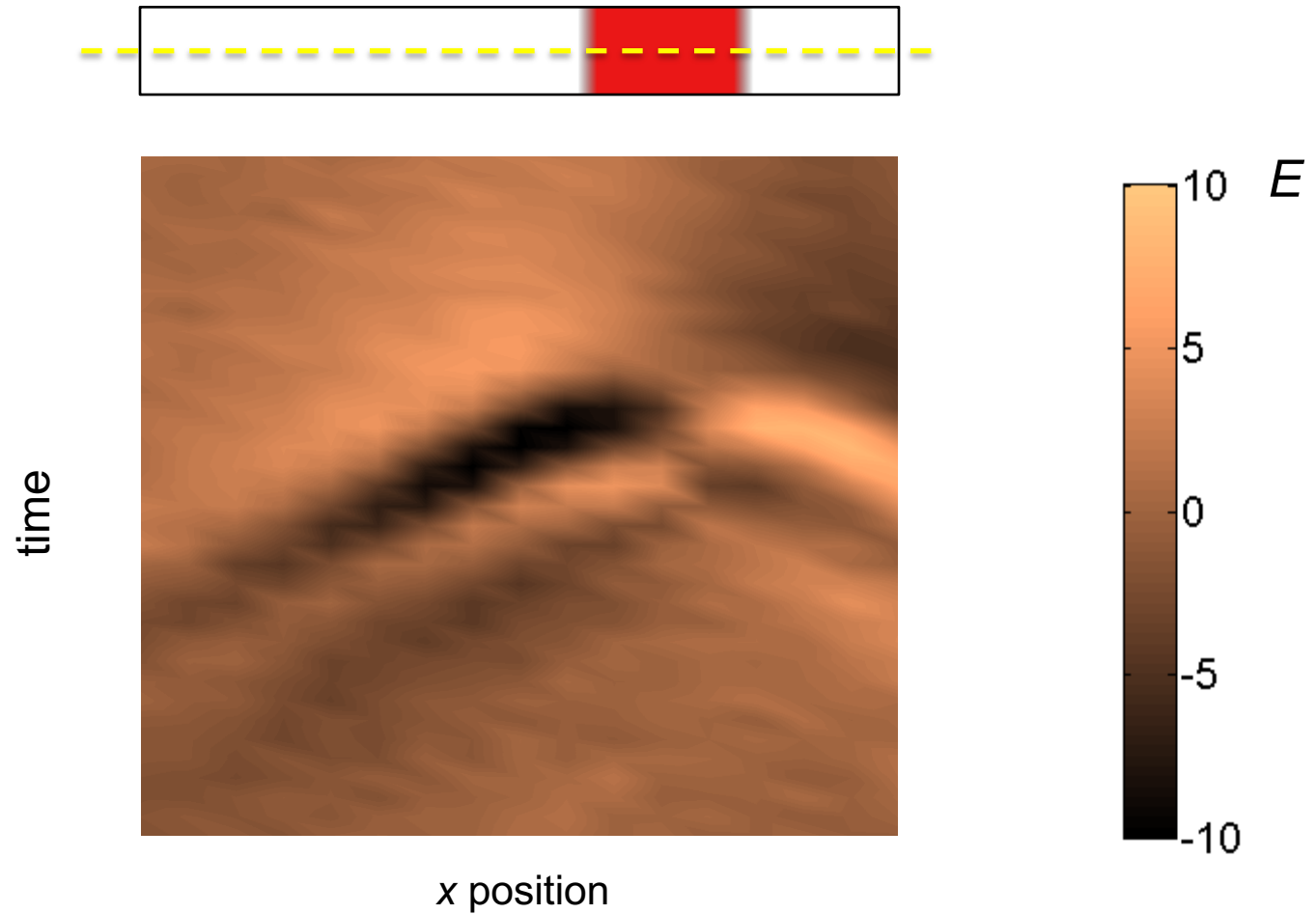
Mitrofanov et al.,  
*IEEE Trans. THz S&T* (2016)

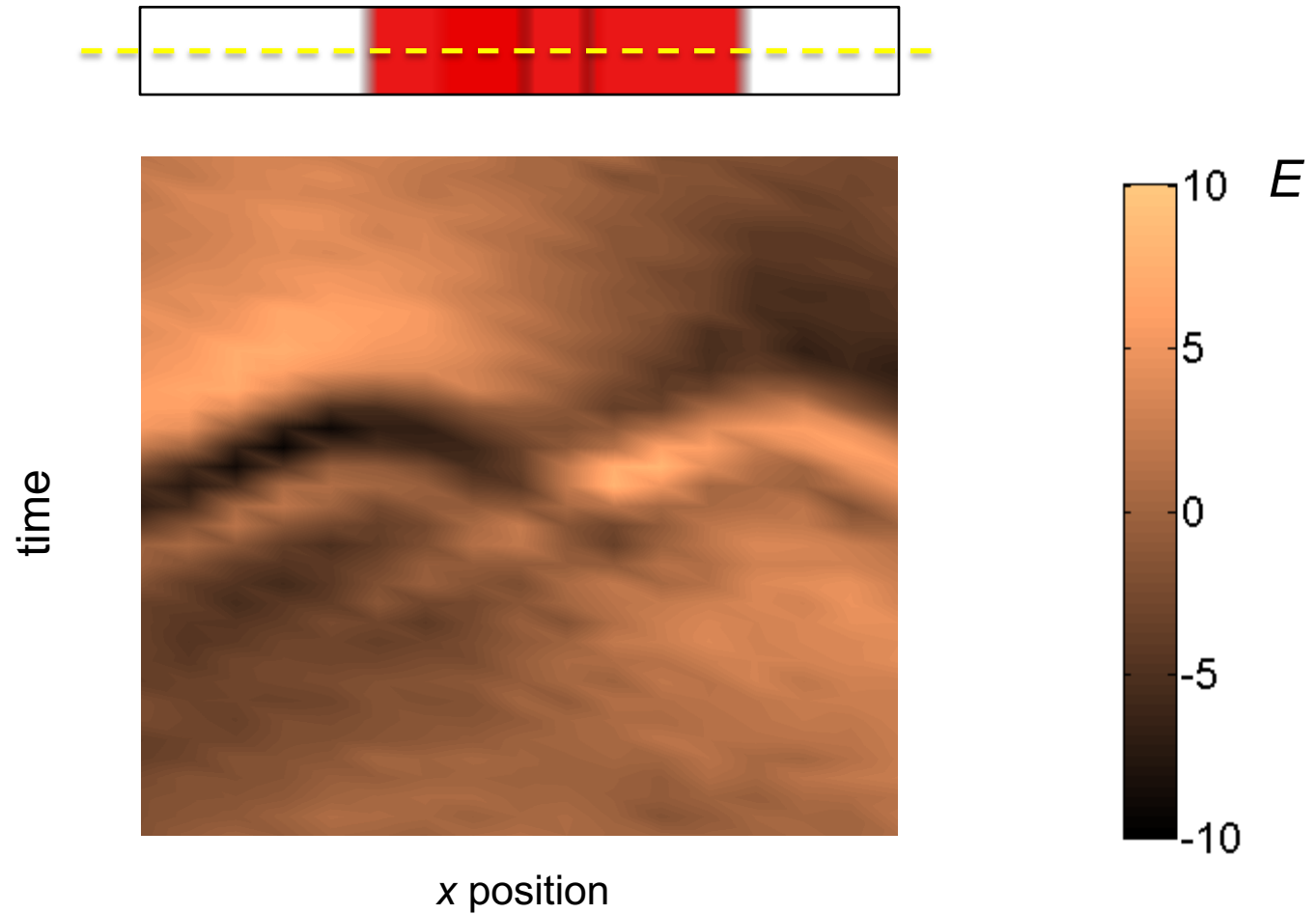


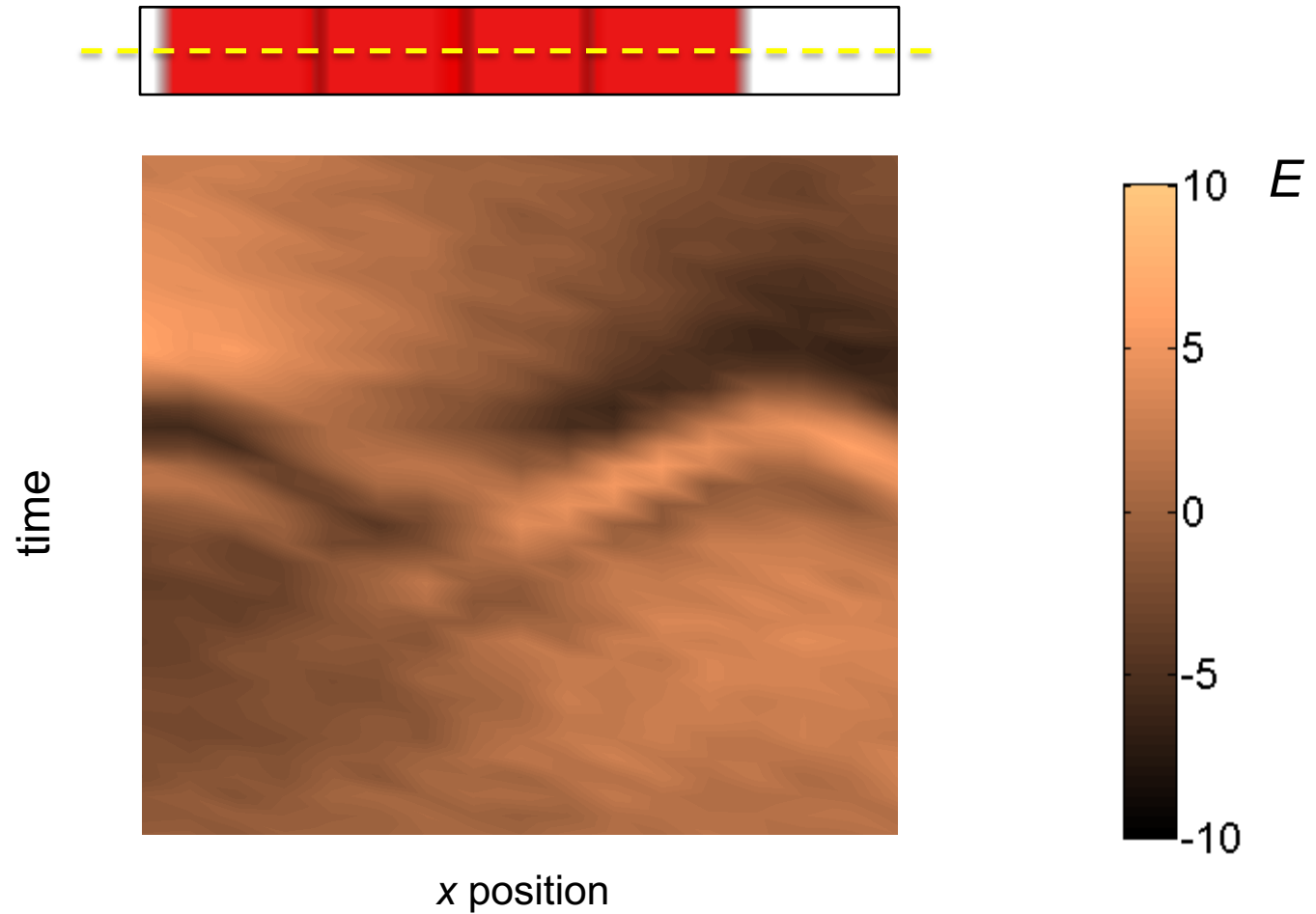
THz emission originates from two distinct points corresponding to the *slit edges*.

The two sources display opposite polarities leading to no emission in forward direction

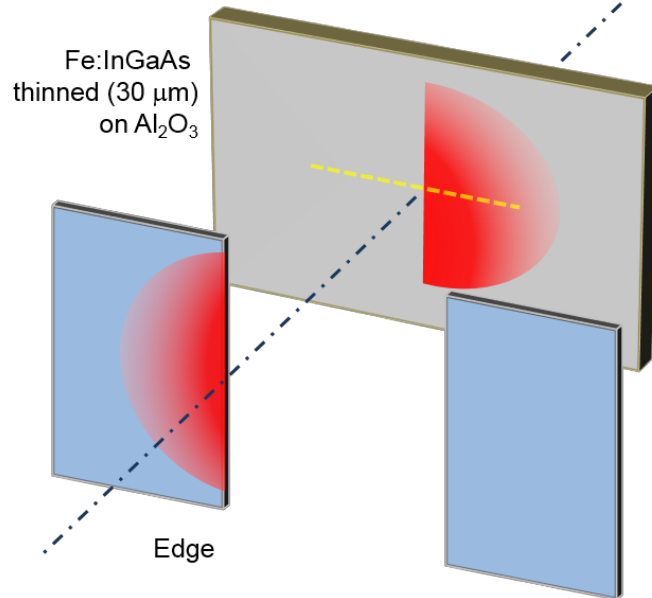




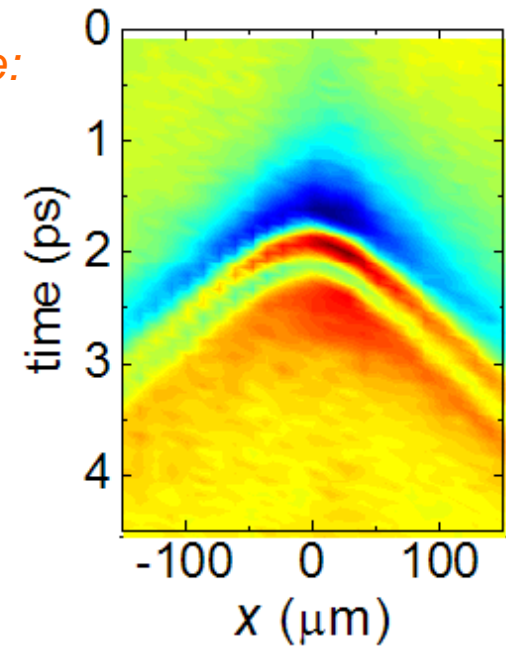




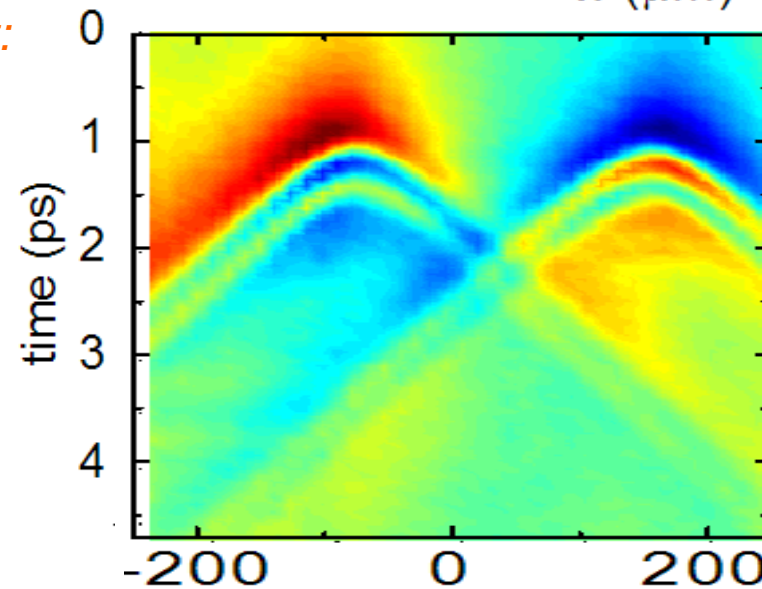


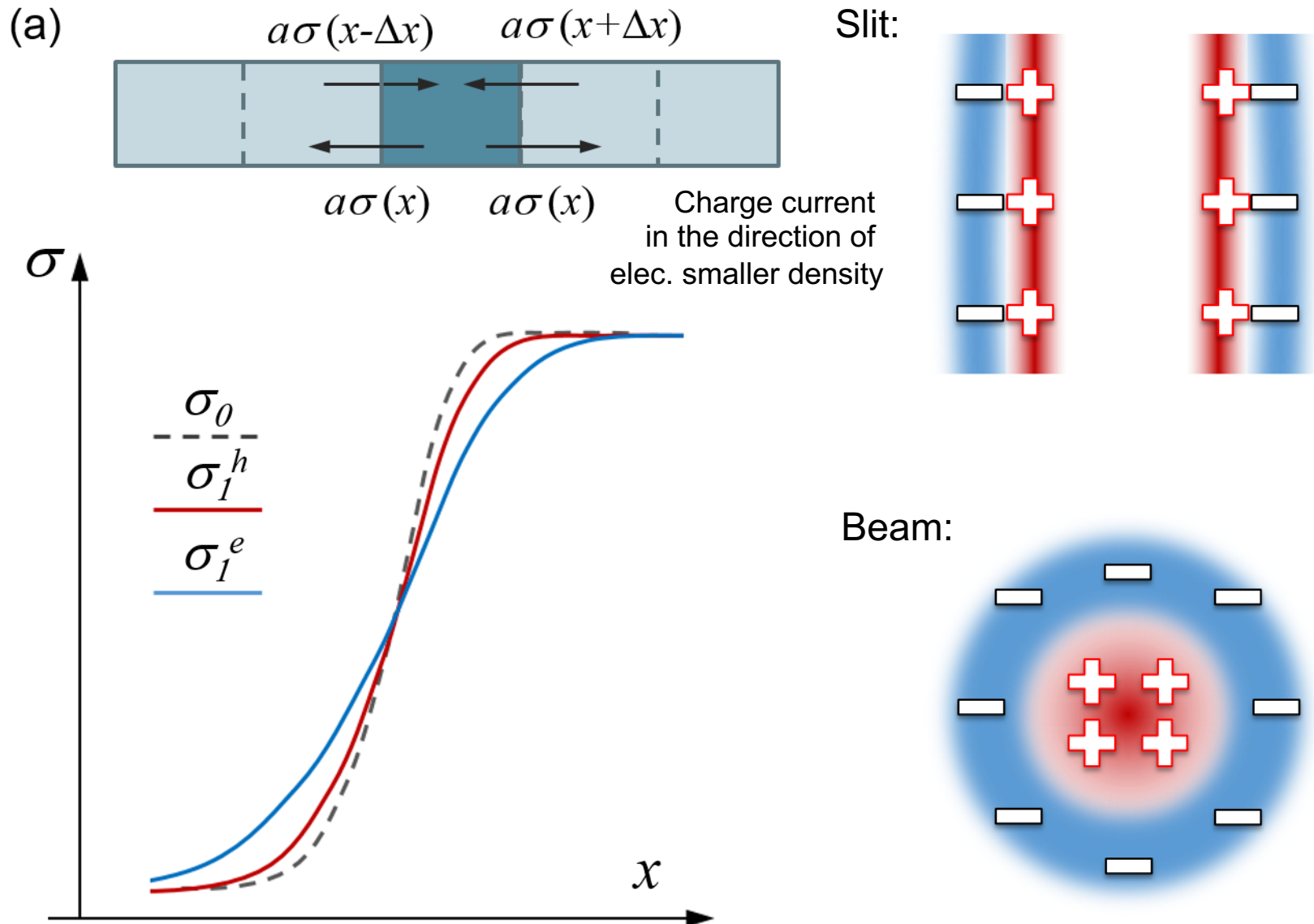


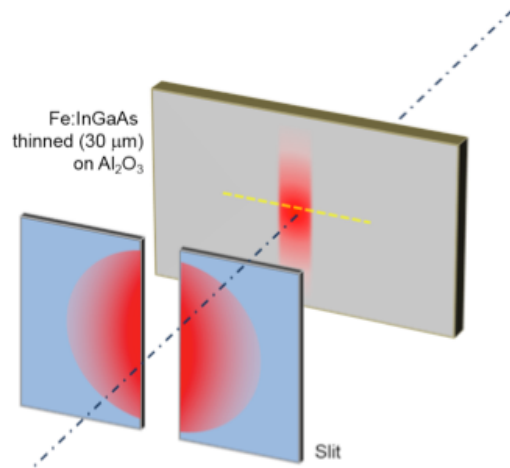
*Single edge:*



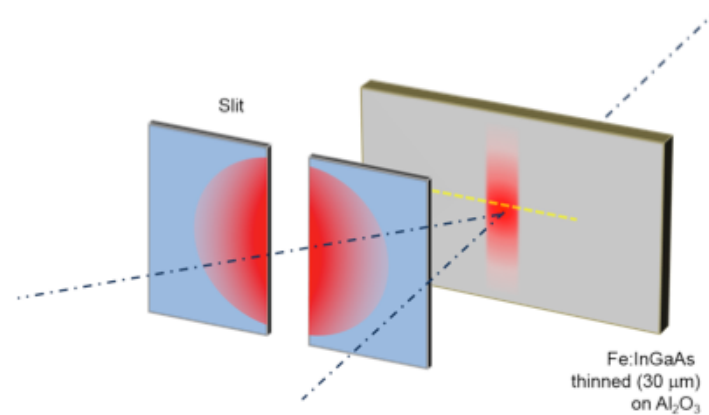
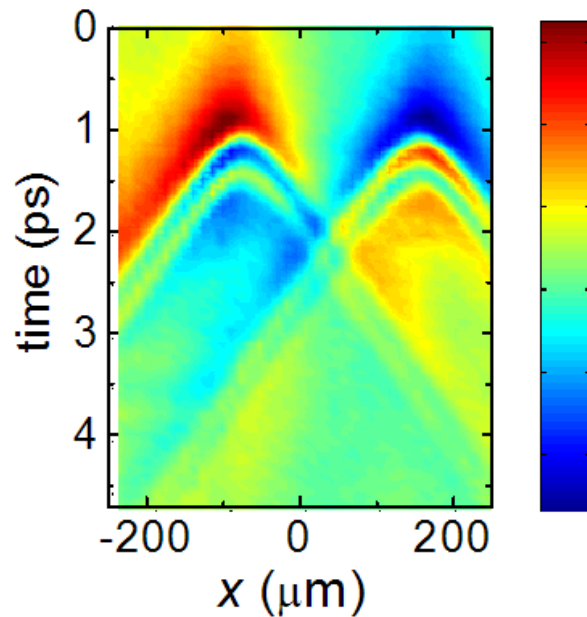
*Slit:*



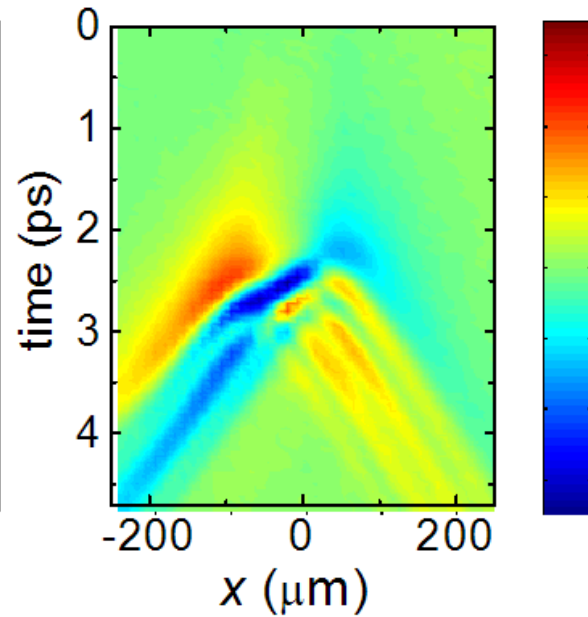




Normal incidence, 220  $\mu\text{m}$



45 deg., 150  $\mu\text{m}$



45 deg., 250  $\mu\text{m}$

